

Neutral-current SMEFT studies at the EIC

Kağan Şimşek

Northwestern
University

High-Energy Physics Seminar
Northwestern University
Monday April 18, 2022

Reference: 2204.07557

Collaborators: R. Boughezal
F. Petriello
D. Wiegand
X. Zheng *et al.* (JLab EIC Group)

Outline

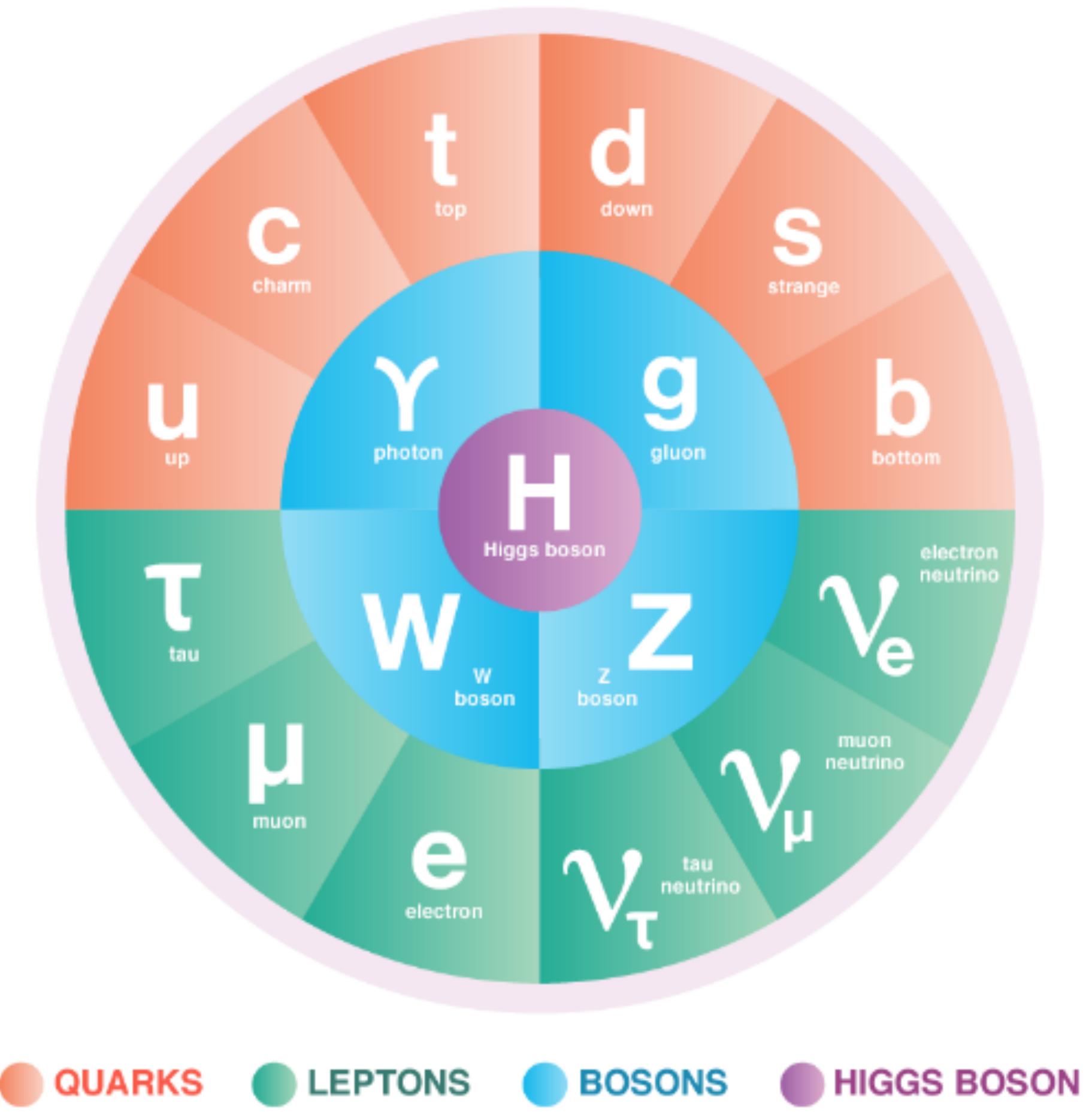
- Introduction
- Neutral-current DIS and SMEFT
- Projections of PV and LC asymmetry data
- Uncertainties
- Framework of the SMEFT analysis
- SMEFT fit results
- Conclusion

Introduction

Standard Model

SM of particle physics:

- Successful in all lab phenomena
- Spectrum completely observed
- No new particles observed so far



Standard Model

SM of particle physics:

- Successful in all lab phenomena
- Spectrum completely observed
- No new particles observed so far

Shortcomings:

- Dark matter
- Baryon-antibaryon asymmetry
- Neutrino mass
- Hierarchy problem

EIC

An experimental program under design: the Electric-Ion Collider (EIC)

- Brookhaven National Laboratory, Upton, NY
- electron + proton/nuclei collisions



bnl.gov

EIC

An experimental program under design: the Electric-Ion Collider (EIC)

- Brookhaven National Laboratory, Upton, NY
- electron + proton/nuclei collisions

Unique features:

- Energy between fixed-target scattering and high-energy colliders
- Luminosity orders of magnitude higher than HERA
- First lepton-ion collider to polarize both beams
- First collider with fast spin-flip capacity

EIC

From these unique features:

- Extraction of $A_{\text{PV}}^{(\ell)}$ and $A_{\text{PV}}^{(H)}$ in EW NC cross section
- Reduced uncertainties from luminosity and detector acceptance/efficiency
- Explore issues in QCD
- Probes of BSM physics

EIC

From these unique features:

- Extraction of $A_{\text{PV}}^{(\ell)}$ and $A_{\text{PV}}^{(H)}$ in EW NC cross section
- Reduced uncertainties from luminosity and detector acceptance/efficiency
- Explore issues in QCD
- Probes of BSM physics

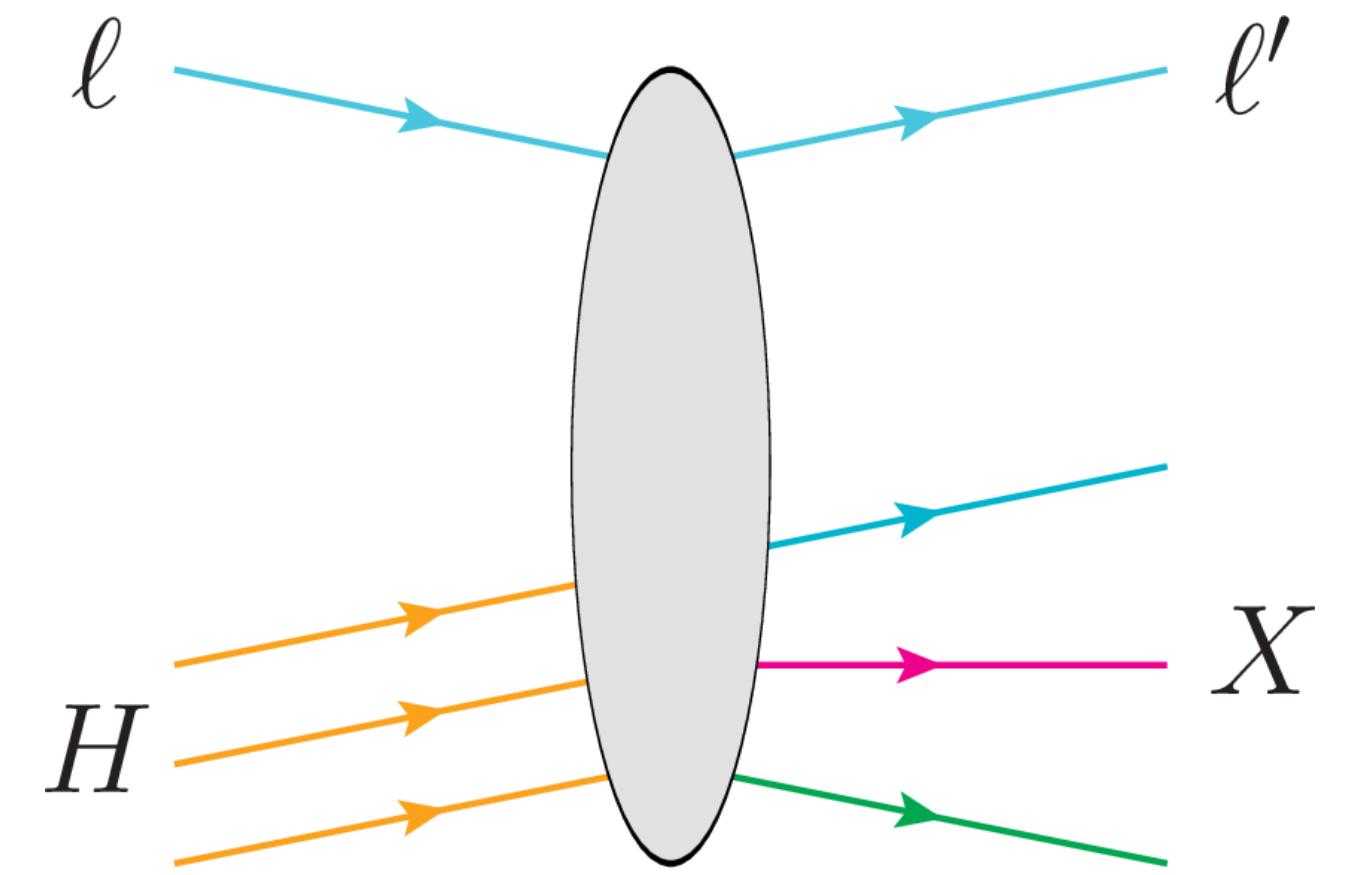
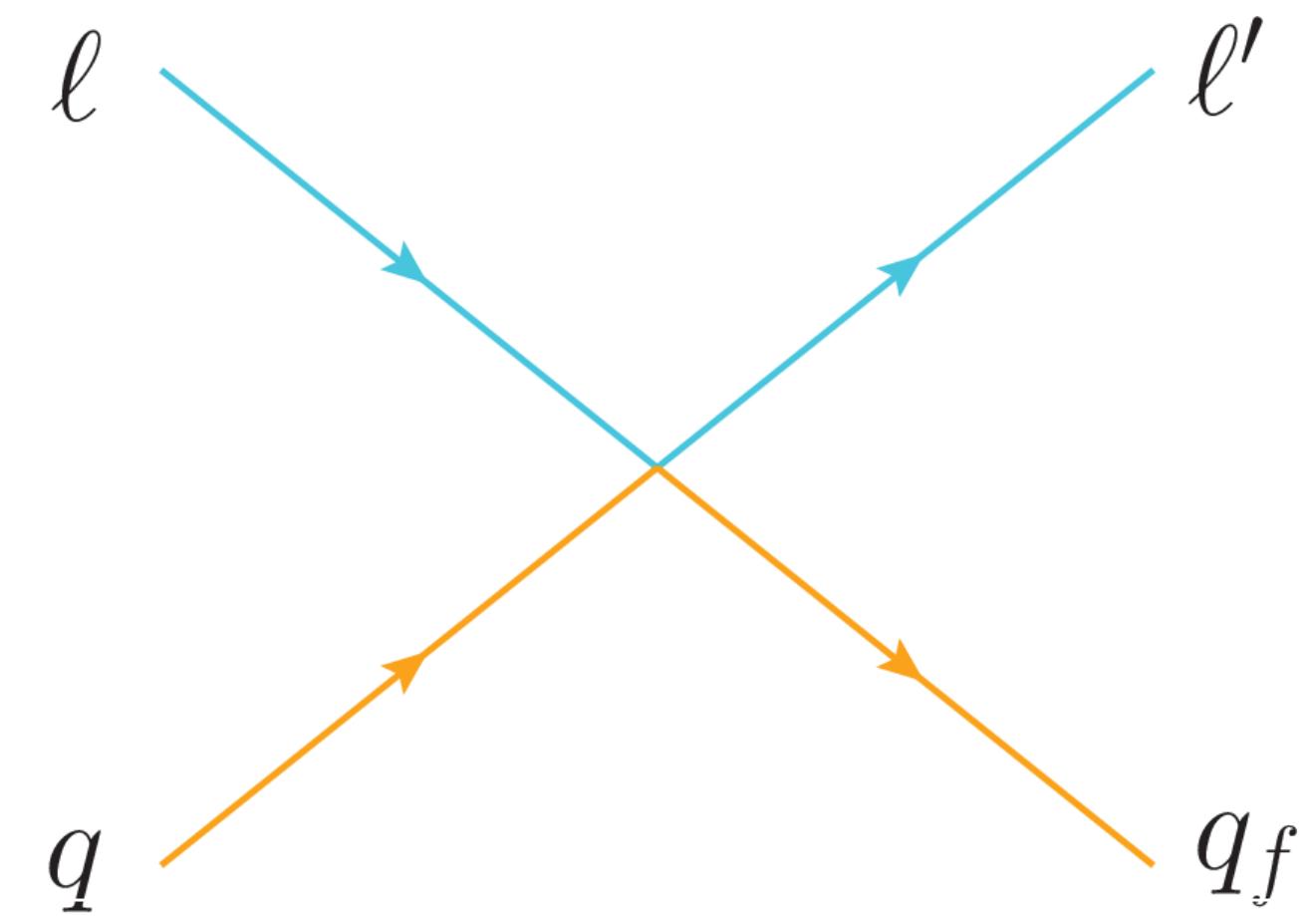
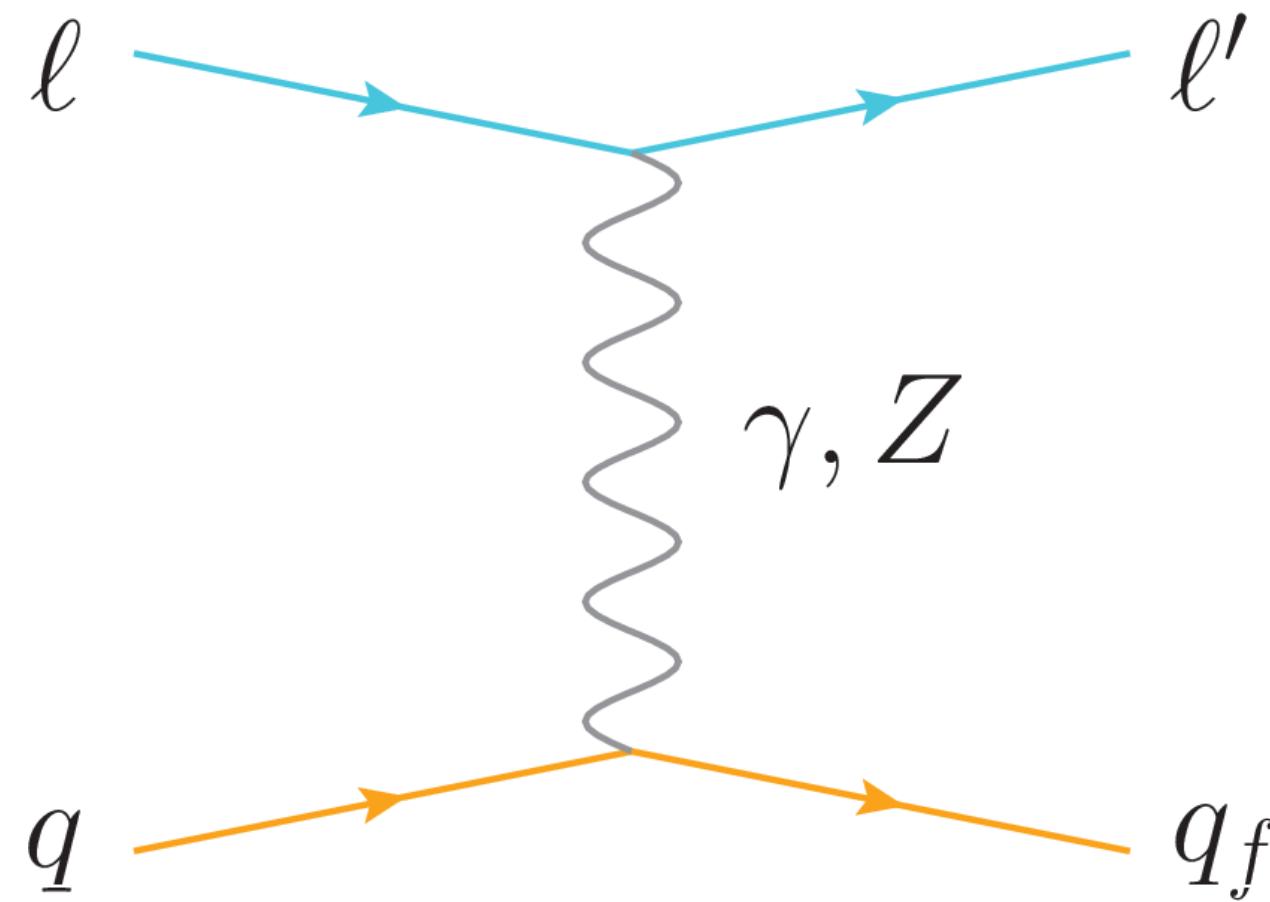
What's new:

- Positron beam in the future: $A_{\text{LC}}^{(H)}$

Neutral-current DIS
and
SMEFT

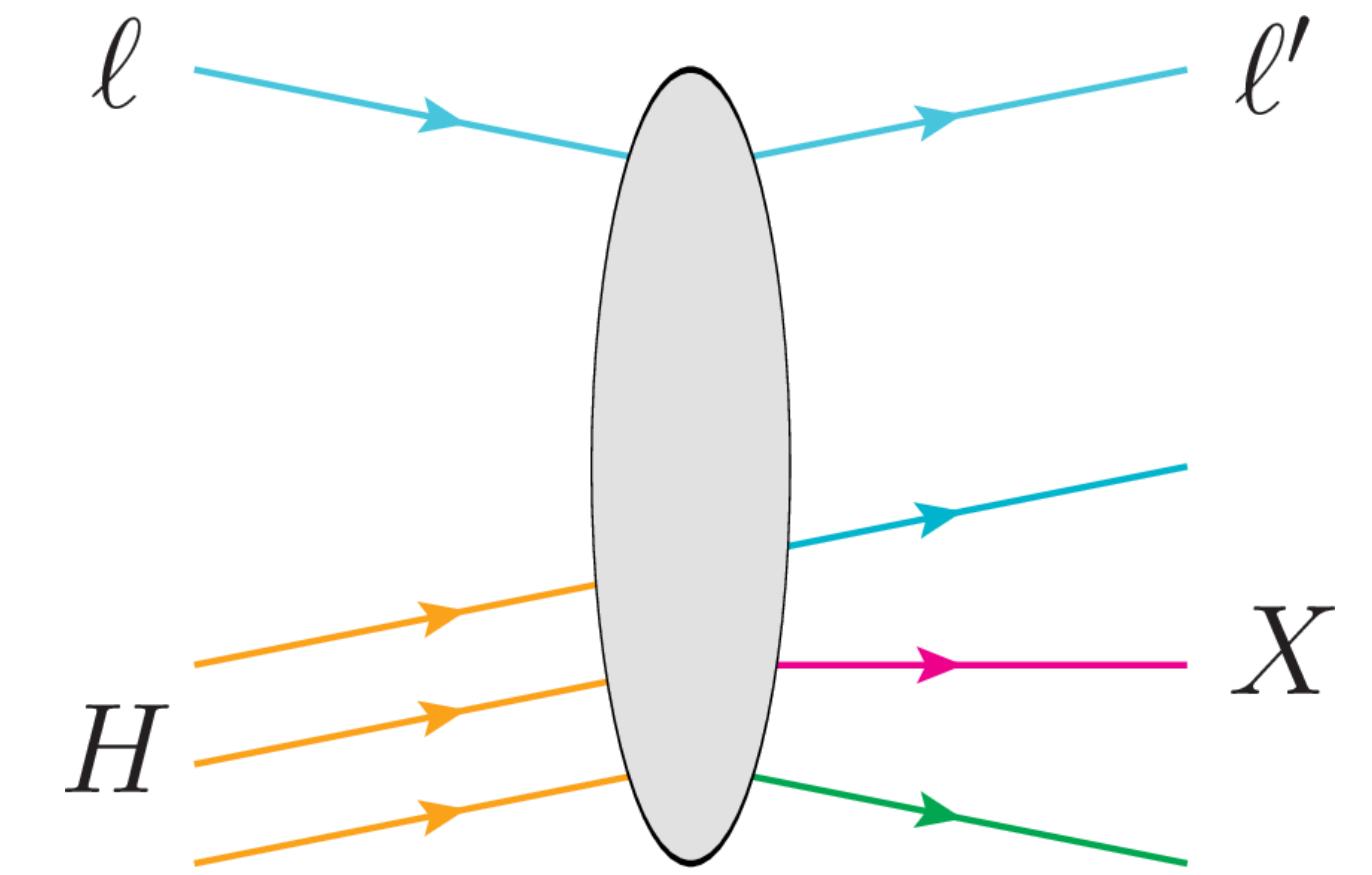
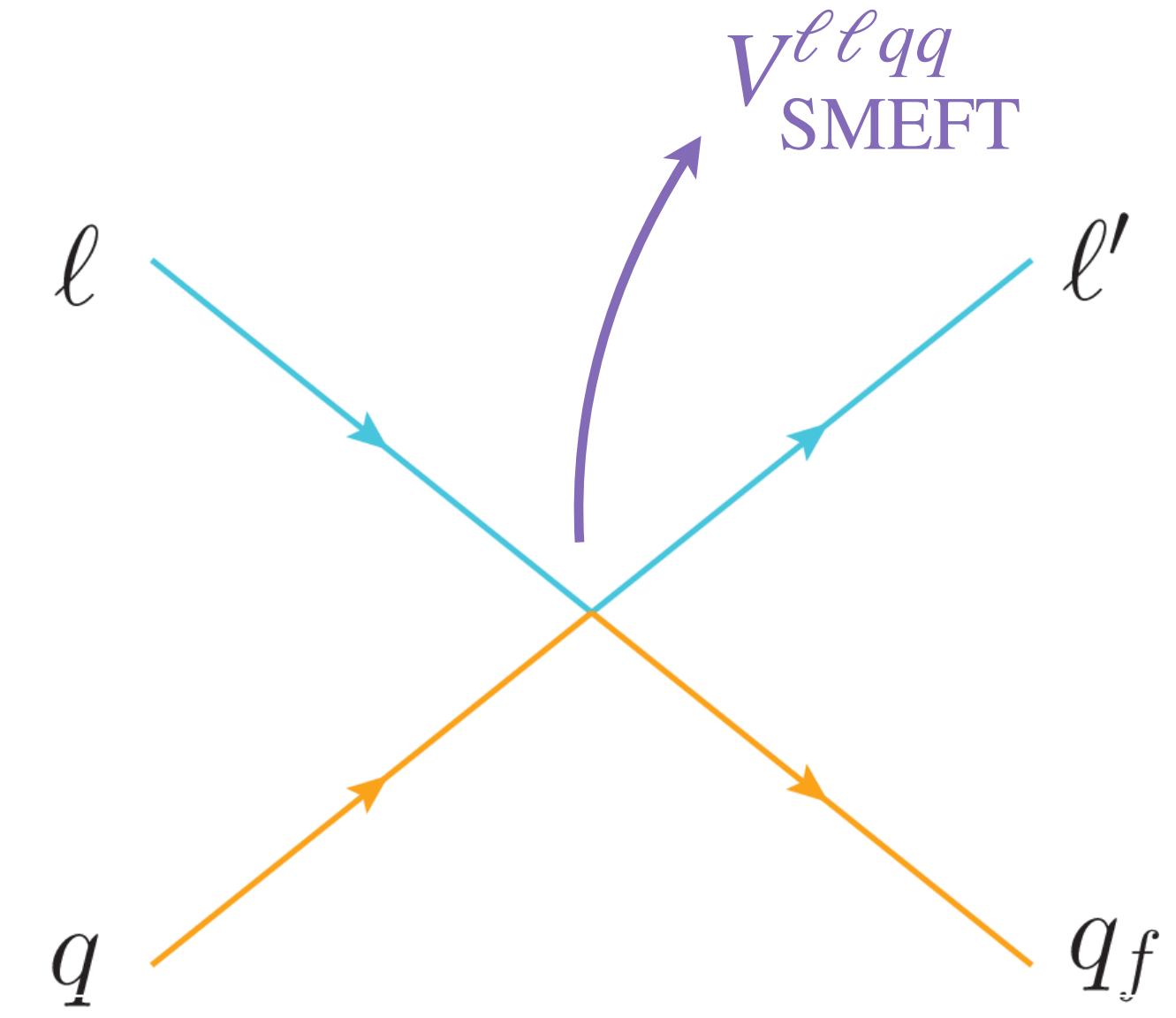
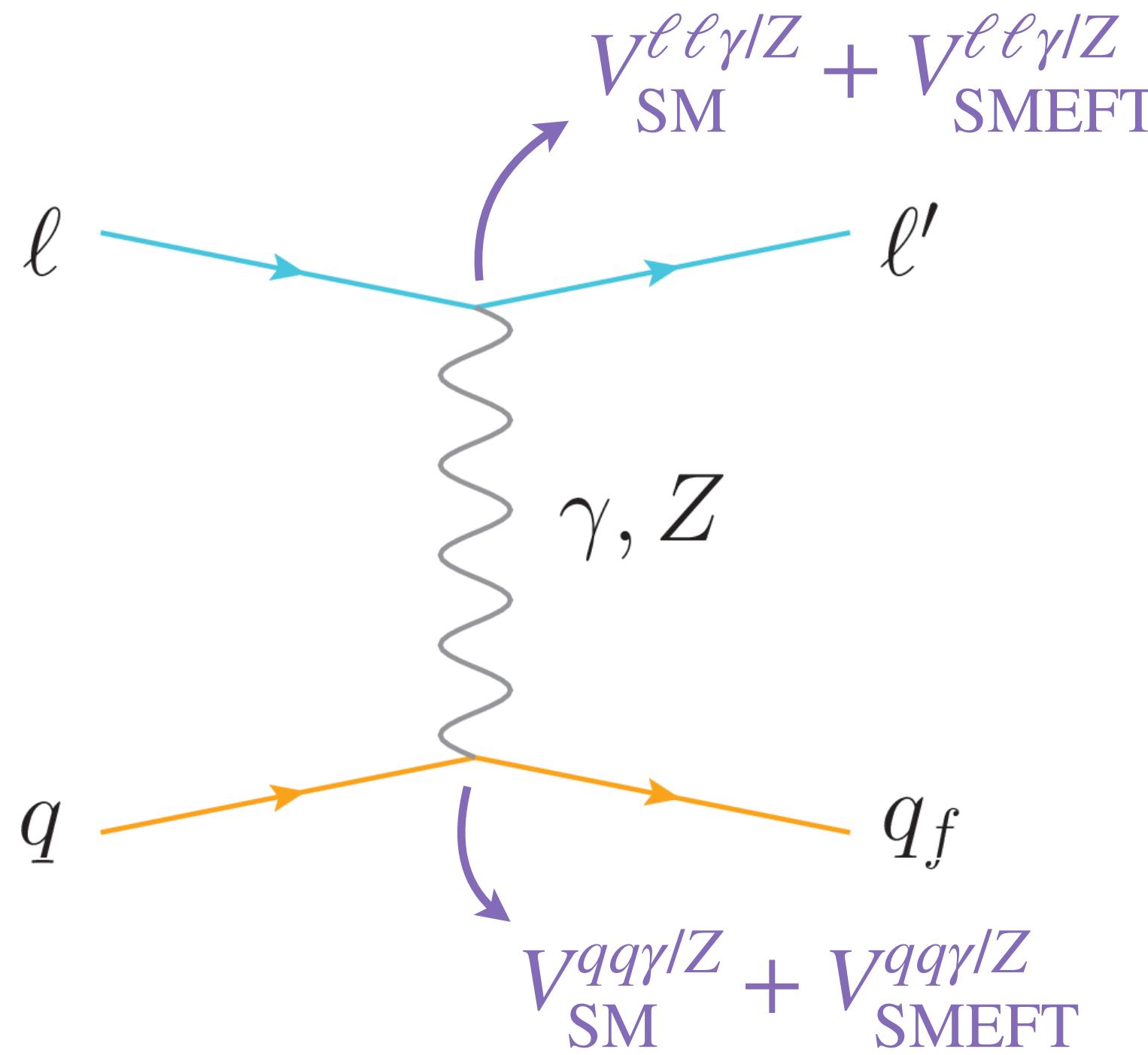
NC DIS and SMEFT

NC DIS in the process $\ell + H \rightarrow \ell' + X$:



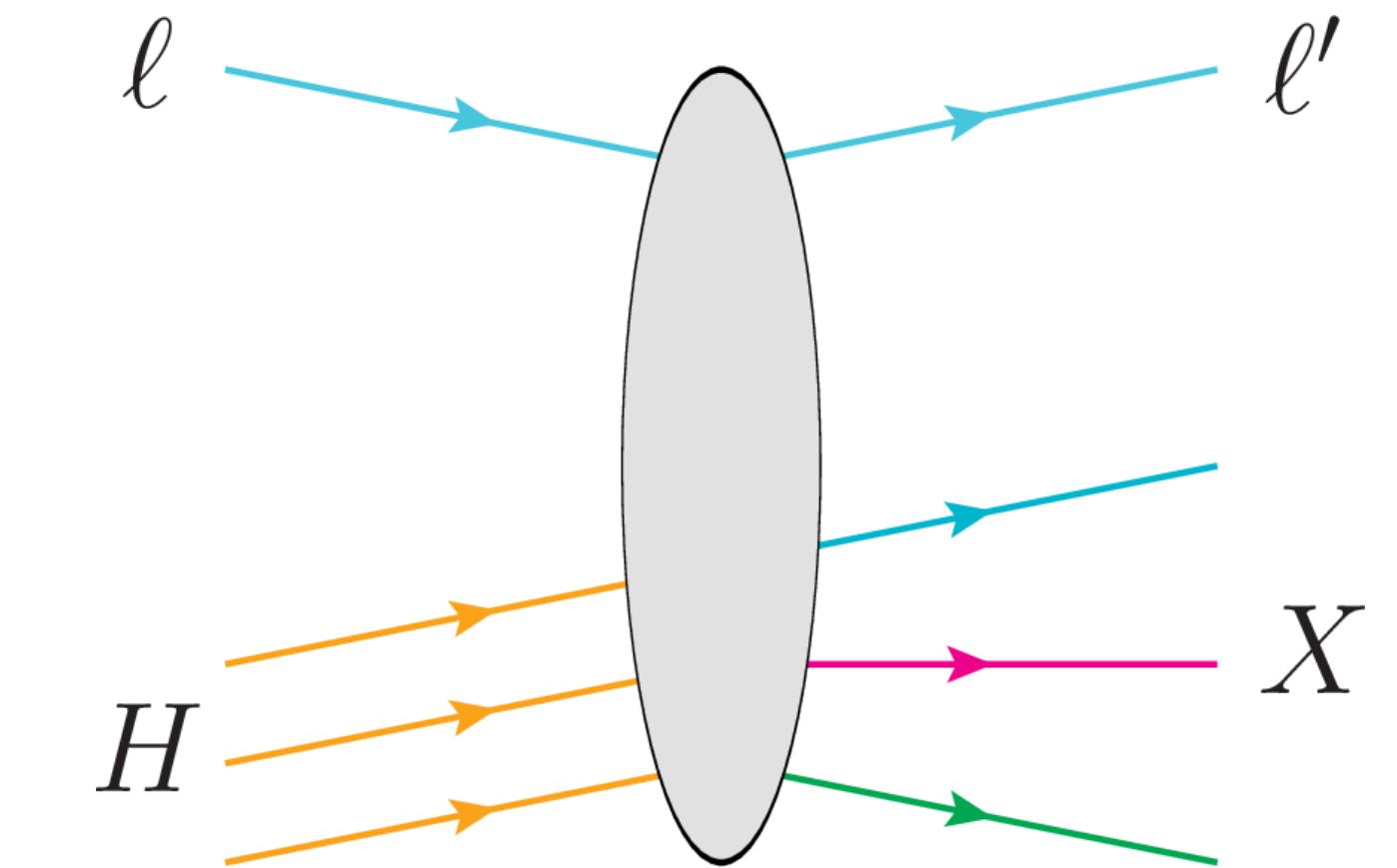
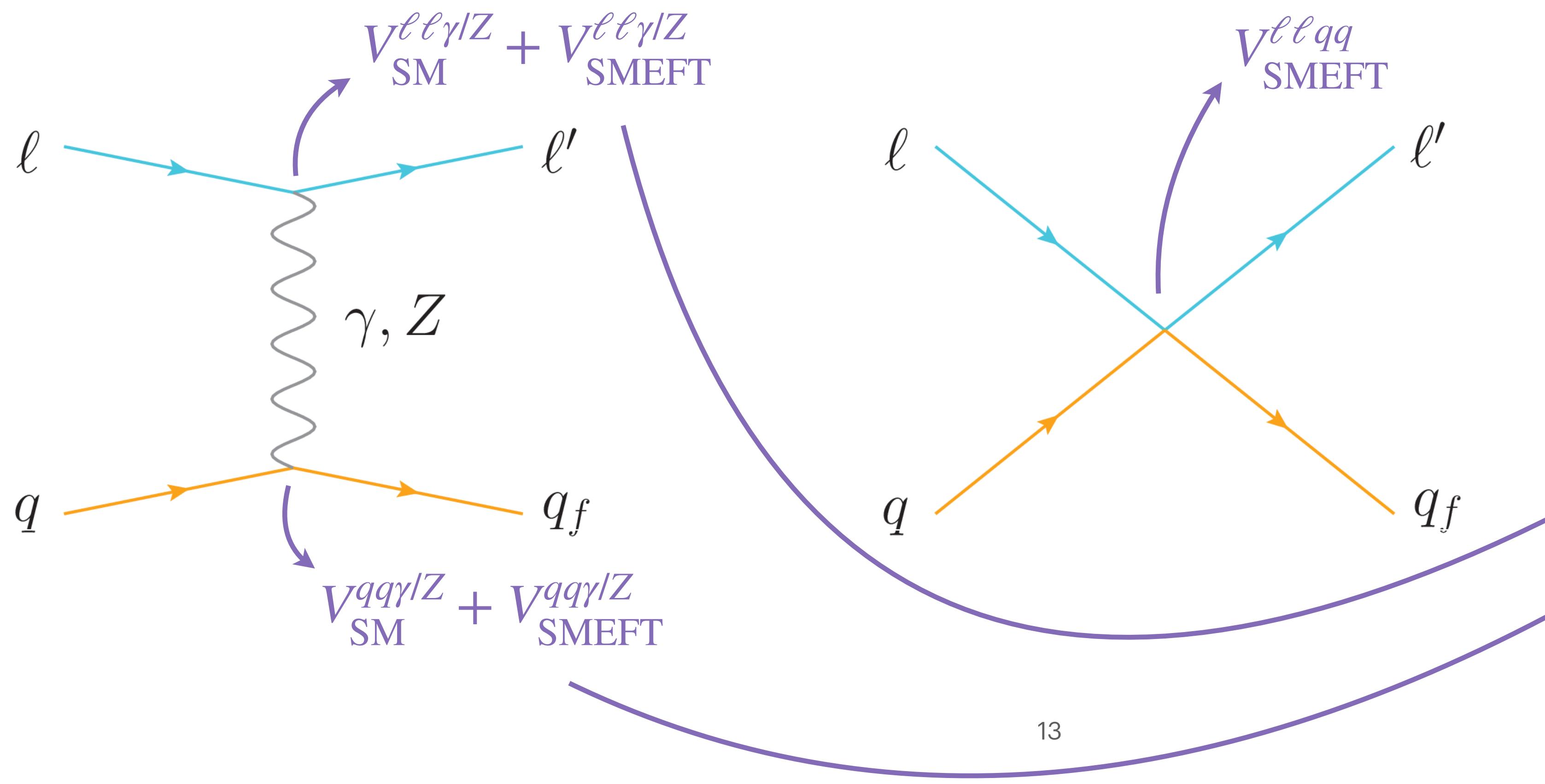
NC DIS and SMEFT

NC DIS in the process $\ell + H \rightarrow \ell' + X$:



NC DIS and SMEFT

NC DIS in the process $\ell + H \rightarrow \ell' + X$:



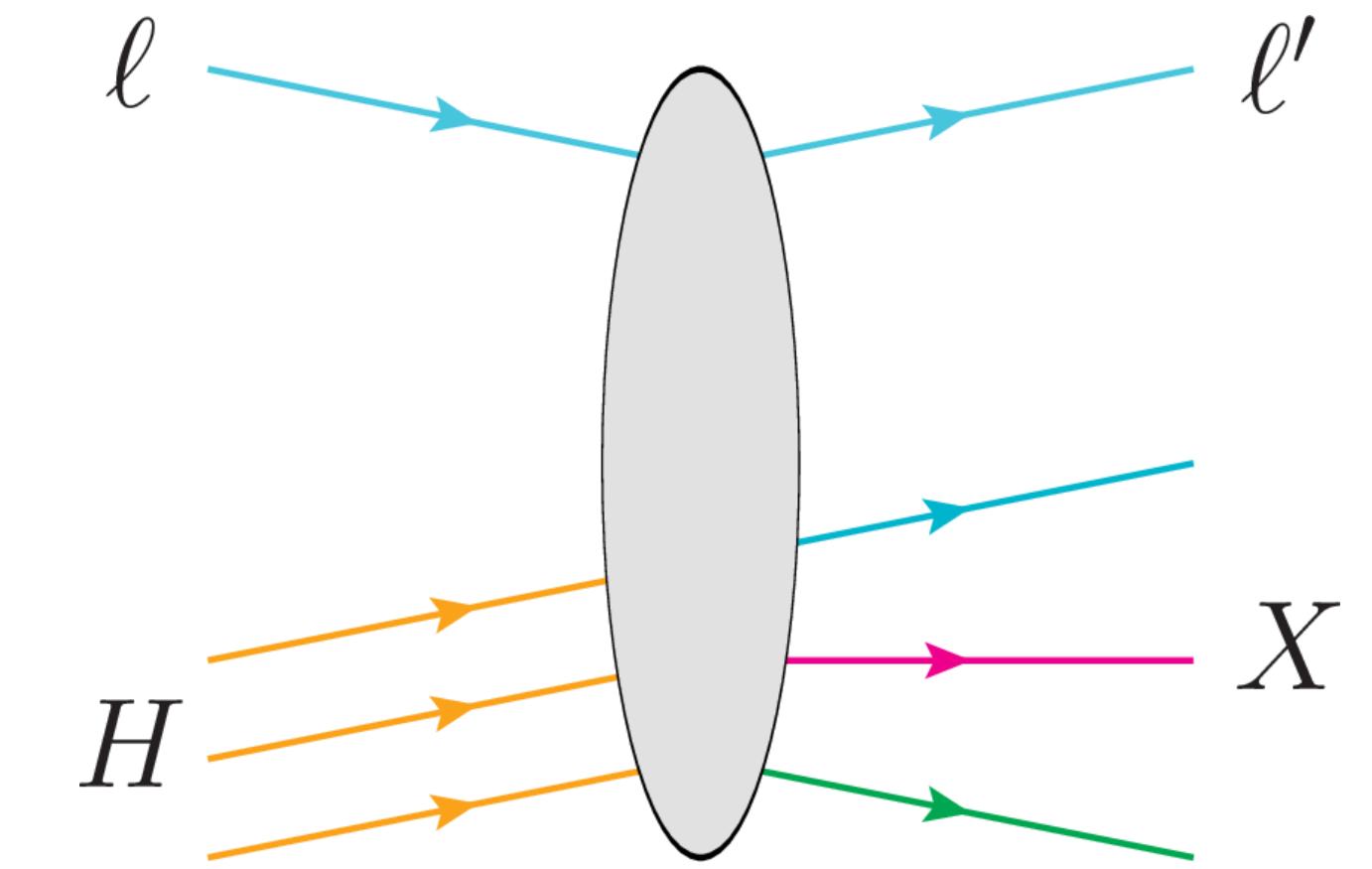
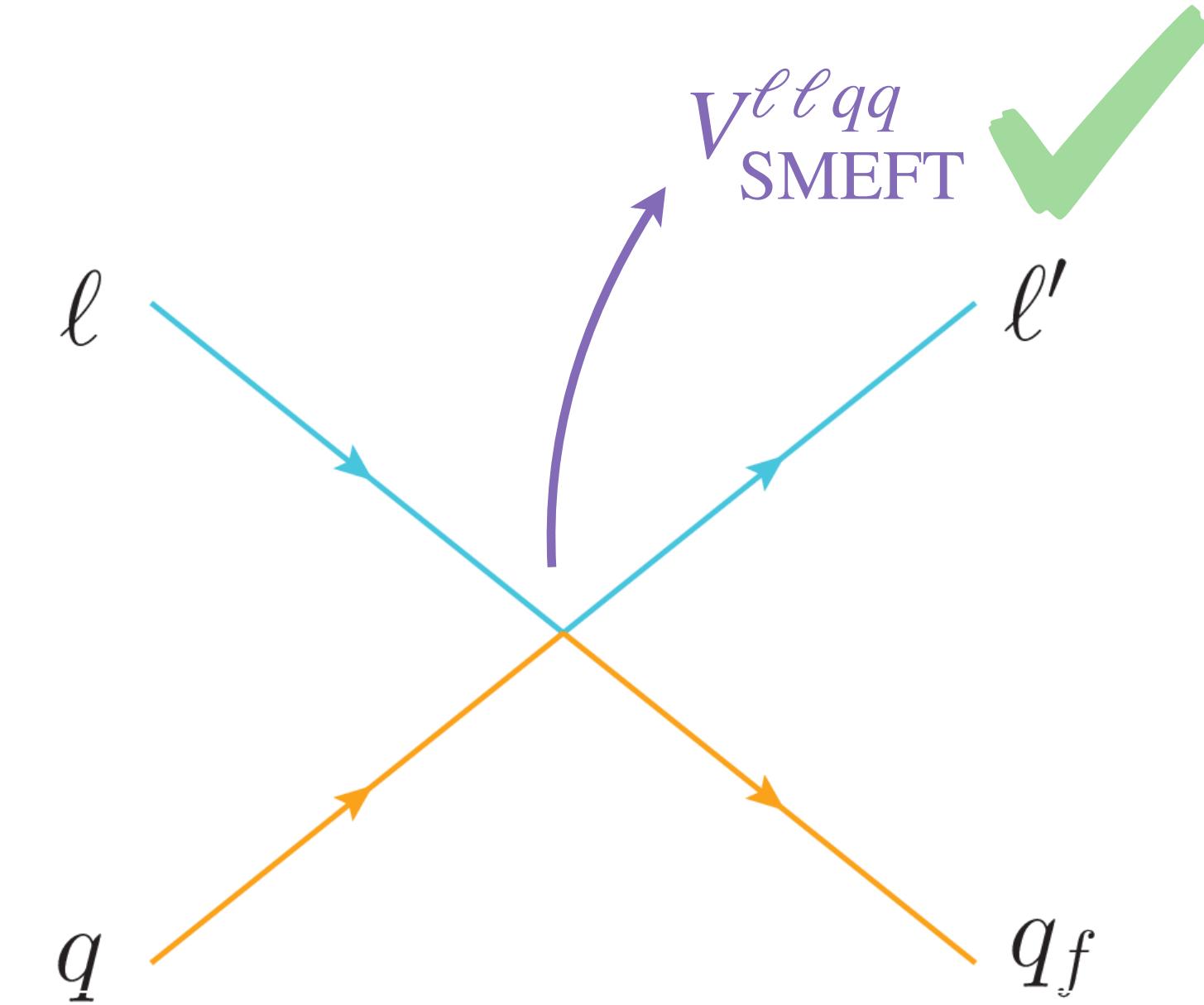
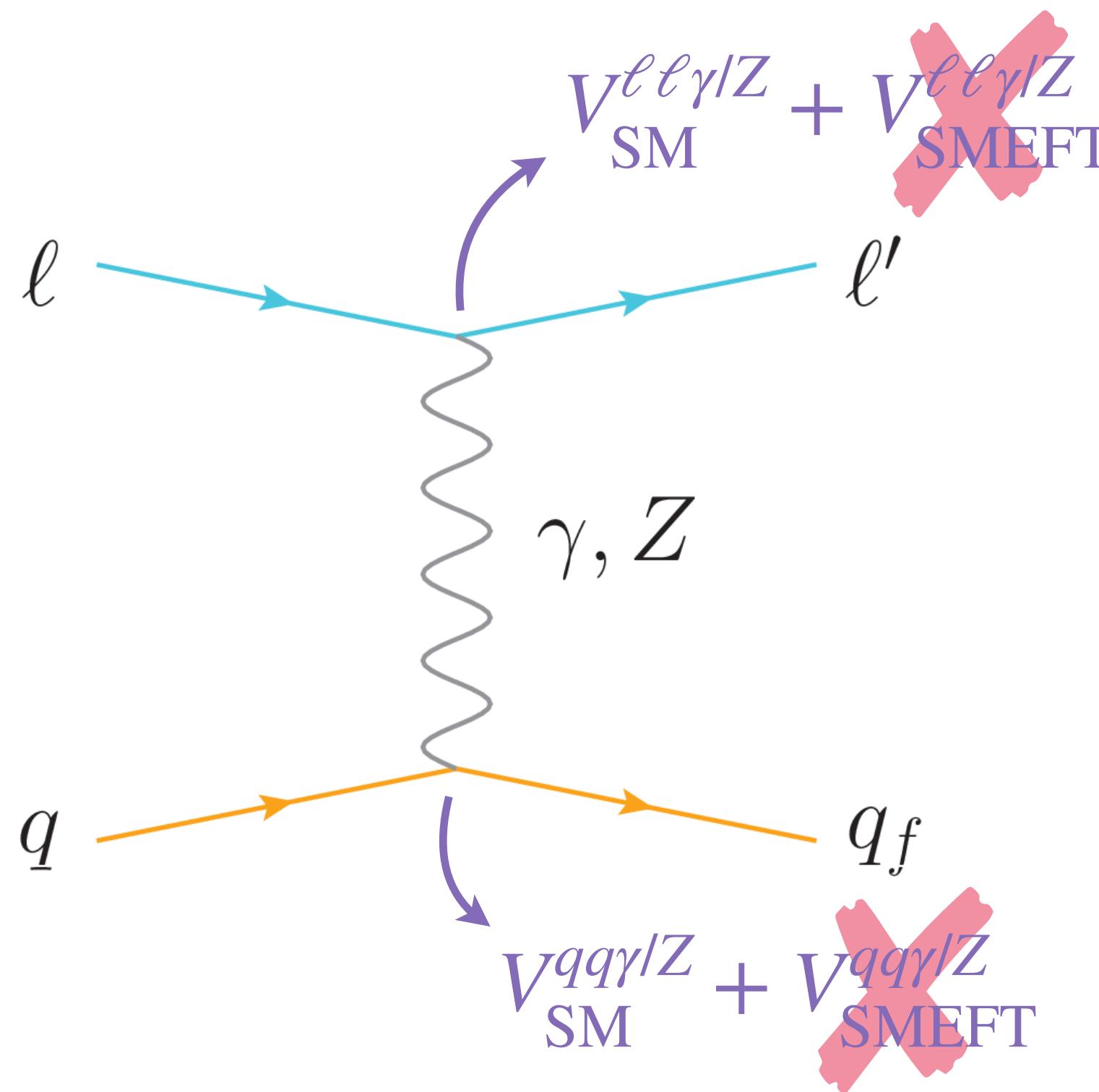
C	95 % CL, $\Lambda = 1$ TeV
$C_{\varphi WB}$	$[-0.0088, 0.0013]$
$C_{\varphi u}$	$[-0.072, 0.091]$
$C_{\varphi q}^{(3)}$	$[-0.011, 0.014]$
$C_{\varphi q}^{(1)}$	$[-0.027, 0.043]$
$C_{\varphi \ell}^{(3)}$	$[-0.012, 0.0029]$
$C_{\varphi \ell}^{(1)}$	$[-0.0043, 0.012]$
$C_{\varphi e}$	$[-0.013, 0.0094]$
$C_{\varphi d}$	$[-0.16, 0.060]$

from Z and W pole observables

Dawson, Giardino 1909.02000

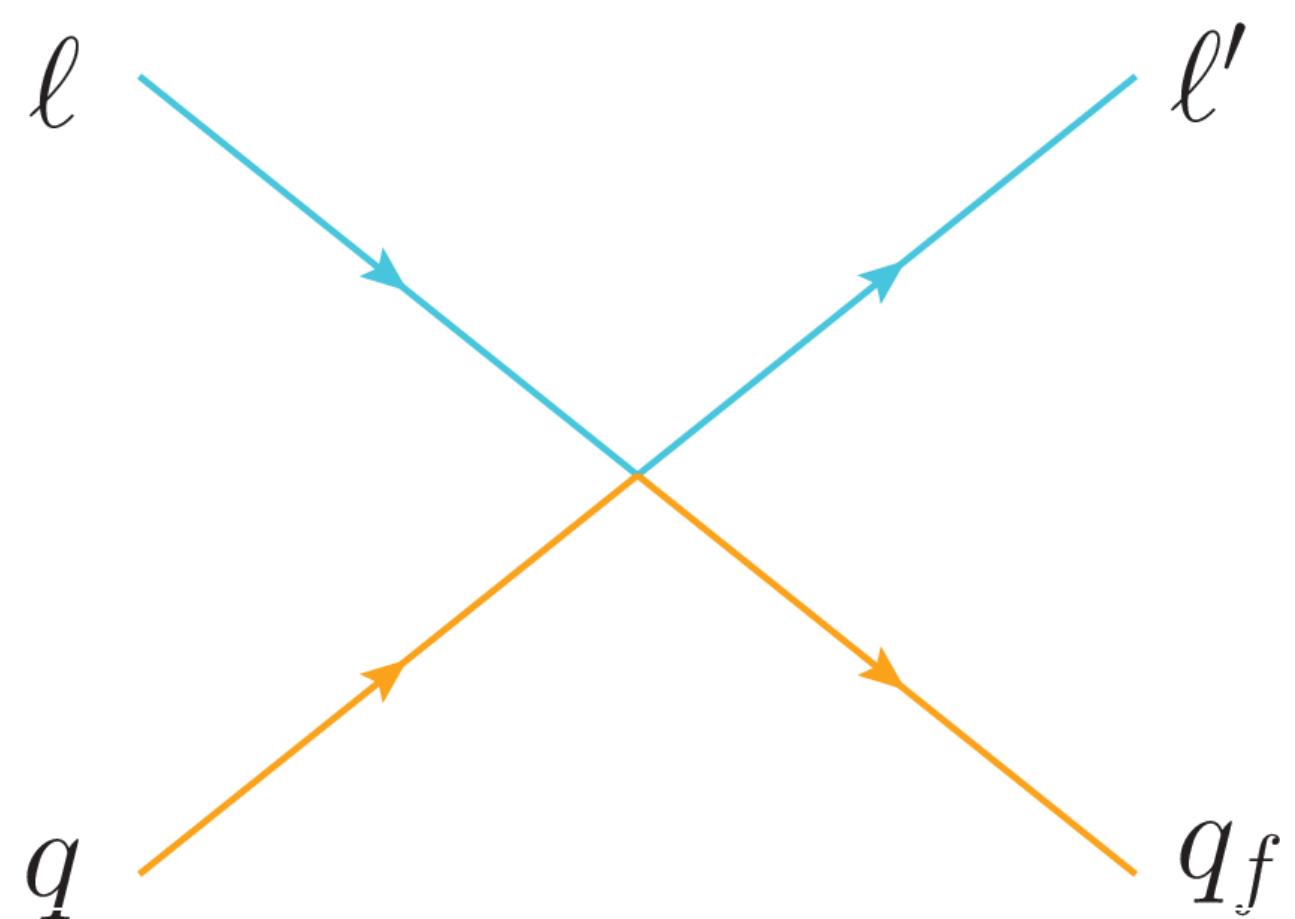
NC DIS and SMEFT

NC DIS in the process $\ell + H \rightarrow \ell' + X$:

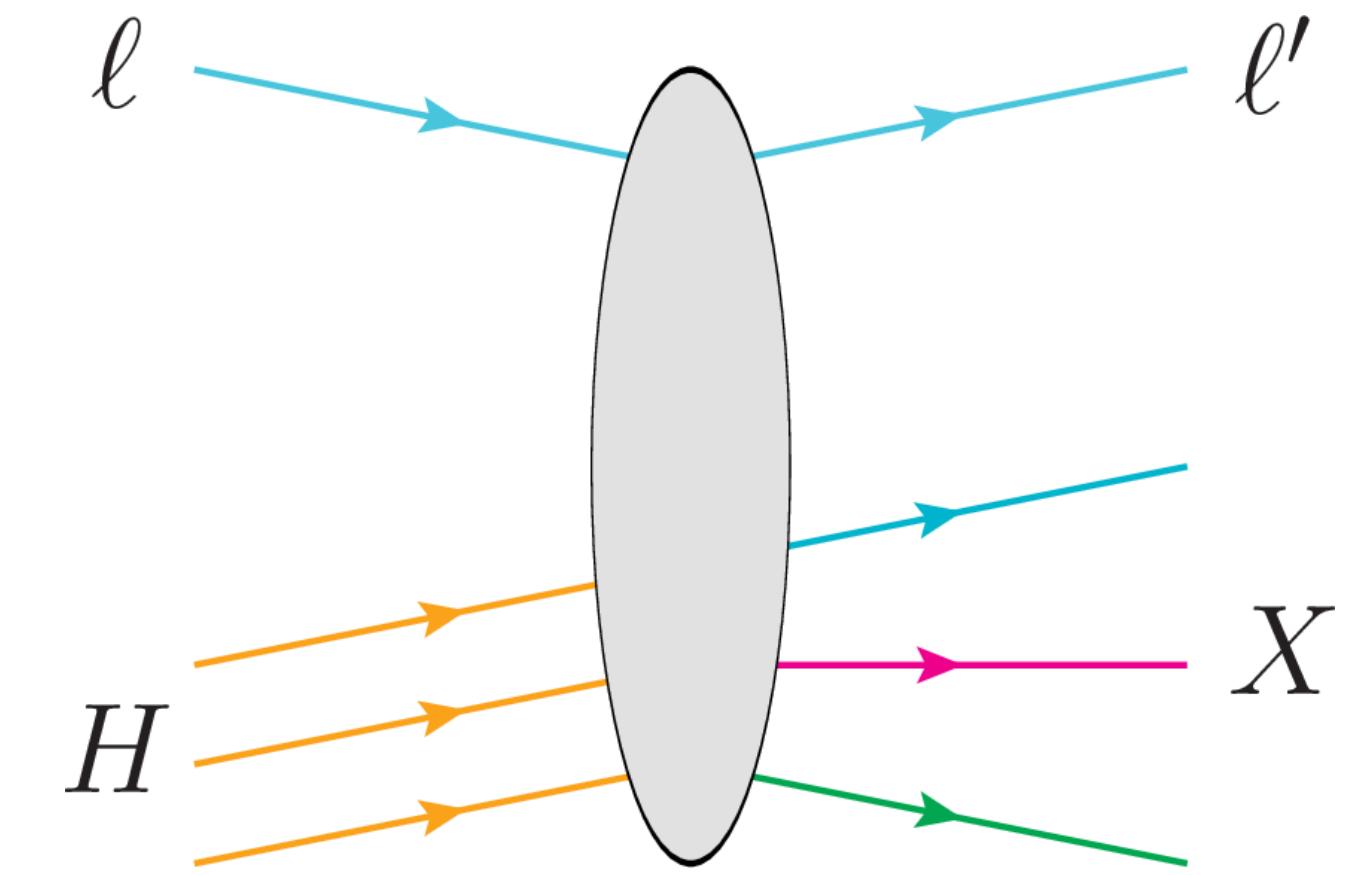


NC DIS and SMEFT

NC DIS in the process $\ell + H \rightarrow \ell' + X$:



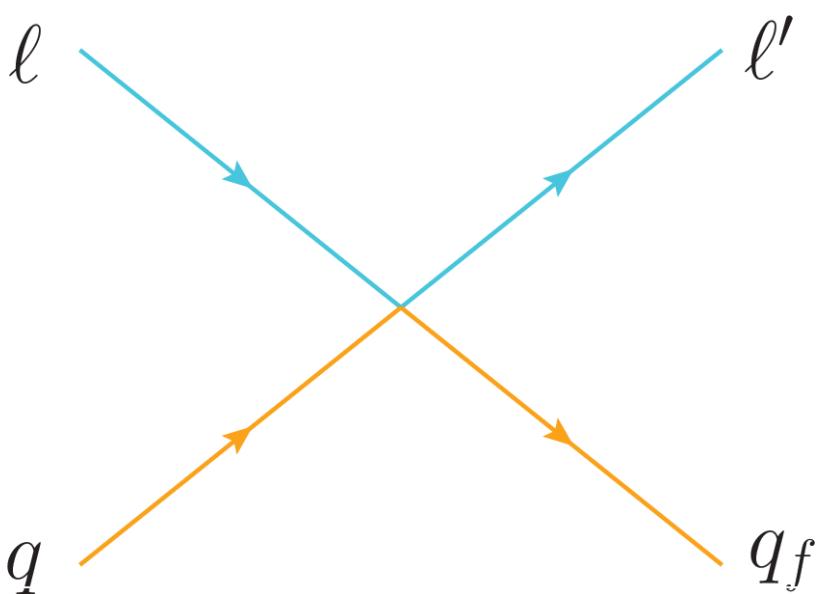
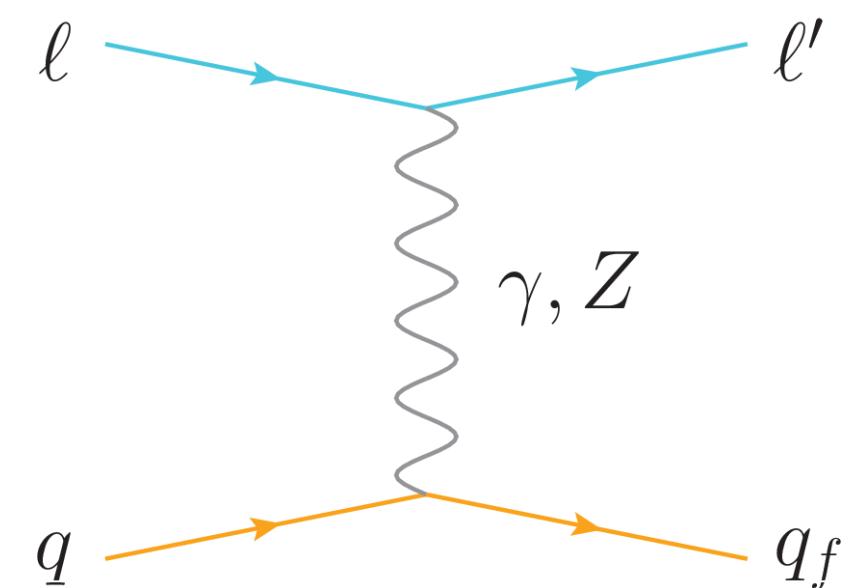
$$\mathcal{L}_{\text{SMEFT}} = \frac{1}{\Lambda^2} \sum_r C_r O_r$$



C_r	O_r
$C_{\ell q}^{(1)}$	$O_{\ell q}^{(1)} = [\bar{\ell} \gamma^\mu \ell][\bar{q} \gamma_\mu q]$
$C_{\ell q}^{(3)}$	$O_{\ell q}^{(3)} = [\bar{\ell} \gamma^\mu \tau^I \ell][\bar{q} \gamma_\mu \tau^I q]$
C_{eu}	$O_{eu} = [\bar{e} \gamma^\mu e][\bar{u} \gamma_\mu u]$
C_{ed}	$O_{ed} = [\bar{e} \gamma^\mu e][\bar{d} \gamma_\mu d]$
$C_{\ell u}$	$O_{\ell u} = [\bar{\ell} \gamma^\mu \ell][\bar{u} \gamma_\mu u]$
$C_{\ell d}$	$O_{\ell d} = [\bar{\ell} \gamma^\mu \ell][\bar{d} \gamma_\mu d]$
C_{qe}	$O_{qe} = [\bar{q} \gamma^\mu q][\bar{e} \gamma_\mu e]$

NC DIS and SMEFT

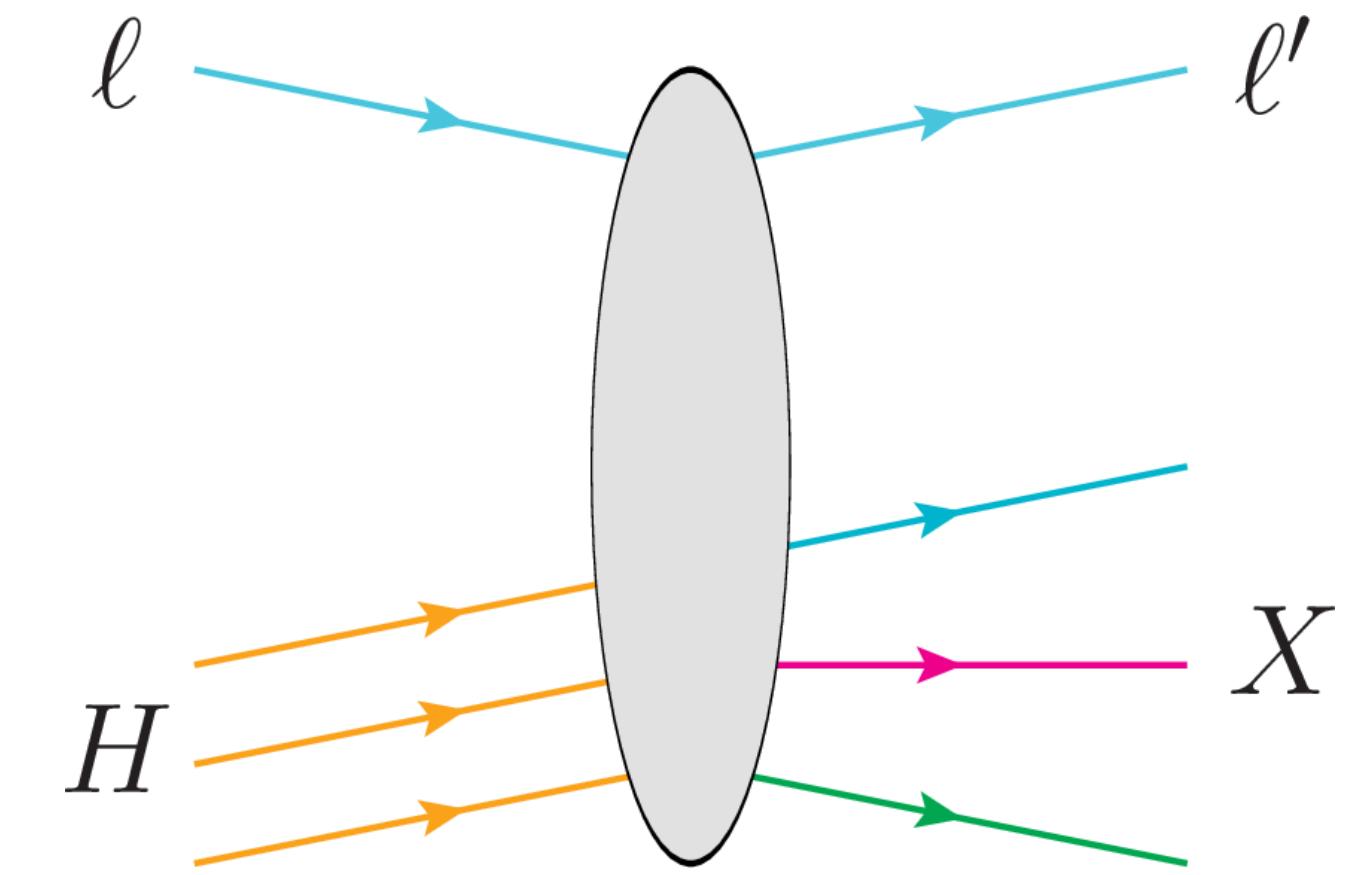
NC DIS in the process $\ell + H \rightarrow \ell' + X$:



$$\mathcal{M} = \mathcal{M}_\gamma + \mathcal{M}_Z + \mathcal{M}_X$$

$$|\mathcal{M}|^2 = \mathcal{M}_{\gamma\gamma} + \mathcal{M}_{\gamma Z} + \mathcal{M}_{ZZ} + \mathcal{M}_{\gamma X} + \mathcal{M}_{ZX} + \mathcal{O}(C^2)$$

$$d\sigma^{\lambda_\ell \lambda_H} = \frac{d^2\sigma(\lambda_\ell, \lambda_H)}{dx \ dQ^2}$$



C_r	O_r
$C_{\ell q}^{(1)}$	$O_{\ell q}^{(1)} = [\bar{\ell} \gamma^\mu \ell][\bar{q} \gamma_\mu q]$
$C_{\ell q}^{(3)}$	$O_{\ell q}^{(3)} = [\bar{\ell} \gamma^\mu \tau^I \ell][\bar{q} \gamma_\mu \tau^I q]$
C_{eu}	$O_{eu} = [\bar{e} \gamma^\mu e][\bar{u} \gamma_\mu u]$
C_{ed}	$O_{ed} = [\bar{e} \gamma^\mu e][\bar{d} \gamma_\mu d]$
$C_{\ell u}$	$O_{\ell u} = [\bar{\ell} \gamma^\mu \ell][\bar{u} \gamma_\mu u]$
$C_{\ell d}$	$O_{\ell d} = [\bar{\ell} \gamma^\mu \ell][\bar{d} \gamma_\mu d]$
C_{qe}	$O_{qe} = [\bar{q} \gamma^\mu q][\bar{e} \gamma_\mu e]$

Polarized and unpolarized cross sections

$$d\sigma^{\lambda_\ell \lambda_H} = \frac{d^2\sigma(\lambda_\ell, \lambda_H)}{dx \ dQ^2}$$

$$d\sigma_0 = \frac{1}{4} [d\sigma^{++} + d\sigma^{+-} + d\sigma^{-+} + d\sigma^{--}]$$

unpolarized lepton + unpolarized hadron

$$d\sigma_\ell = \frac{1}{4} [d\sigma^{++} + d\sigma^{+-} - d\sigma^{-+} - d\sigma^{--}]$$

polarized lepton + unpolarized hadron

$$d\sigma_H = \frac{1}{4} [d\sigma^{++} - d\sigma^{+-} + d\sigma^{-+} - d\sigma^{--}]$$

unpolarized lepton + polarized hadron

$$d\sigma_{\ell H} = \frac{1}{4} [d\sigma^{++} - d\sigma^{+-} - d\sigma^{-+} + d\sigma^{--}]$$

polarized lepton + polarized hadron

Polarized and unpolarized cross sections

$$d\sigma^{\lambda_\ell \lambda_H} = \frac{d^2\sigma(\lambda_\ell, \lambda_H)}{dx \ dQ^2}$$

$$d\sigma_0 = \frac{1}{4} [d\sigma^{++} + d\sigma^{+-} + d\sigma^{-+} + d\sigma^{--}]$$

$$d\sigma_\ell = \frac{1}{4} [d\sigma^{++} + d\sigma^{+-} - d\sigma^{-+} - d\sigma^{--}]$$

$$d\sigma_H = \frac{1}{4} [d\sigma^{++} - d\sigma^{+-} + d\sigma^{-+} - d\sigma^{--}]$$

$$d\sigma_{\ell H} = \frac{1}{4} [d\sigma^{++} - d\sigma^{+-} - d\sigma^{-+} + d\sigma^{--}]$$

$$A_{\text{PV}}^{(\ell)} := \frac{d\sigma_\ell}{d\sigma_0} \quad \text{unpolarized } A_{\text{PV}}$$

$$A_{\text{PV}}^{(H)} := \frac{d\sigma_H}{d\sigma_0} \quad \text{polarized } A_{\text{PV}}$$

$$A_{\text{LC}}^{(H)} := \frac{d\sigma_0(e^+H) - d\sigma_0(e^-H)}{d\sigma_0(e^+H) + d\sigma_0(e^-H)}$$

lepton-charge A

Measurement of PV asymmetries at the EIC

$$d\sigma = d\sigma_0 + P_\ell d\sigma_e + P_H d\sigma_H + P_\ell P_H d\sigma_{\ell H}$$

$$\begin{array}{c} P_\ell, P_H \leftrightarrow \lambda_\ell, \lambda_H \\ \downarrow \\ -1 \leq \quad \leq 1 \end{array}$$

Measurement of PV asymmetries at the EIC

$$d\sigma = d\sigma_0 + P_\ell d\sigma_e + P_H d\sigma_H + P_\ell P_H d\sigma_{\ell H}$$

integrated
luminosity

$$N^{++} = a_{\text{det}} L^{++} \left[d\sigma_0 + |P_\ell^{++}| d\sigma_\ell + |P_H^{++}| d\sigma_H + |P_\ell^{++}| |P_H^{++}| d\sigma_{\ell H} \right]$$

$$N^{+-} = a_{\text{det}} L^{+-} \left[d\sigma_0 + |P_\ell^{+-}| d\sigma_\ell - |P_H^{+-}| d\sigma_H - |P_\ell^{+-}| |P_H^{+-}| d\sigma_{\ell H} \right]$$

$$N^{-+} = a_{\text{det}} L^{-+} \left[d\sigma_0 - |P_\ell^{-+}| d\sigma_\ell + |P_H^{-+}| d\sigma_H - |P_\ell^{-+}| |P_H^{-+}| d\sigma_{\ell H} \right]$$

$N^{--} = a_{\text{det}} L^{--} \left[d\sigma_0 - |P_\ell^{--}| d\sigma_\ell - |P_H^{--}| d\sigma_H + |P_\ell^{--}| |P_H^{--}| d\sigma_{\ell H} \right]$

event count

detector phase space
acceptance and efficiency

$$P_\ell, P_H \leftrightarrow \lambda_\ell, \lambda_H$$

$$-1 \leq \lambda \leq 1$$

Measurement of PV asymmetries at the EIC

$$d\sigma = d\sigma_0 + P_\ell d\sigma_e + P_H d\sigma_H + P_\ell P_H d\sigma_{\ell H}$$

$$P_\ell, P_H \leftrightarrow \lambda_\ell, \lambda_H$$

$$-1 \leq \quad \leq 1$$

Assuming P_ℓ^{ij} , P_H^{ij} , L^{ij} , and a_{det} are constant:

$$\xrightarrow{\hspace{1cm}} d\sigma^{ij} = N^{ij}/L^{ij}/a_{\text{det}}$$

$$d\sigma_0 = \frac{1}{4} [d\sigma^{++} + d\sigma^{+-} + d\sigma^{-+} + d\sigma^{--}]$$

$$d\sigma_\ell = \frac{1}{4|P_\ell|} [d\sigma^{++} + d\sigma^{+-} - d\sigma^{-+} - d\sigma^{--}]$$

$$d\sigma_H = \frac{1}{4|P_H|} [d\sigma^{++} - d\sigma^{+-} + d\sigma^{-+} - d\sigma^{--}]$$

$$d\sigma_{\ell H} = \frac{1}{4|P_\ell||P_H|} [d\sigma^{++} - d\sigma^{+-} - d\sigma^{-+} + d\sigma^{--}]$$

Measurement of PV asymmetries at the EIC

$$d\sigma = d\sigma_0 + P_\ell d\sigma_e + P_H d\sigma_H + P_\ell P_H d\sigma_{\ell H}$$

$$P_\ell, P_H \leftrightarrow \lambda_\ell, \lambda_H$$

$$-1 \leq \quad \leq 1$$

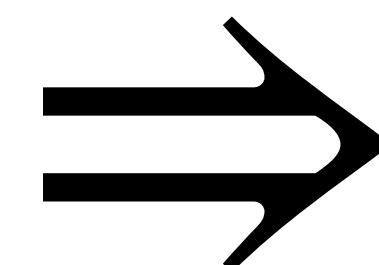
Assuming P_ℓ^{ij} , P_H^{ij} , L^{ij} , and a_{det} are constant:

$$d\sigma_0 = \frac{1}{4} [d\sigma^{++} + d\sigma^{+-} + d\sigma^{-+} + d\sigma^{--}]$$

$$d\sigma_\ell = \frac{1}{4|P_\ell|} [d\sigma^{++} + d\sigma^{+-} - d\sigma^{-+} - d\sigma^{--}]$$

$$d\sigma_H = \frac{1}{4|P_H|} [d\sigma^{++} - d\sigma^{+-} + d\sigma^{-+} - d\sigma^{--}]$$

$$d\sigma_{\ell H} = \frac{1}{4|P_\ell||P_H|} [d\sigma^{++} - d\sigma^{+-} - d\sigma^{-+} + d\sigma^{--}]$$



$\curvearrowright Y^{ij} = N^{ij}/L^{ij}$

$$A_{\text{PV}}^{(\ell)} = \frac{d\sigma_\ell}{d\sigma_0} = \frac{1}{|P_\ell|} \frac{Y^{++} + Y^{+-} - Y^{-+} - Y^{--}}{Y^{++} + Y^{+-} + Y^{-+} + Y^{--}}$$

$$A_{\text{PV}}^{(H)} = \frac{d\sigma_H}{d\sigma_0} = \frac{1}{|P_H|} \frac{Y^{++} - Y^{+-} + Y^{-+} - Y^{--}}{Y^{++} + Y^{+-} + Y^{-+} + Y^{--}}$$

Measurement of PV asymmetries at the EIC

$$d\sigma = d\sigma_0 + P_\ell d\sigma_e + P_H d\sigma_H + P_\ell P_H d\sigma_{\ell H}$$

$$P_\ell, P_H \leftrightarrow \lambda_\ell, \lambda_H$$

$$-1 \leq \quad \leq 1$$

Assuming P_ℓ^{ij} , P_H^{ij} , L^{ij} , and a_{det} are constant:

$$A_{\text{PV}}^{(\ell)} = \frac{d\sigma_\ell}{d\sigma_0} = \frac{1}{|P_\ell|} \frac{Y^{++} + Y^{+-} - Y^{-+} - Y^{--}}{Y^{++} + Y^{+-} + Y^{-+} + Y^{--}}$$

point-to-point luminosity
uncertainty $\sim 10^{-4}$

$$A_{\text{PV}}^{(H)} = \frac{d\sigma_H}{d\sigma_0} = \frac{1}{|P_H|} \frac{Y^{++} - Y^{+-} + Y^{-+} - Y^{--}}{Y^{++} + Y^{+-} + Y^{-+} + Y^{--}}$$

dominant uncertainty:
polarimetry

Measurement of LC asymmetries at the EIC

PV asymmetries:

- Compare scattering yields of LH and RH beams
- Short-time scale
- Dominant uncertainty: **polarimetry**

Measurement of LC asymmetries at the EIC

PV asymmetries:

- Compare scattering yields of LH and RH beams
- Short-time scale
- Dominant uncertainty: **polarimetry**

LC asymmetries:

- Compare between e^- and e^+ runs
- Two independent cross-section measurements
- Reverse detector magnet polarity to minimize systematic errors in e^- and e^+ detection
- Dominant uncertainty: **luminosity difference**

Projections of PV and LC asymmetry data

ECCE detector configuration for inclusive NC study

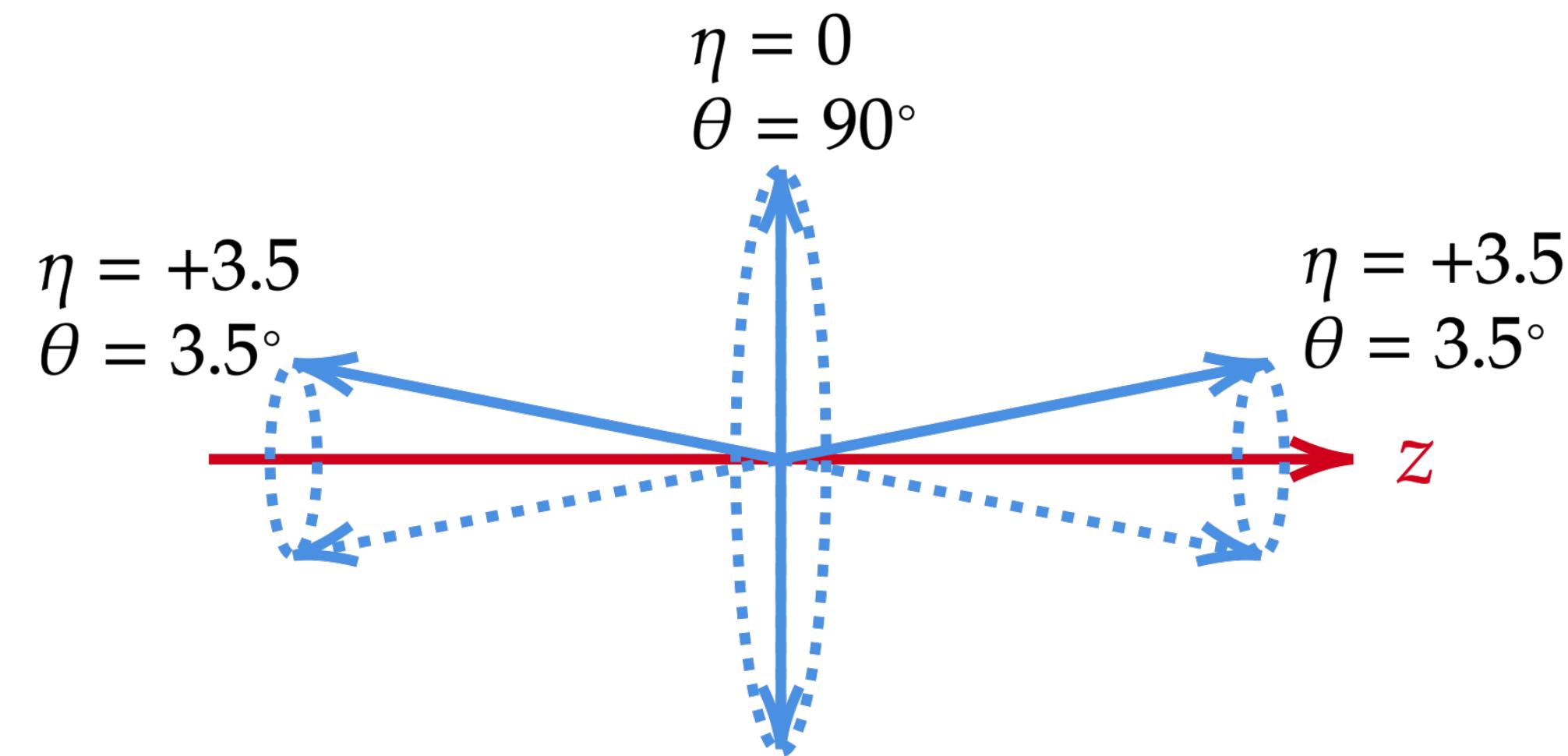
Accardi *et al.* 1212.1701

- Hybrid tracking detector and EM calorimetry: $|\eta| \leq 3.5$ with full azimuthal coverage

- Determination of inclusive DIS kinematics:

- Single- e^- simulations in full detector: p_e, θ_e, φ_e
- Detected vs. true values: smearing
- Apply smearing to simulated events without involving full detector

⇒ “fast-smearing”



Simulation with fast-smearing

- Djangoh MC event generator: full EW radiative effects
- Apply fast-smearing to inclusive e^- events
- Get number of e^- using σ and L

Simulation with fast-smearing

- Djangoh MC event generator: full EW radiative effects
- Apply fast-smearing to inclusive e^- events
- Get number of e^- using σ and L
- Bin migration: [PHENIX Collaboration 1402.1209](#)
 - Due to radiative effects
 - Unfold

Simulation with fast-smearing

- Bin migration: [PHENIX Collaboration 1402.1209](#)
 - Due to radiative effects
 - Unfold
- Background reactions:
 - Due to hadronic final state
 - High at large values of y

Event selection

$Q > 1 \text{ GeV}$	to avoid nonperturbative region of QCD
$y > 0.1$	to avoid bin migration and unfolding uncertainty
$y < 0.9$	to avoid high photoproduction background
$-3.5 < \eta < 3.5$	to restrict events in main acceptance of ECCE detector
$E' > 2 \text{ GeV}$	to ensure high purity of e^- samples

Event selection

$Q > 1 \text{ GeV}$	to avoid nonperturbative region of QCD
$y > 0.1$	to avoid bin migration and unfolding uncertainty
$y < 0.9$	to avoid high photoproduction background
$-3.5 < \eta < 3.5$	to restrict events in main acceptance of ECCE detector
$E' > 2 \text{ GeV}$	to ensure high purity of e^- samples
$x < 0.5$	additional cuts for SMEFT analysis:
$Q > 10 \text{ GeV}$	to remove <i>large</i> uncertainties from non-perturbative QCD and nuclear dynamics

Data sets

Beam energy, beam type, and nominal annual luminosity (NL) assumed for the EIC analysis:

<i>eD</i> scattering		<i>ep</i> scattering	
D1	$5 \text{ GeV} \times 41 \text{ GeV } eD, 4.4 \text{ fb}^{-1}$	P1	$5 \text{ GeV} \times 41 \text{ GeV } ep, 4.4 \text{ fb}^{-1}$
D2	$5 \text{ GeV} \times 100 \text{ GeV } eD, 36.8 \text{ fb}^{-1}$	P2	$5 \text{ GeV} \times 100 \text{ GeV } ep, 36.8 \text{ fb}^{-1}$
D3	$10 \text{ GeV} \times 100 \text{ GeV } eD, 44.8 \text{ fb}^{-1}$	P3	$10 \text{ GeV} \times 100 \text{ GeV } ep, 44.8 \text{ fb}^{-1}$
D4	$10 \text{ GeV} \times 137 \text{ GeV } eD, 100 \text{ fb}^{-1}$	P4	$10 \text{ GeV} \times 275 \text{ GeV } ep, 100 \text{ fb}^{-1}$
D5	$18 \text{ GeV} \times 137 \text{ GeV } eD, 15.4 \text{ fb}^{-1}$	P5	$18 \text{ GeV} \times 275 \text{ GeV } ep, 15.4 \text{ fb}^{-1}$
		P6	$18 \text{ GeV} \times 275 \text{ GeV } ep, 100 \text{ fb}^{-1}$

CM energy

energy of electron beam

energy of hadron beam

YR reference setting [2103.05419]

Statistical uncertainty projection for PV asymmetries

For a given value of integrated luminosity:

$$\delta A_{\text{stat}}^{\text{measured}} = \frac{1}{\sqrt{N}}$$

Statistical uncertainty projection for PV asymmetries

For a given value of integrated luminosity:

$$\delta A_{\text{stat}}^{\text{measured}} = \frac{1}{\sqrt{N}}$$

For PV asymmetries, corrections for beam polarization:

$$[\delta A_{\text{PV}}^{(\ell)}]_{\text{stat}} = \frac{1}{|P_\ell|} \frac{1}{\sqrt{N}}, \quad [\delta A_{\text{PV}}^{(H)}]_{\text{stat}} = \frac{1}{|P_H|} \frac{1}{\sqrt{N}}$$

Assume:

- $P_\ell = 80\%$ with 1% rel. systematic uncertainty
- $P_H = 70\%$ with 2% rel. systematic uncertainty

Statistical uncertainty projection for PV asymmetries

For a given value of integrated luminosity:

$$\delta A_{\text{stat}}^{\text{measured}} = \frac{1}{\sqrt{N}}$$

For PV asymmetries, corrections for beam polarization:

$$[\delta A_{\text{PV}}^{(\ell)}]_{\text{stat}} = \frac{1}{|P_\ell|} \frac{1}{\sqrt{N}}, \quad [\delta A_{\text{PV}}^{(H)}]_{\text{stat}} = \frac{1}{|P_H|} \frac{1}{\sqrt{N}}$$

Assume:

- $P_\ell = 80\%$ with 1% rel. systematic uncertainty
- $P_H = 70\%$ with 2% rel. systematic uncertainty

large because
 $A_{\text{PV}}^{(H)} \ll A_{\text{PV}}^{(\ell)}$

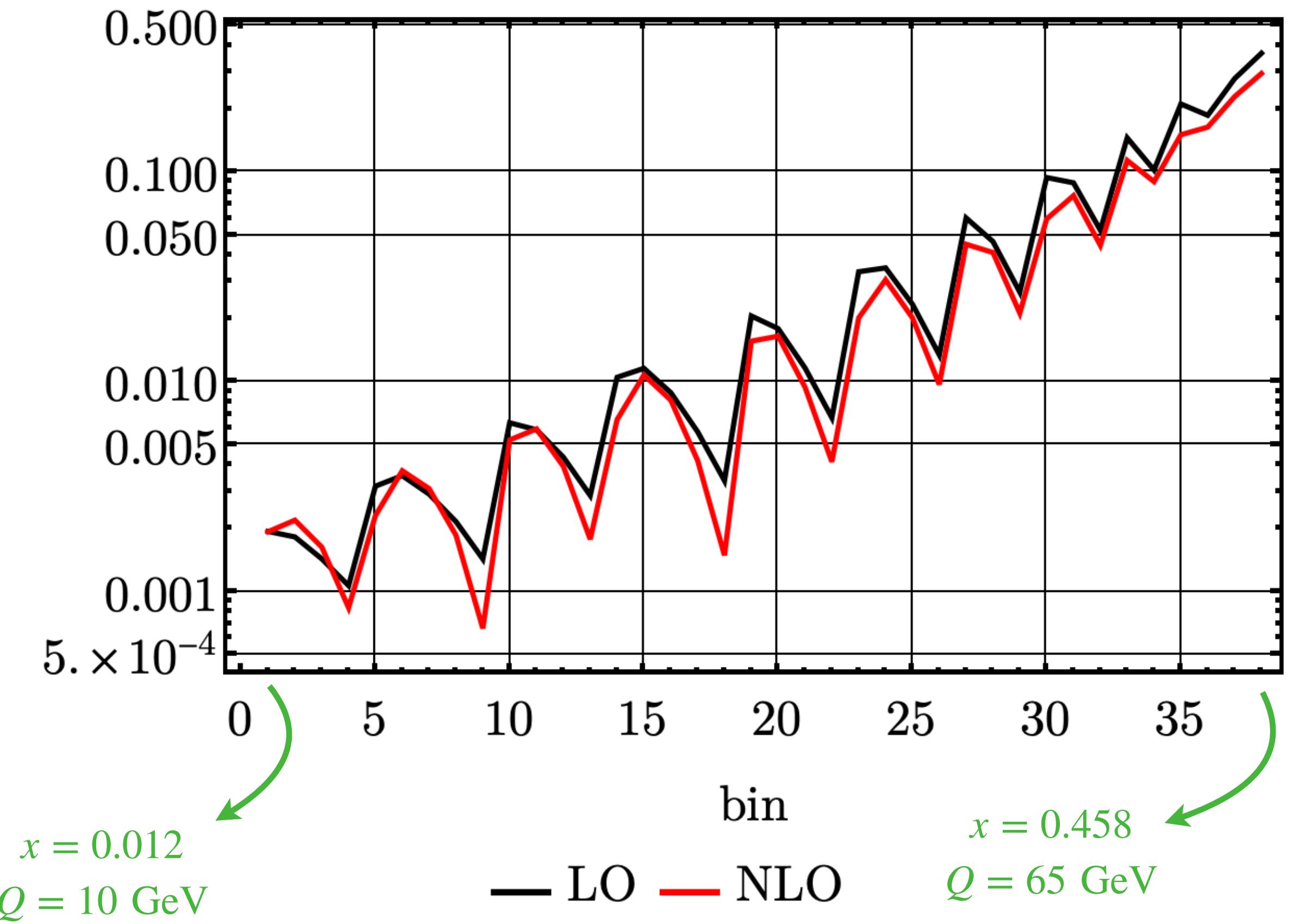
Statistical uncertainty projection for LC asymmetries

- Dominant uncertainty: luminosity difference between e^- and e^+ runs
- Assume 2 % : relative in luminosity, absolute in $A_{\text{LC}}^{(H)}$
- Reverse detector magnet polarity:
 - ⇒ detection of DIS $e^+ \sim$ detection of DIS e^-
 - ⇒ $\left[\delta A_{\text{LC}}^{(H)} \right]_{\text{stat}}$ determined by luminosity of e^+ run
- Assume $L^{e^+} = \frac{1}{10} L^{e^-}$.
- **Polarimetry** is irrelevant for A_{LC} measurements.

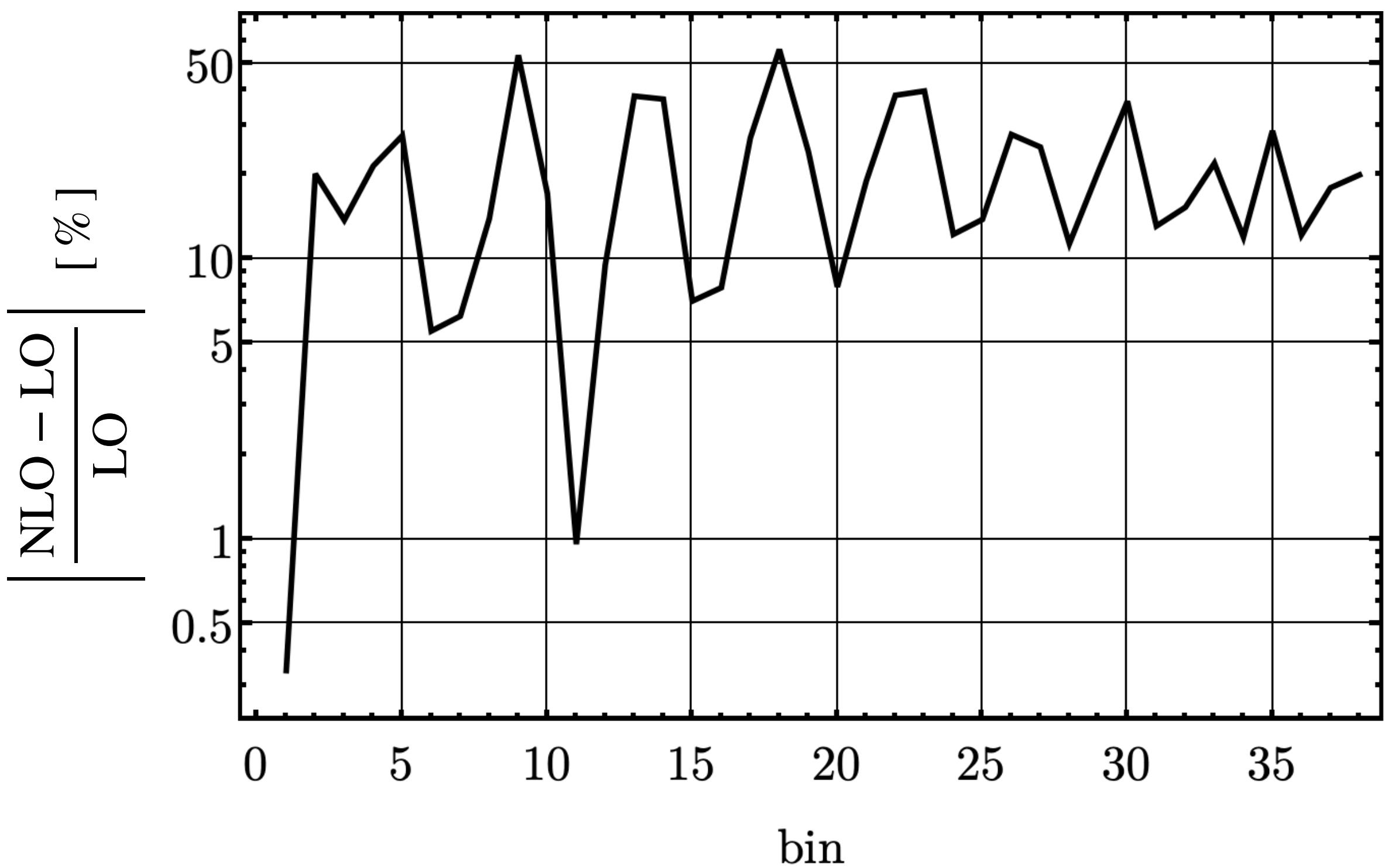
QED uncertainty projection for LC asymmetries

- Higher-order QED effects in e^- and e^+ DIS cross sections
- LO and NLO $A_{\text{LC}}^{(H)}$ computed using Djangoh

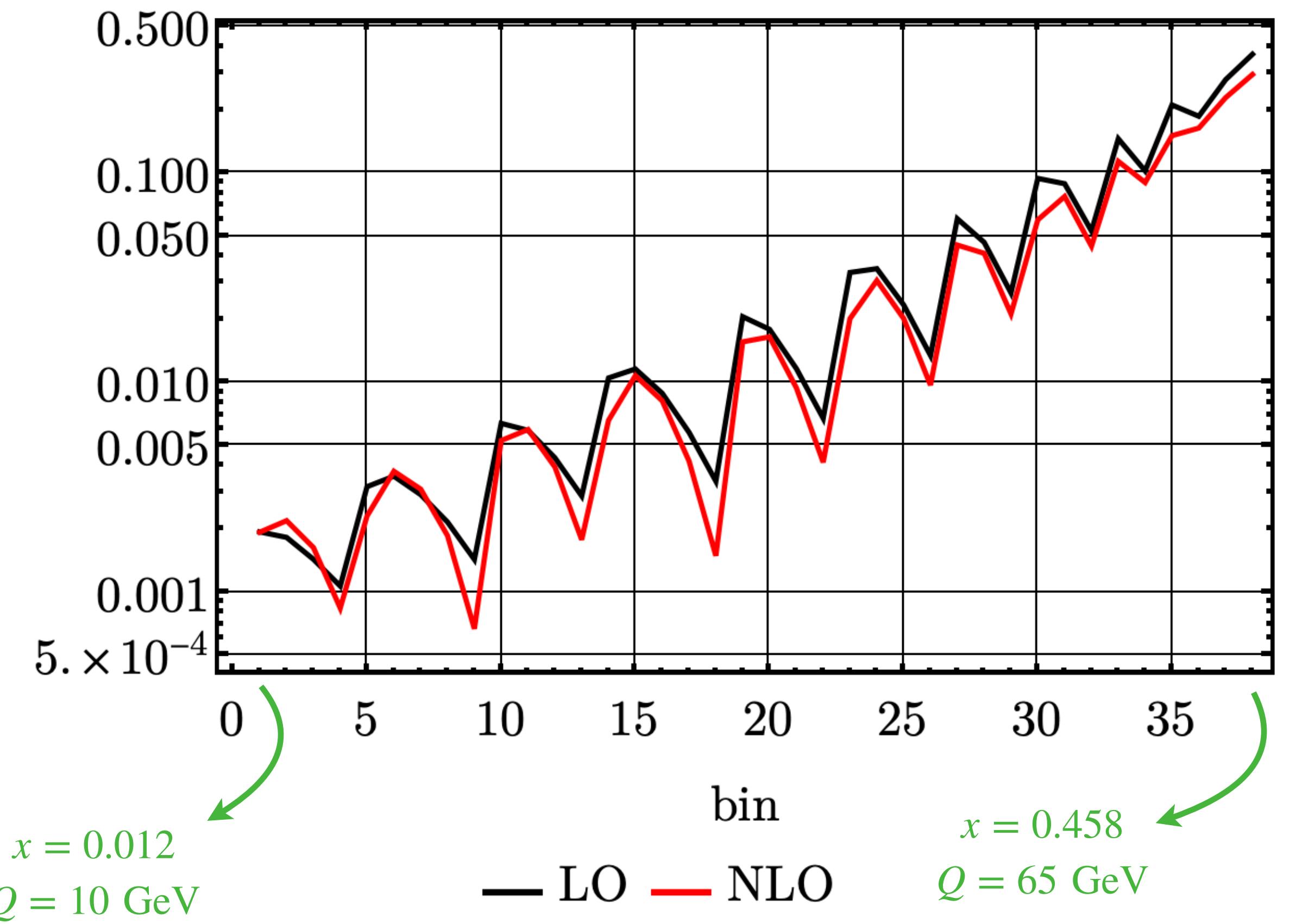
ep 10 GeV \times 275 GeV 100 fb^{-1}
lepton-charge A of proton



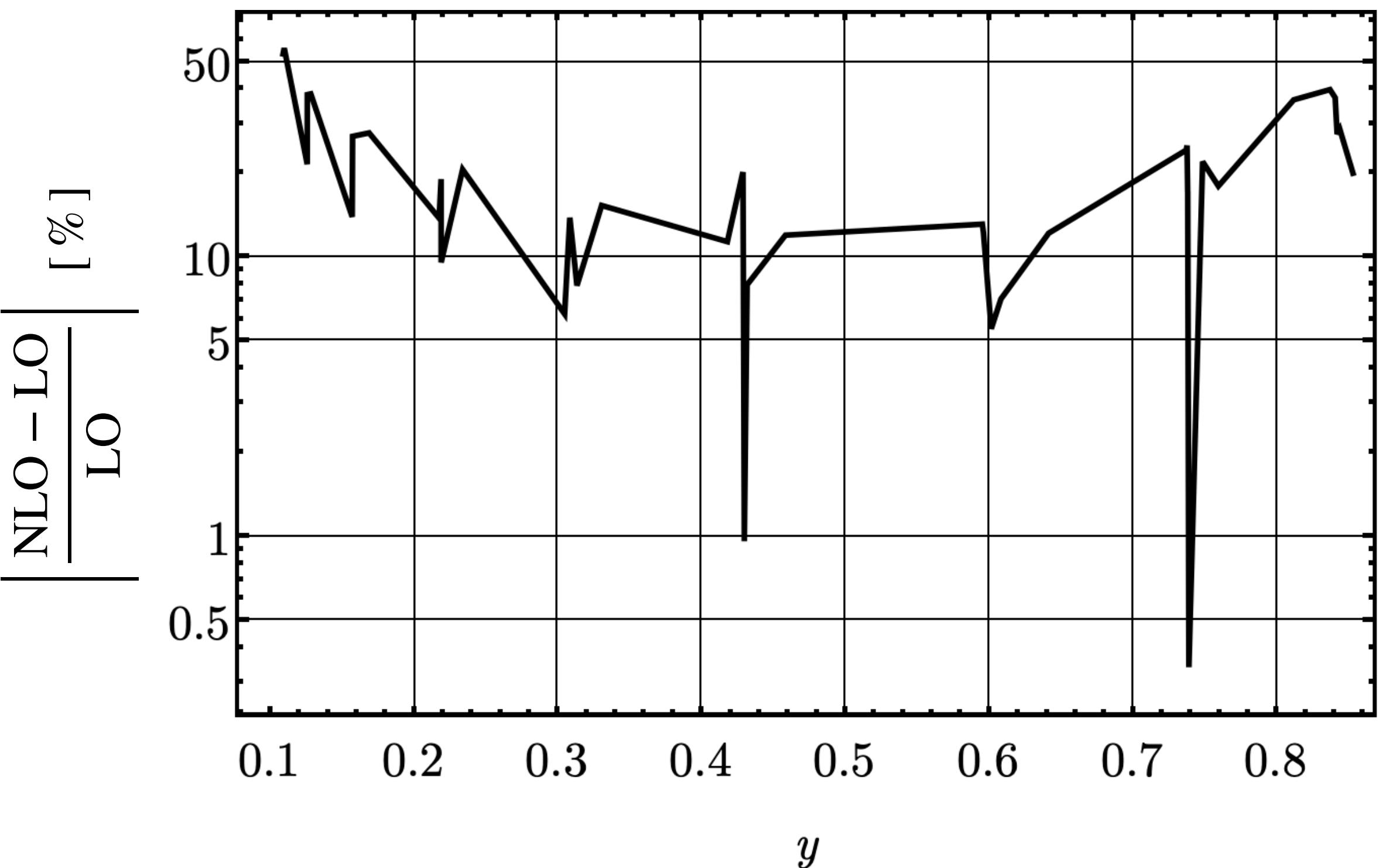
ep 10 GeV \times 275 GeV 100 fb^{-1}
lepton-charge A of proton



ep 10 GeV \times 275 GeV 100 fb $^{-1}$
lepton-charge A of proton



ep 10 GeV \times 275 GeV 100 fb $^{-1}$
lepton-charge A of proton



QED uncertainty projection for LC asymmetries

- Higher-order QED effects in e^- and e^+ DIS cross sections
- LO and NLO $A_{\text{LC}}^{(H)}$ computed using Djangoh
- Introduce 5 % of $[A_{\text{LC}}^{(H)}]_{\text{NLO}} - [A_{\text{LC}}^{(H)}]_{\text{LO}}$ as QED NLO uncertainty

Projection for High-Luminosity EIC

- Possibility of additional $10 \times$ increase in annual luminosity beyond initial run
- Assuming all experimental systematic effects remain the same:

$$\sigma_{\text{stat}} \rightarrow \frac{1}{\sqrt{10}} \sigma_{\text{stat}}$$

Uncertainties

Anticipated uncertainties

PV asymmetries:

- σ_{stat}
- $\frac{\sigma_{\text{sys}}^{\text{unc}}}{A} = 1\% \text{ rel.}$
- $\frac{\sigma_{\text{pol}}^{\text{cor}}}{A} = \begin{cases} 1\% \text{ rel. for polarized } \ell \\ 2\% \text{ rel. for polarized } H \end{cases}$
- $\sigma_{\text{pdf}}^{\text{cor}}$

Anticipated uncertainties

PV asymmetries:

- σ_{stat}
- $\frac{\sigma_{\text{sys}}^{\text{unc}}}{A} = 1\% \text{ rel.}$
- $\frac{\sigma_{\text{pol}}^{\text{cor}}}{A} = \begin{cases} 1\% \text{ rel. for polarized } \ell \\ 2\% \text{ rel. for polarized } H \end{cases}$
- $\sigma_{\text{pdf}}^{\text{cor}}$

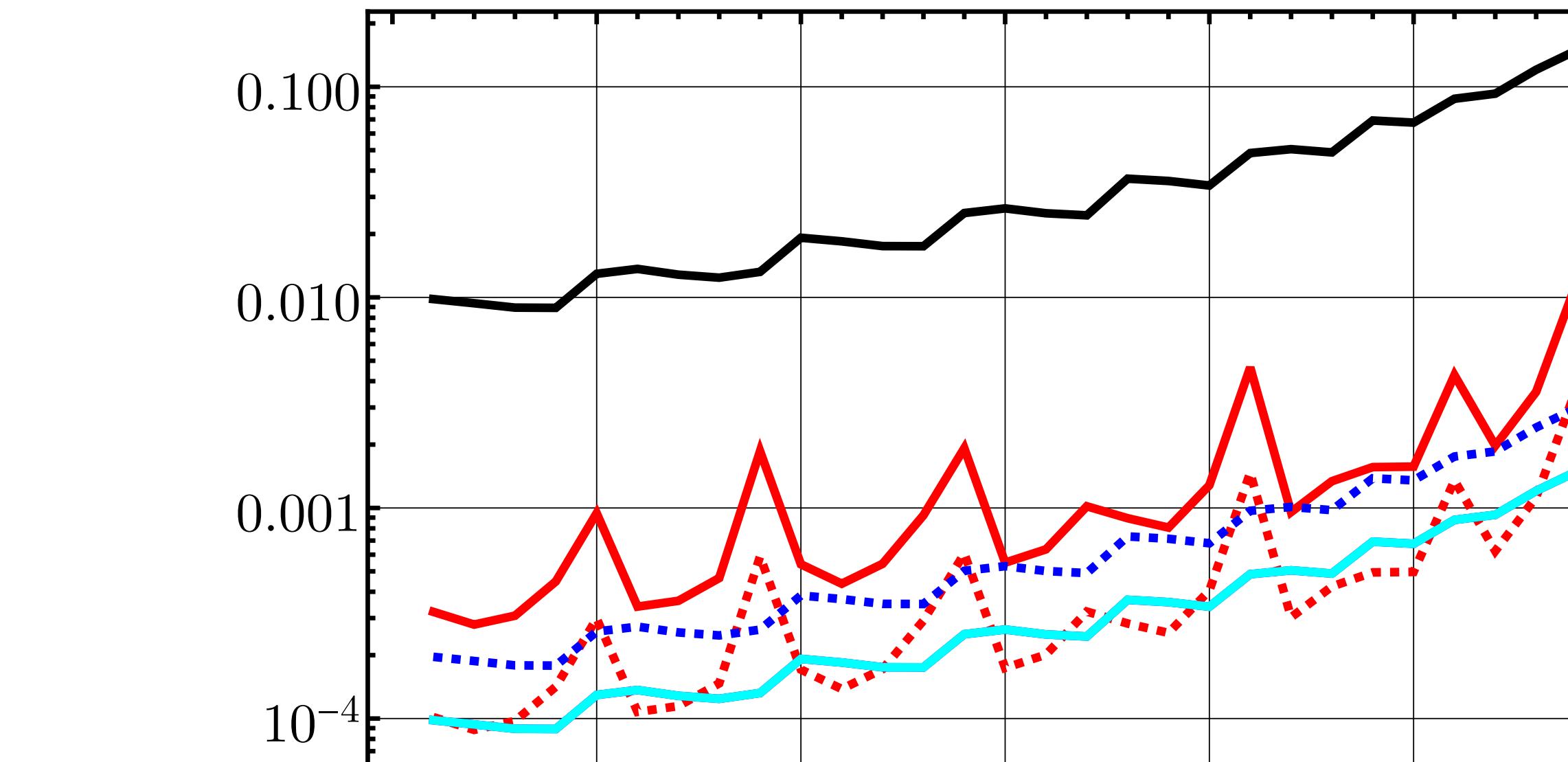
LC asymmetries:

- σ_{stat}
- $\frac{\sigma_{\text{sys}}^{\text{unc}}}{A} = 1\% \text{ rel.}$
- $\sigma_{\text{lum}}^{\text{cor}} = 2\% \text{ abs.}$
- $\sigma_{\text{nlo}}^{\text{unc}} = 5\% \cdot (A_{\text{LC}}^{\text{NLO}} - A_{\text{LC}}^{\text{Born}})$
- $\sigma_{\text{pdf}}^{\text{cor}}$

eD 10 GeV \times 137 GeV 100 fb $^{-1}$

unpolarized A_{PV} of deuteron

$q_a = u, d$



$x = 0.026$
 $Q = 10 \text{ GeV}$

$A_{SM,0}^{\text{theo}}$

$\sigma_{\text{stat}} (\text{NL})$

$\sigma_{\text{stat}} (\text{HL})$

$1\% \text{ sys (rel)}$ $2\% \text{ sys (rel)}$

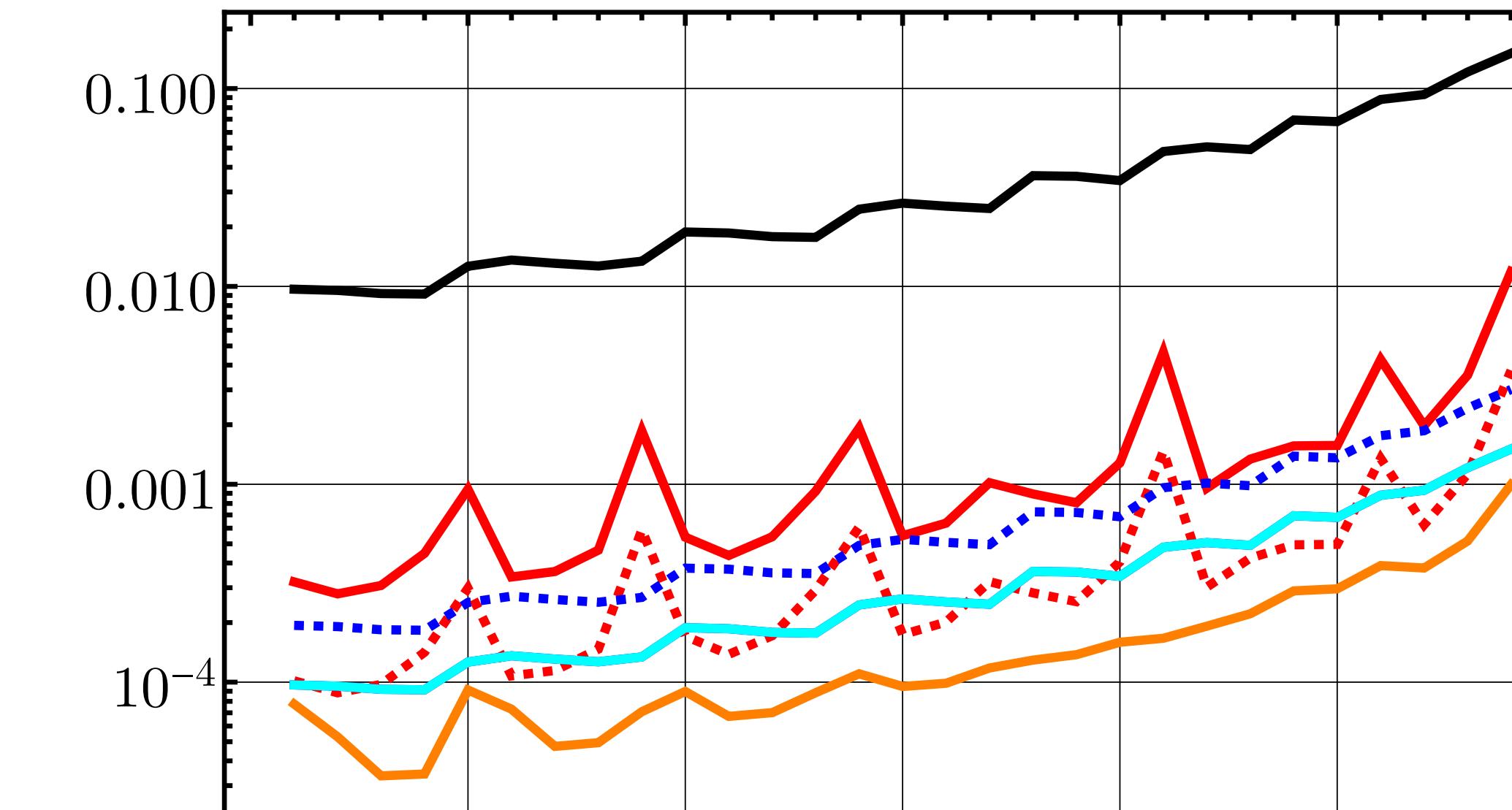
$1\% \text{ pol (rel)}$

σ_{pdf}

eD 10 GeV \times 137 GeV 100 fb $^{-1}$

unpolarized A_{PV} of deuteron

$q_a = u, \bar{u}, d, \bar{d}, s, \bar{s}$



$x = 0.461$
 $Q = 46 \text{ GeV}$

$A_{SM,0}^{\text{theo}}$

$\sigma_{\text{stat}} (\text{NL})$

$\sigma_{\text{stat}} (\text{HL})$

$1\% \text{ sys (rel)}$ $2\% \text{ sys (rel)}$

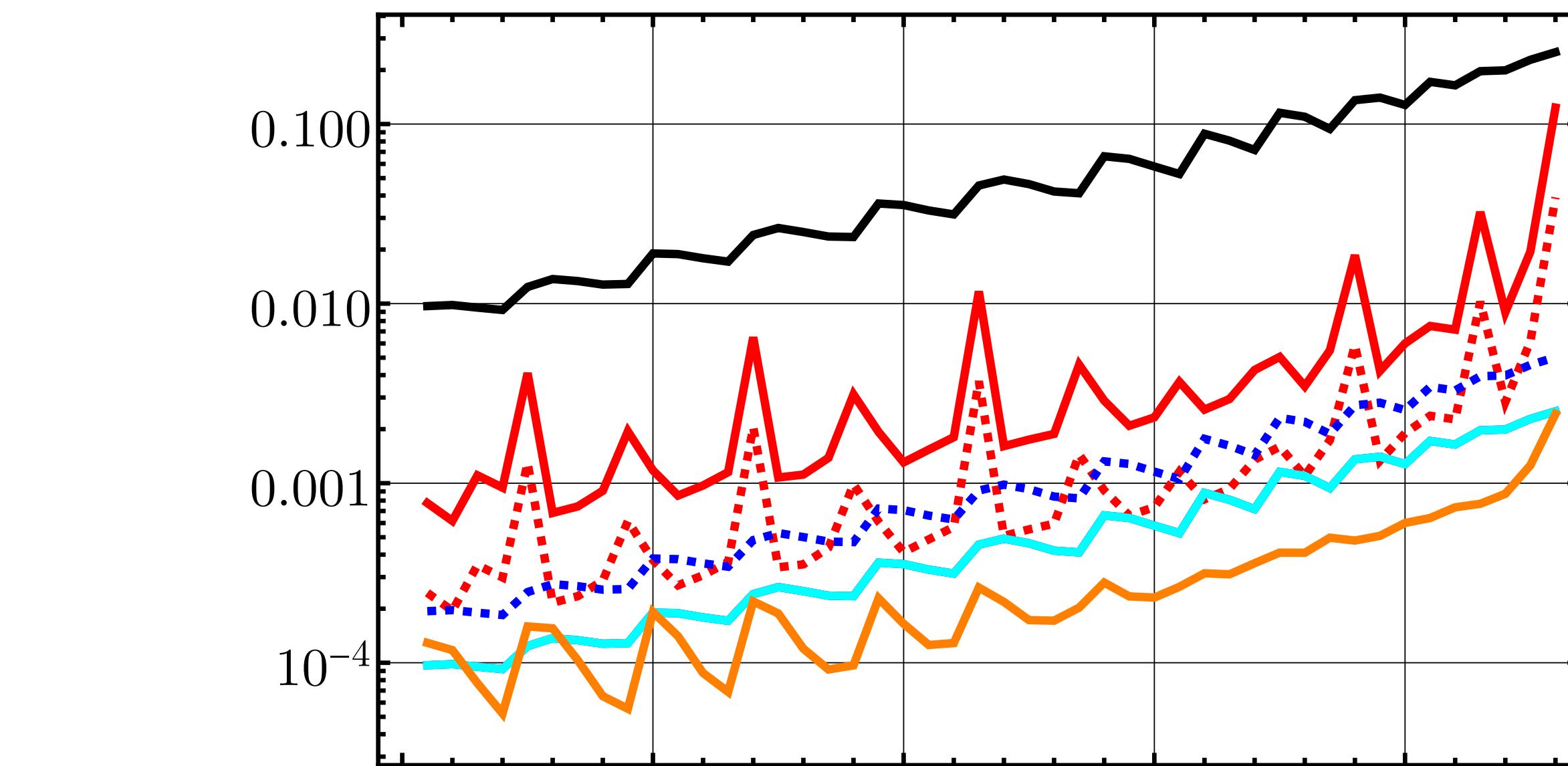
$1\% \text{ pol (rel)}$

σ_{pdf}

ep 18 GeV \times 275 GeV 15.4 fb $^{-1}$

unpolarized A_{PV} of proton

$q_a = u, \bar{u}, d, \bar{d}, s, \bar{s}$



$x = 0.007$
 $Q = 10$ GeV

$x = 0.487$
 $Q = 92$ GeV

— $A_{SM,0}^{\text{theo}}$

— $\sigma_{\text{stat}} (\text{NL})$ ····· $\sigma_{\text{stat}} (\text{HL})$

— 1% sys (rel) ····· 2% sys (rel)

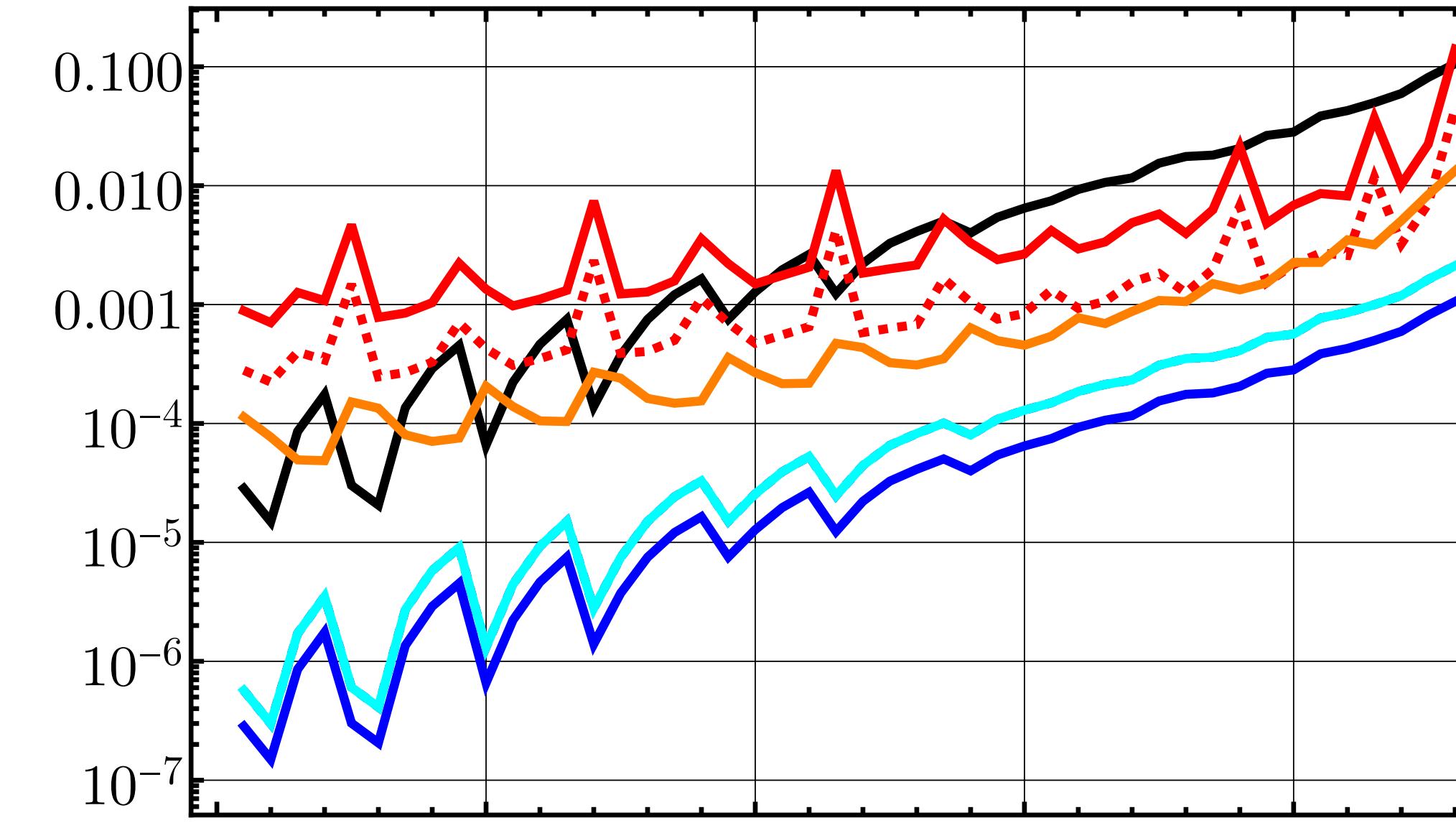
— 1% pol (rel)

— σ_{pdf}

ep 18 GeV \times 275 GeV 15.4 fb $^{-1}$

polarized A_{PV} of proton

$q_a = u, \bar{u}, d, \bar{d}, s, \bar{s}$



— $A_{SM,0}^{\text{theo}}$

— $\sigma_{\text{stat}} (\text{NL})$ ····· $\sigma_{\text{stat}} (\text{HL})$

— 1% sys (rel) ····· 2% sys (rel)

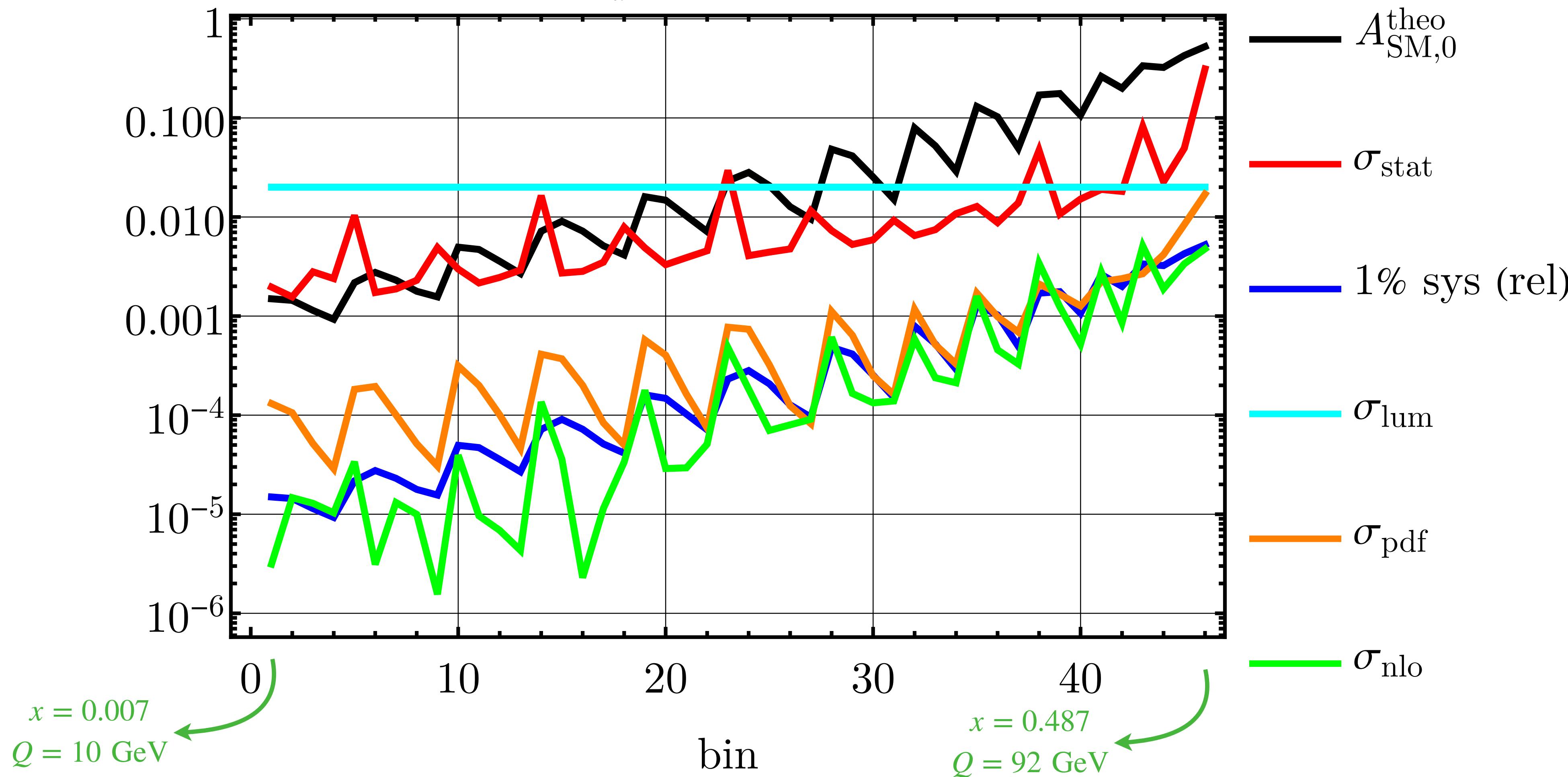
— 2% pol (rel)

— σ_{pdf}

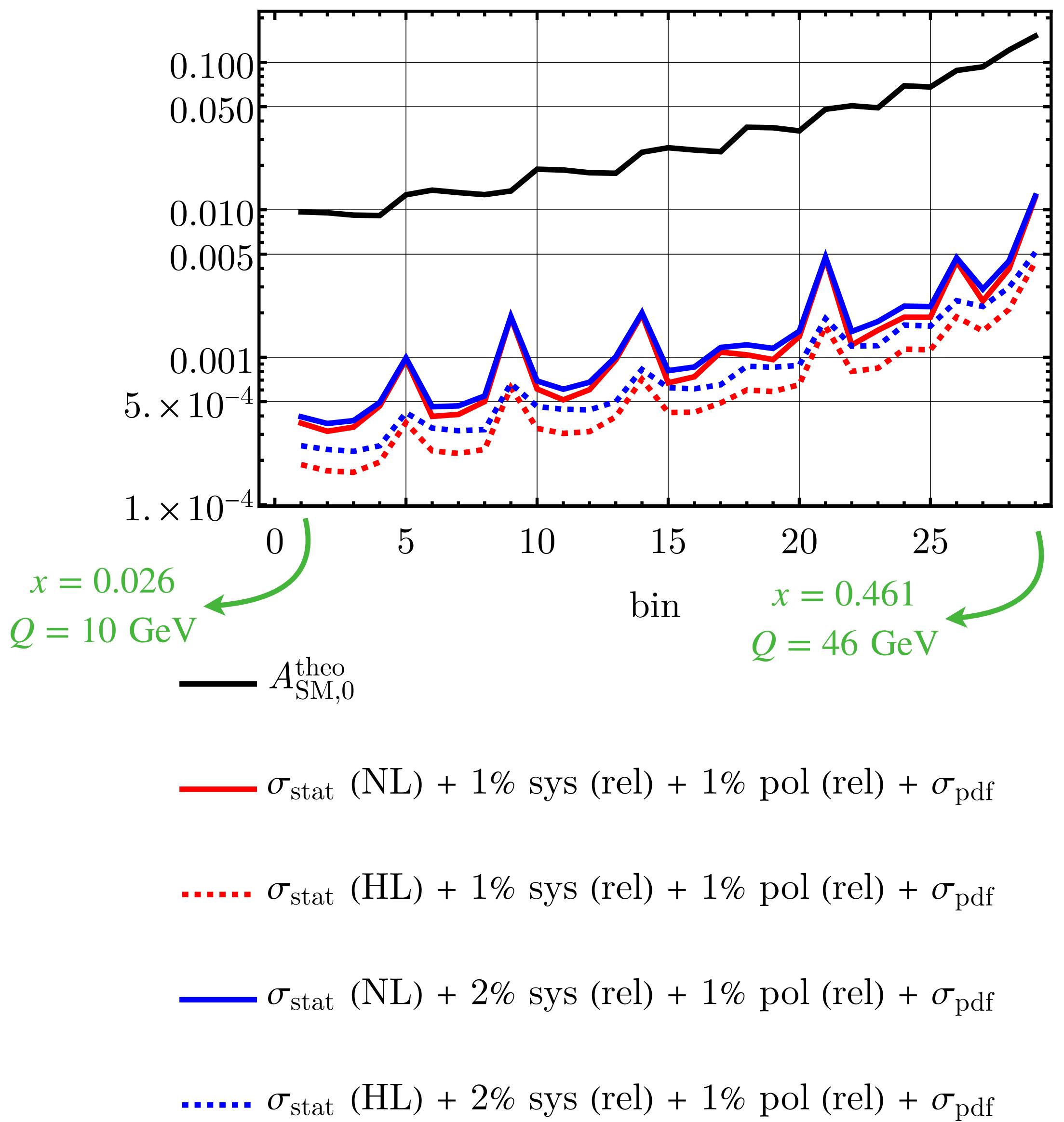
ep 18 GeV \times 275 GeV 15.4 fb^{-1}

lepton-charge A of proton

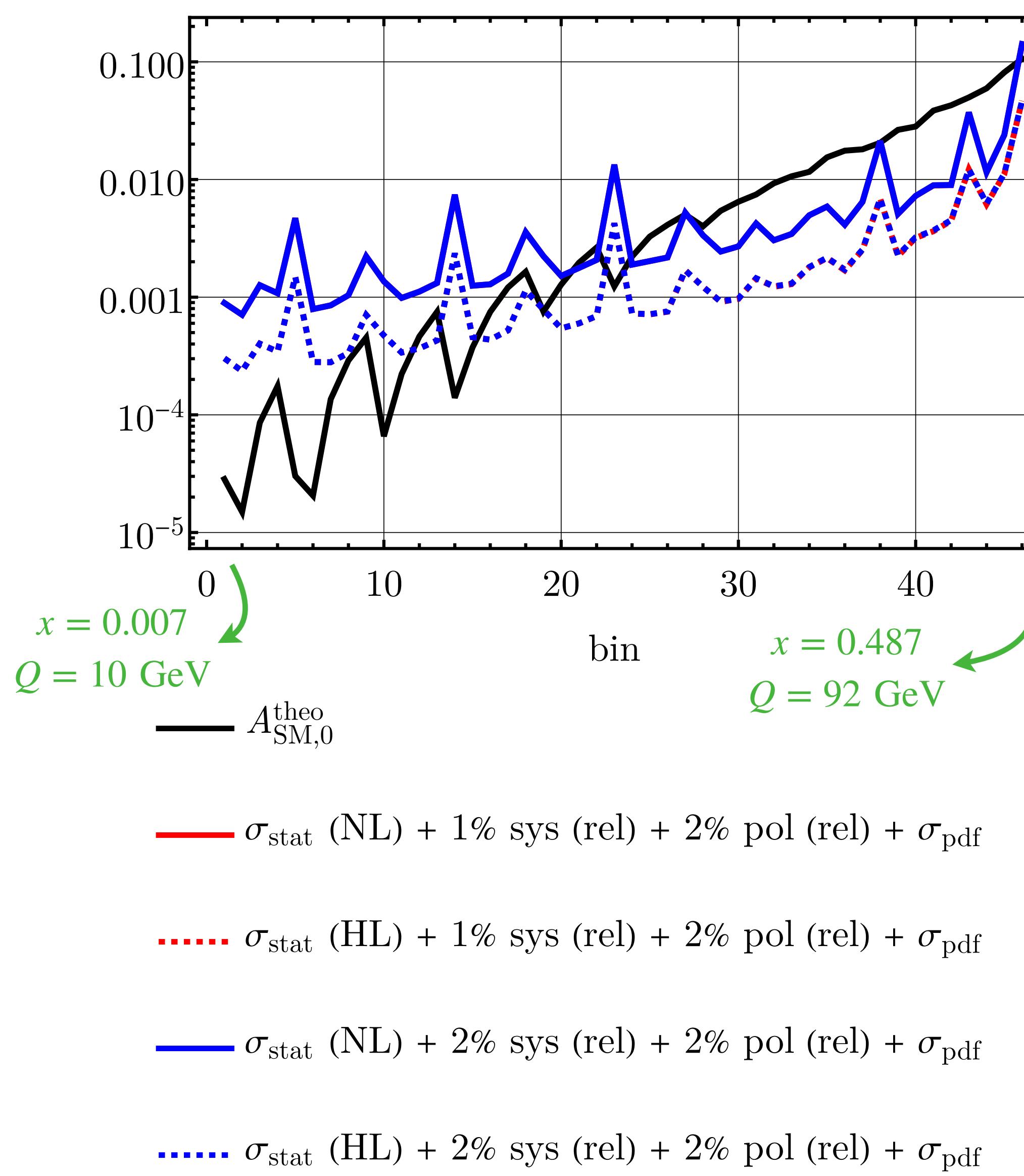
$q_a = u, \bar{u}, d, \bar{d}, s, \bar{s}$



eD 10 GeV \times 137 GeV 100 fb $^{-1}$
unpolarized A_{PV} of deuteron



ep 18 GeV \times 275 GeV 15.4 fb $^{-1}$
polarized A_{PV} of proton



Uncertainty assumptions

Brief summary:

- Quark flavors: up to strange or only at up and down.
- σ_{stat} : dominant for **unpolarized** A_{PV} in NL; comparable to σ_{sys} in HL.
- PDF uncertainties: negligible for **unpolarized** A_{PV} , significant for **polarized** A_{PV} .
- **Luminosity** effects $> \sigma_{\text{stat}}$.
- Higher-order QED corrections to **lepton-charge** A are negligible.

Framework of the SMEFT analysis

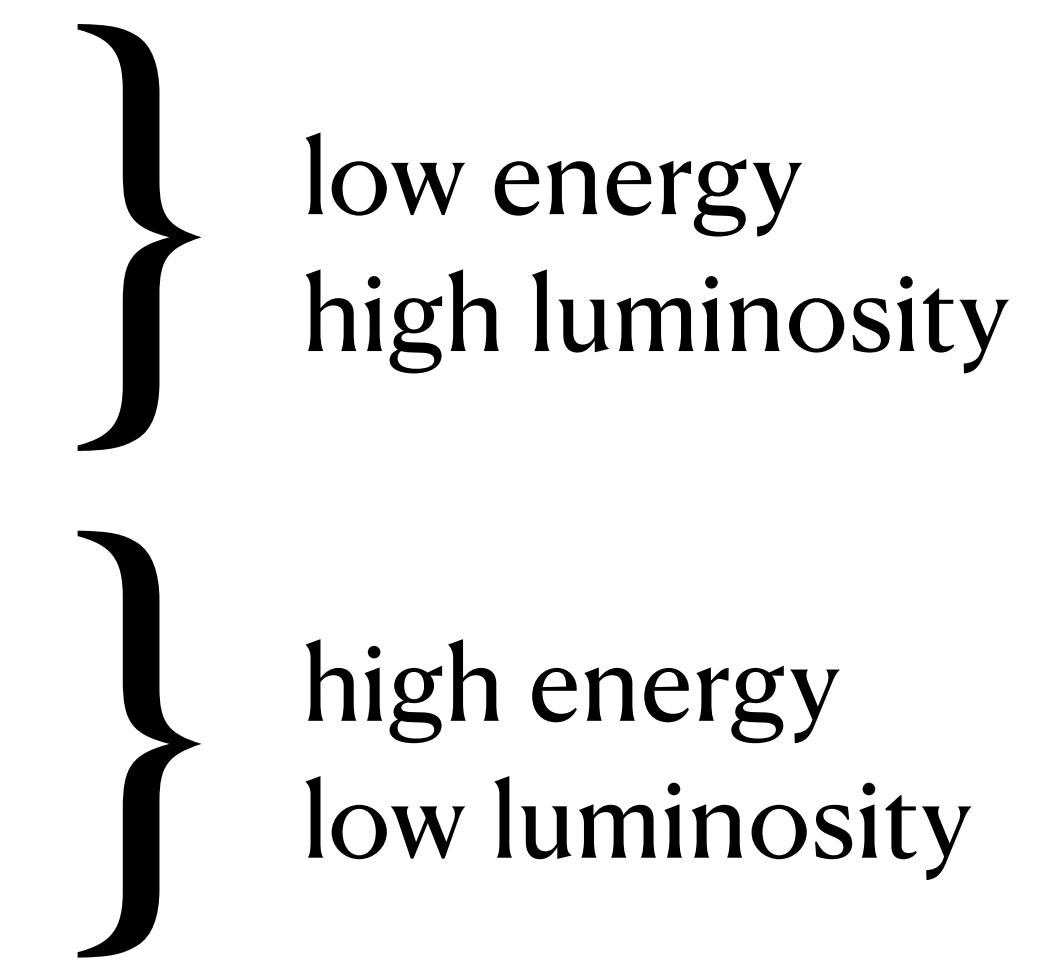
Data sets

Concentrate on the two highest-energy settings:

- $10 \text{ GeV} \times 137 \text{ GeV } eD \text{ } 100 \text{ fb}^{-1}$
- $10 \text{ GeV} \times 275 \text{ GeV } ep \text{ } 100 \text{ fb}^{-1}$
- $18 \text{ GeV} \times 137 \text{ GeV } eD \text{ } 15.4 \text{ fb}^{-1}$
- $18 \text{ GeV} \times 275 \text{ GeV } ep \text{ } 15.4 \text{ fb}^{-1}$

Data sets

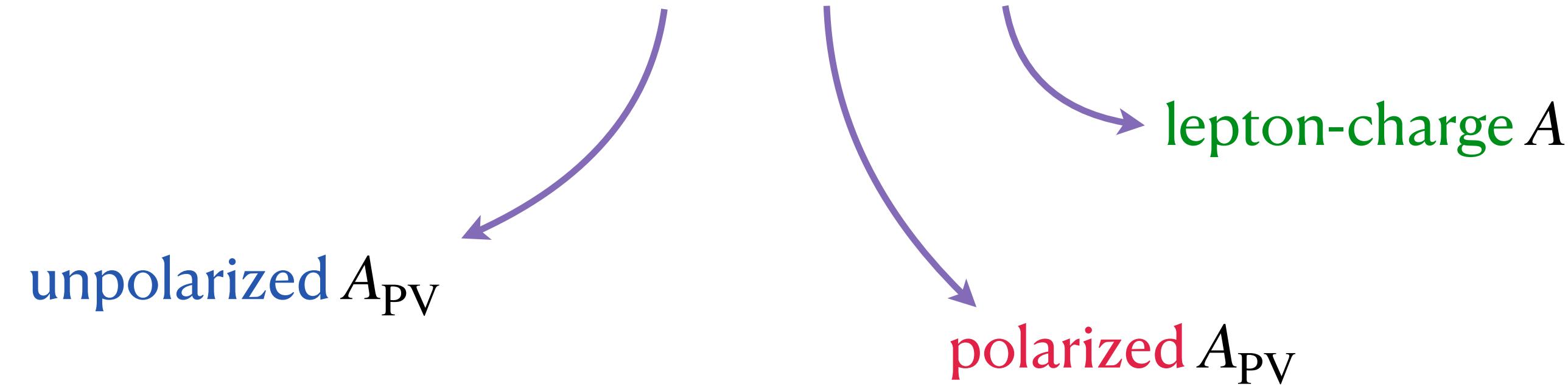
Concentrate on the two highest-energy settings:

- $10 \text{ GeV} \times 137 \text{ GeV } eD \text{ } 100 \text{ fb}^{-1}$
 - $10 \text{ GeV} \times 275 \text{ GeV } ep \text{ } 100 \text{ fb}^{-1}$
 - $18 \text{ GeV} \times 137 \text{ GeV } eD \text{ } 15.4 \text{ fb}^{-1}$
 - $18 \text{ GeV} \times 275 \text{ GeV } ep \text{ } 15.4 \text{ fb}^{-1}$
- 
- low energy
high luminosity
- high energy
low luminosity

Data sets

Concentrate on the two highest-energy settings:

- $10 \text{ GeV} \times 137 \text{ GeV } eD \text{ } 100 \text{ fb}^{-1}$: D4, Δ D4, LD4
- $10 \text{ GeV} \times 275 \text{ GeV } ep \text{ } 100 \text{ fb}^{-1}$: P4, Δ P4, LP4
- $18 \text{ GeV} \times 137 \text{ GeV } eD \text{ } 15.4 \text{ fb}^{-1}$: D5, Δ D5, LD5
- $18 \text{ GeV} \times 275 \text{ GeV } ep \text{ } 15.4 \text{ fb}^{-1}$: P5, Δ P5, LP5



Pseudodata generation

For b^{th} bin:

$$A_b^{\text{pseudo}} = A_b^{\text{theo}} + r_b \sigma_b^{\text{unc}} + r' \sigma_b^{\text{cor}}$$

where

A_b^{theo} : Born-level SM prediction

$$\sigma_b^{\text{unc}} = \begin{cases} \sigma_{\text{stat},b} \oplus \sigma_{\text{sys},b} \\ \sigma_{\text{stat},b} \oplus \sigma_{\text{sys},b}^{\text{unc}} \oplus \sigma_{\text{nlo},b}^{\text{unc}} \end{cases}, \quad \sigma_b^{\text{cor}} = \begin{cases} \sigma_{\text{pol},b}^{\text{cor}} & (\text{PV}) \\ \sigma_{\text{lum},b}^{\text{cor}} & (\text{LC}) \end{cases}$$

$r_b, r' \sim \mathcal{N}(0,1)$: random numbers

SMEFT asymmetry corrections

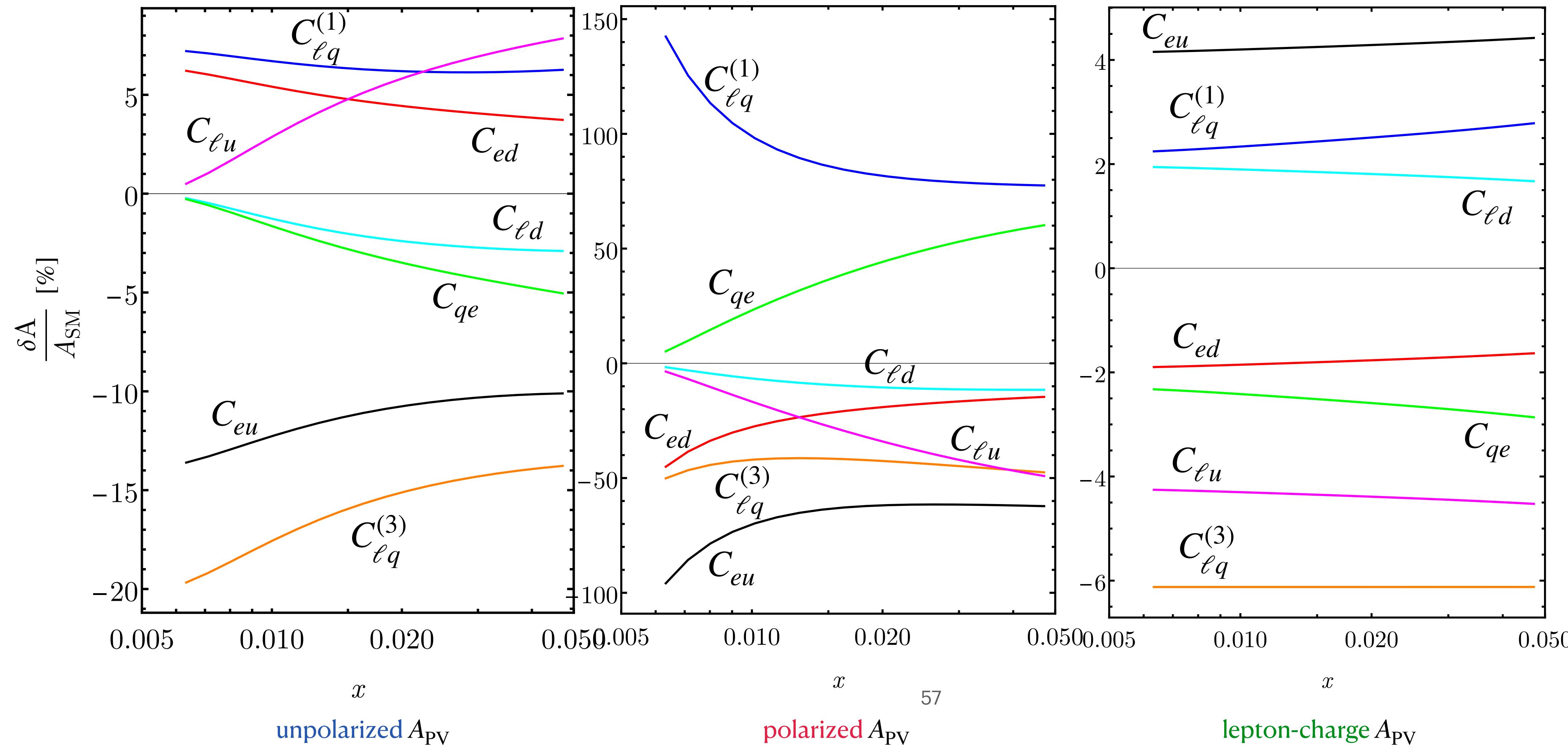
- Set $\Lambda = 1$ TeV.
- Turn on only one or two Wilson coefficients at a time and linearize:

$$A_{\text{SMEFT}}(x, Q^2, C) = A_{\text{SM}}(x, Q^2) + C \delta A(x, Q^2)$$

$$A_{\text{SMEFT}}(x, Q^2, C_1, C_2) = A_{\text{SM}}(x, Q^2) + C_1 \delta A_1(x, Q^2) + C_2 \delta A_2(x, Q^2)$$

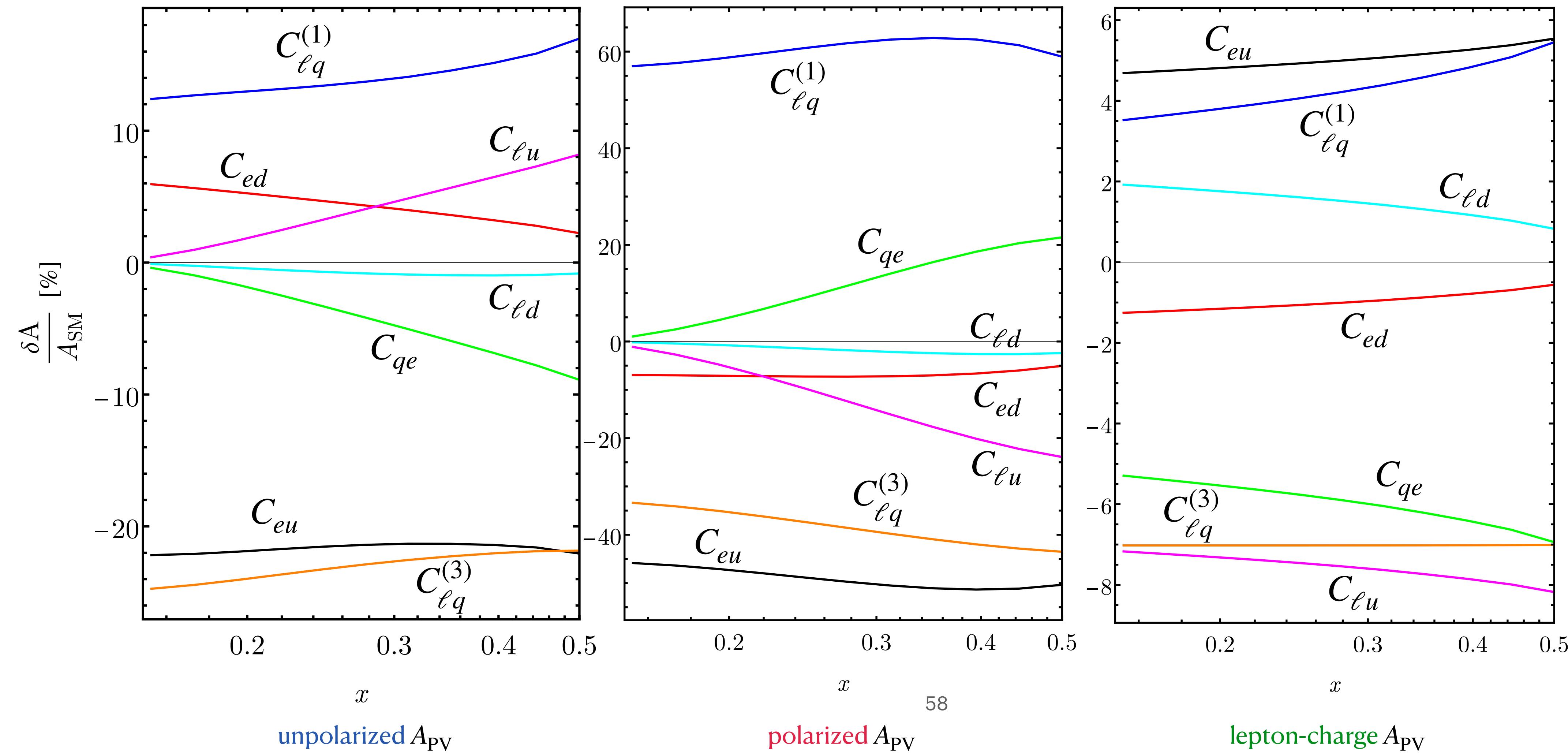
Sensitivity of SMEFT terms

$$A_{\text{SMEFT}}(x, Q^2, C) = A_{\text{SM}}(x, Q^2) + C \delta A(x, Q^2)$$



Sensitivity of SMEFT terms

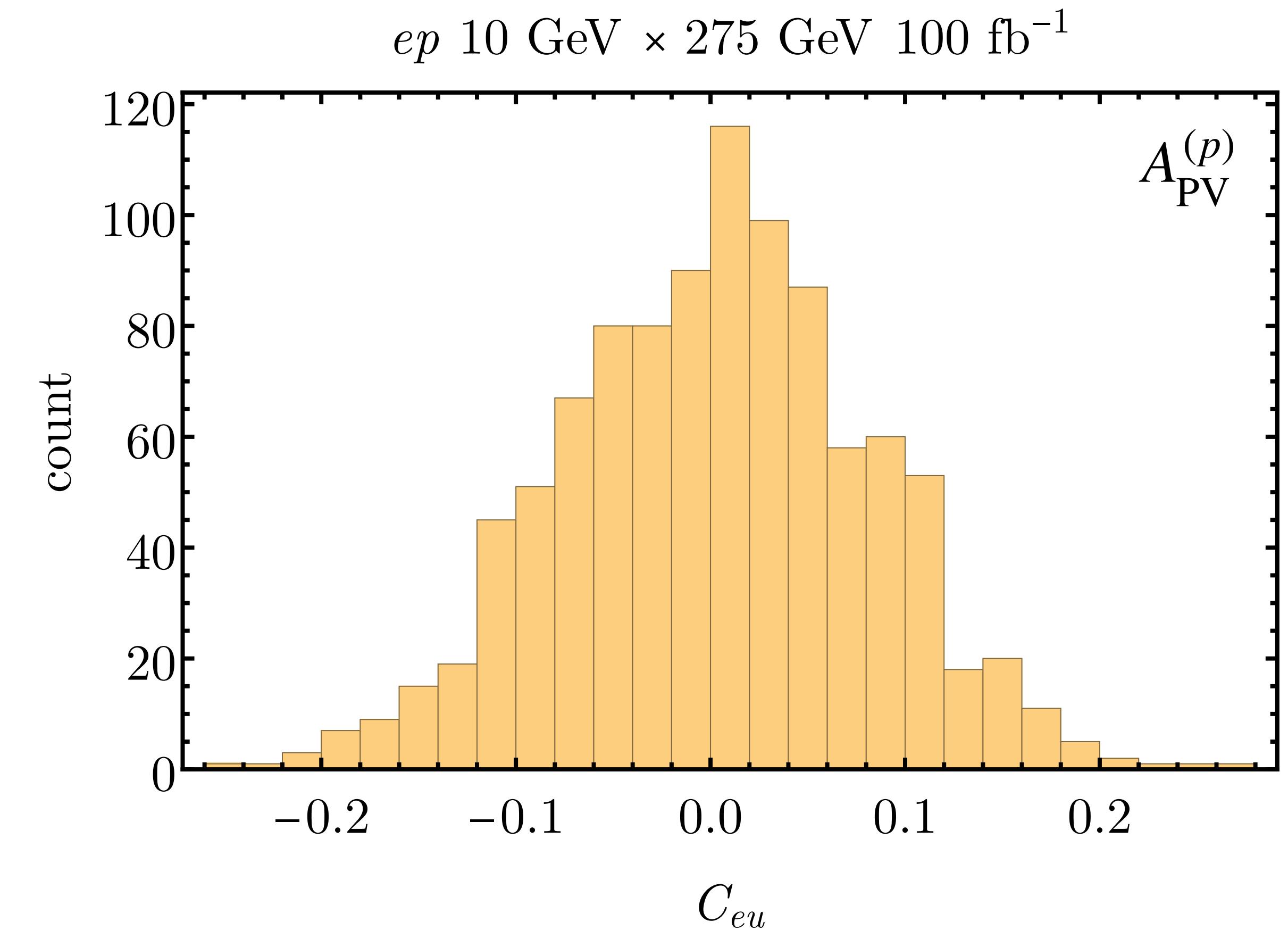
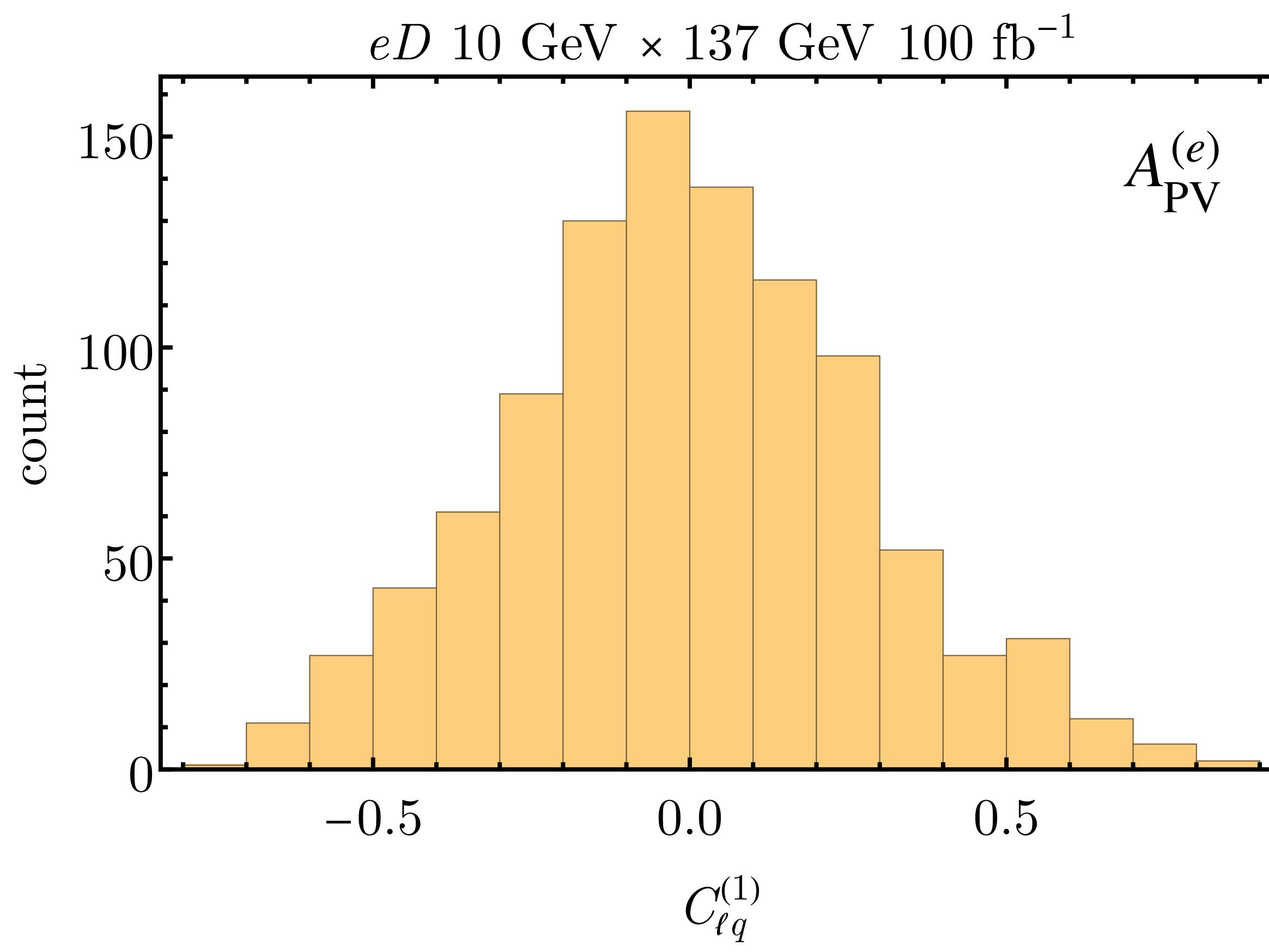
$$A_{\text{SMEFT}}(x, Q^2, C) = A_{\text{SM}}(x, Q^2) + C \delta A(x, Q^2)$$



Comments:

- ep
- $Q = 50 \text{ GeV}$
- $\sqrt{s} = 140 \text{ GeV}$

Distribution of best-fit values of Wilson coefficients



Best-fit analysis

χ^2 test statistic for the fits of Wilson coefficients:

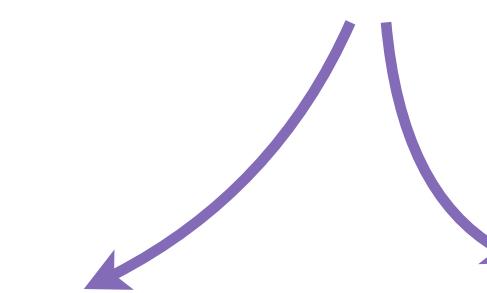
$$\chi^2 = \sum_{b=1}^{N_{\text{bin}}} \sum_{b'=1}^{N_{\text{bin}}} [A_{\text{SMEFT}} - A^{\text{pseudo}}]_b H_{bb'} [A_{\text{SMEFT}} - A^{\text{pseudo}}]_{b'}$$

PV asymmetries:

$$\sigma_{\text{stat}}, \sigma_{\text{sys}}^{\text{unc}}, \sigma_{\text{pol}}^{\text{cor}}, \sigma_{\text{pdf}}^{\text{cor}}$$

LC asymmetries:

$$\sigma_{\text{stat}}, \sigma_{\text{sys}}^{\text{unc}}, \sigma_{\text{pol}}^{\text{lum}}, \sigma_{\text{nlo}}^{\text{unc}}, \sigma_{\text{pdf}}^{\text{cor}}$$



Best-fit analysis

χ^2 test statistic for the fits of Wilson coefficients:

$$\chi^2 = \sum_{b=1}^{N_{\text{bin}}} \sum_{b'=1}^{N_{\text{bin}}} [A_{\text{SMEFT}} - A^{\text{pseudo}}]_b H_{bb'} [A_{\text{SMEFT}} - A^{\text{pseudo}}]_{b'}$$

Polarimetry and luminosity difference can be limiting factors.

- ⇒ use data itself to constrain these systematic effects
- ⇒ simultaneous fits of Wilson coefficients with beam polarization, P
- ⇒ simultaneous fits of Wilson coefficients with luminosity difference, A_{lum}

Best-fit analysis

χ^2 test statistic for the fits of Wilson coefficients:

$$\chi^2 = \sum_{b=1}^{N_{\text{bin}}} \sum_{b'=1}^{N_{\text{bin}}} [A_{\text{SMEFT}} - A^{\text{pseudo}}]_b H_{bb'} [A_{\text{SMEFT}} - A^{\text{pseudo}}]_{b'}$$

assumed 1.0
for simplicity

$$\chi^2 = \sum_{b=1}^{N_{\text{bin}}} \sum_{b'=1}^{N_{\text{bin}}} [P(A_{\text{PV}})_{\text{SMEFT}} - A^{\text{pseudo}}]_b \left[H \Big|_{\sigma_{\text{pol}} \rightarrow 0} \right]_{bb'} [P(A_{\text{PV}})_{\text{SMEFT}} - A^{\text{pseudo}}]_{b'} + \frac{(P - \bar{P})^2}{\delta P^2}$$

beam polarization
as a nuance parameter

beam polarization uncertainty
1 % for lepton beam
2 % for hadron beam

Best-fit analysis

χ^2 test statistic for the fits of Wilson coefficients:

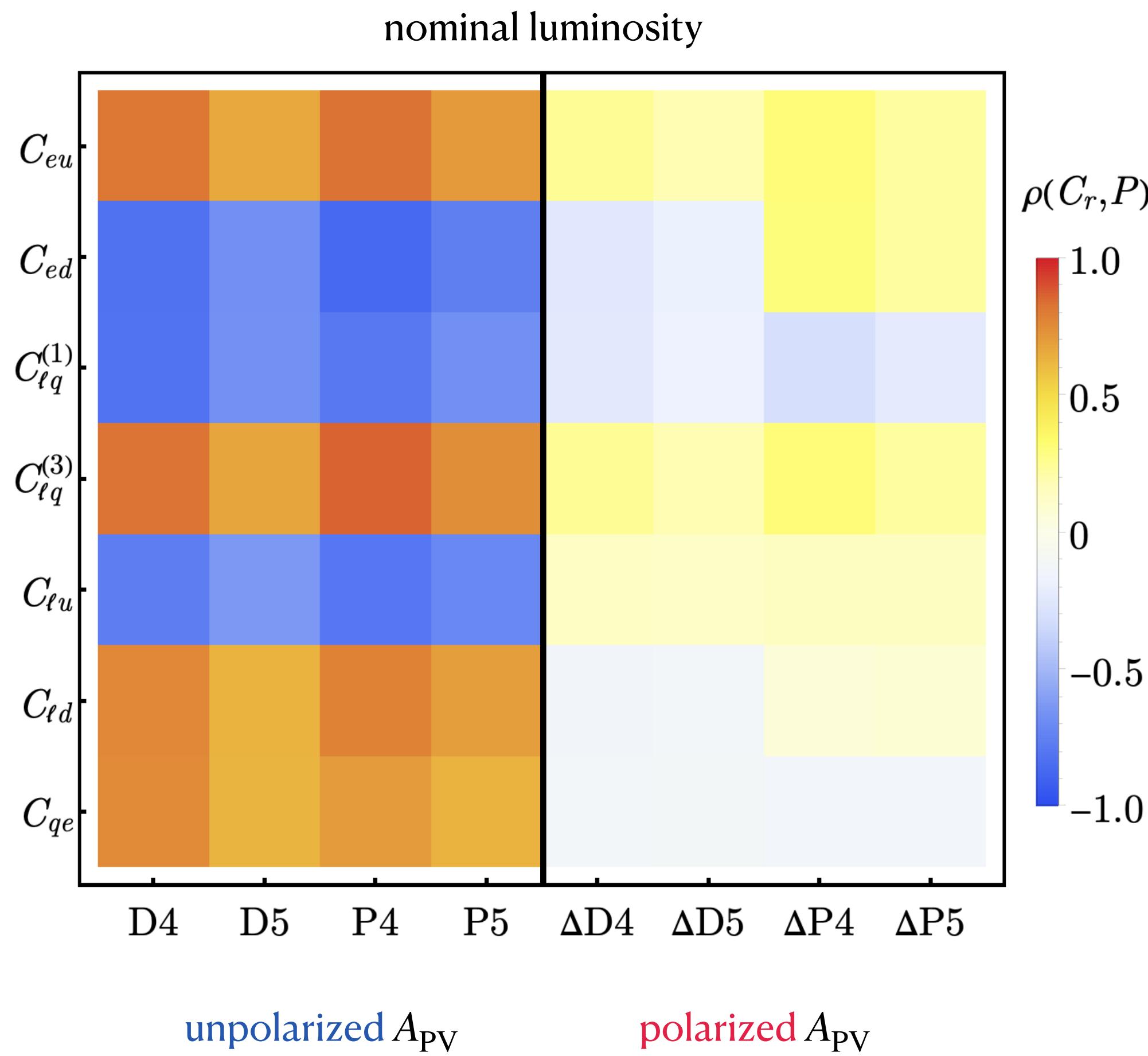
$$\chi^2 = \sum_{b=1}^{N_{\text{bin}}} \sum_{b'=1}^{N_{\text{bin}}} [A_{\text{SMEFT}} - A^{\text{pseudo}}]_b H_{bb'} [A_{\text{SMEFT}} - A^{\text{pseudo}}]_{b'}$$

$$\chi^2 = \sum_{b=1}^{N_{\text{bin}}} \sum_{b'=1}^{N_{\text{bin}}} [P(A_{\text{PV}})_{\text{SMEFT}} - A^{\text{pseudo}}]_b \left[H \Big|_{\sigma_{\text{pol}} \rightarrow 0} \right]_{bb'} [P(A_{\text{PV}})_{\text{SMEFT}} - A^{\text{pseudo}}]_{b'} + \frac{(P - \bar{P})^2}{\delta P^2}$$

$$\chi^2 = \sum_{b=1}^{N_{\text{bin}}} \sum_{b'=1}^{N_{\text{bin}}} [(A_{\text{LC}})_{\text{SMEFT}} - A^{\text{pseudo}} - A_{\text{lum}}]_b \left[H \Big|_{\sigma_{\text{lum}} \rightarrow 0} \right]_{bb'} [(A_{\text{LC}})_{\text{SMEFT}} - A^{\text{pseudo}} - A_{\text{lum}}]_{b'}$$

shift in pseudodata
by a residual amount

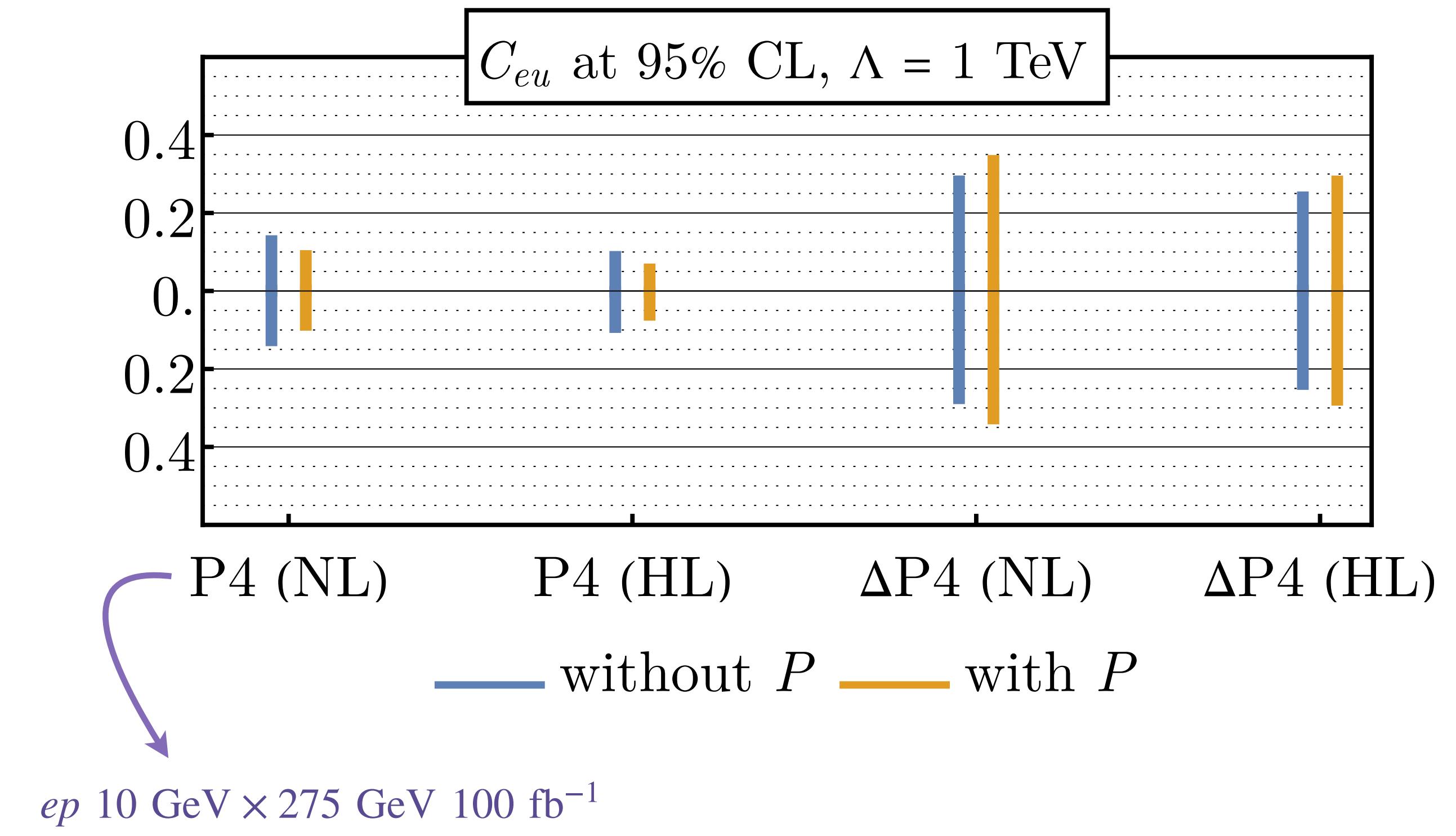
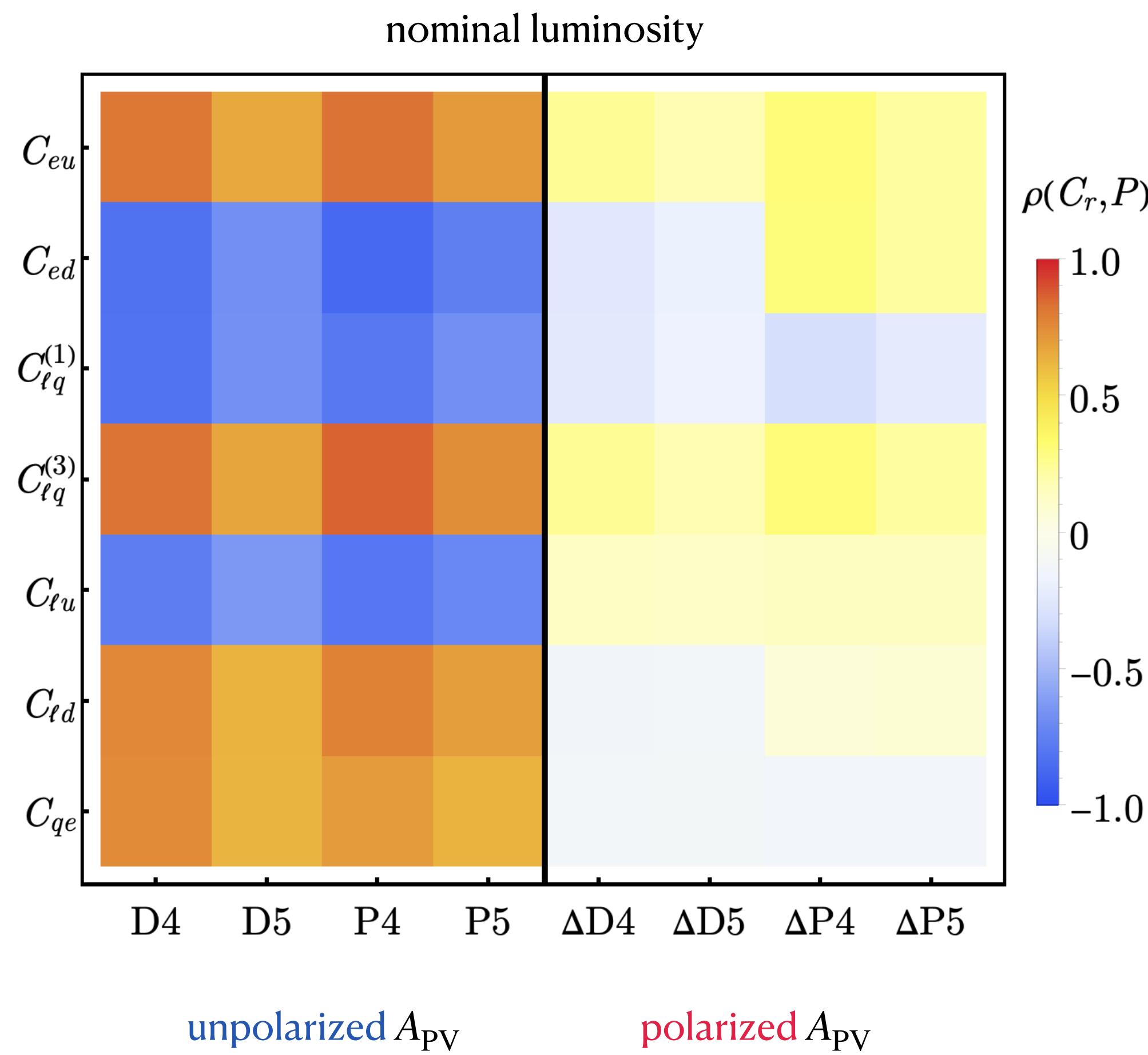
Simultaneous fits with polarization parameter



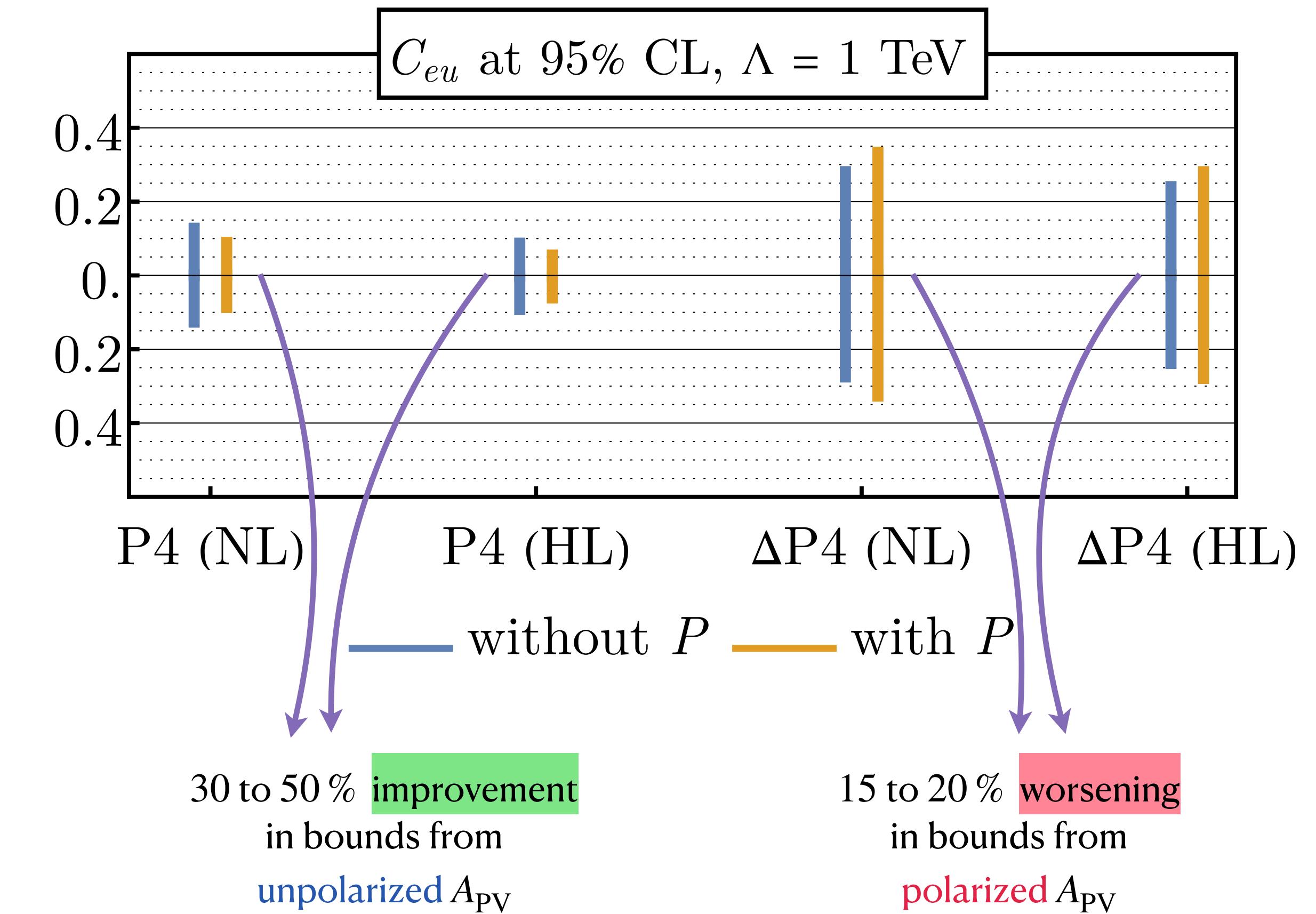
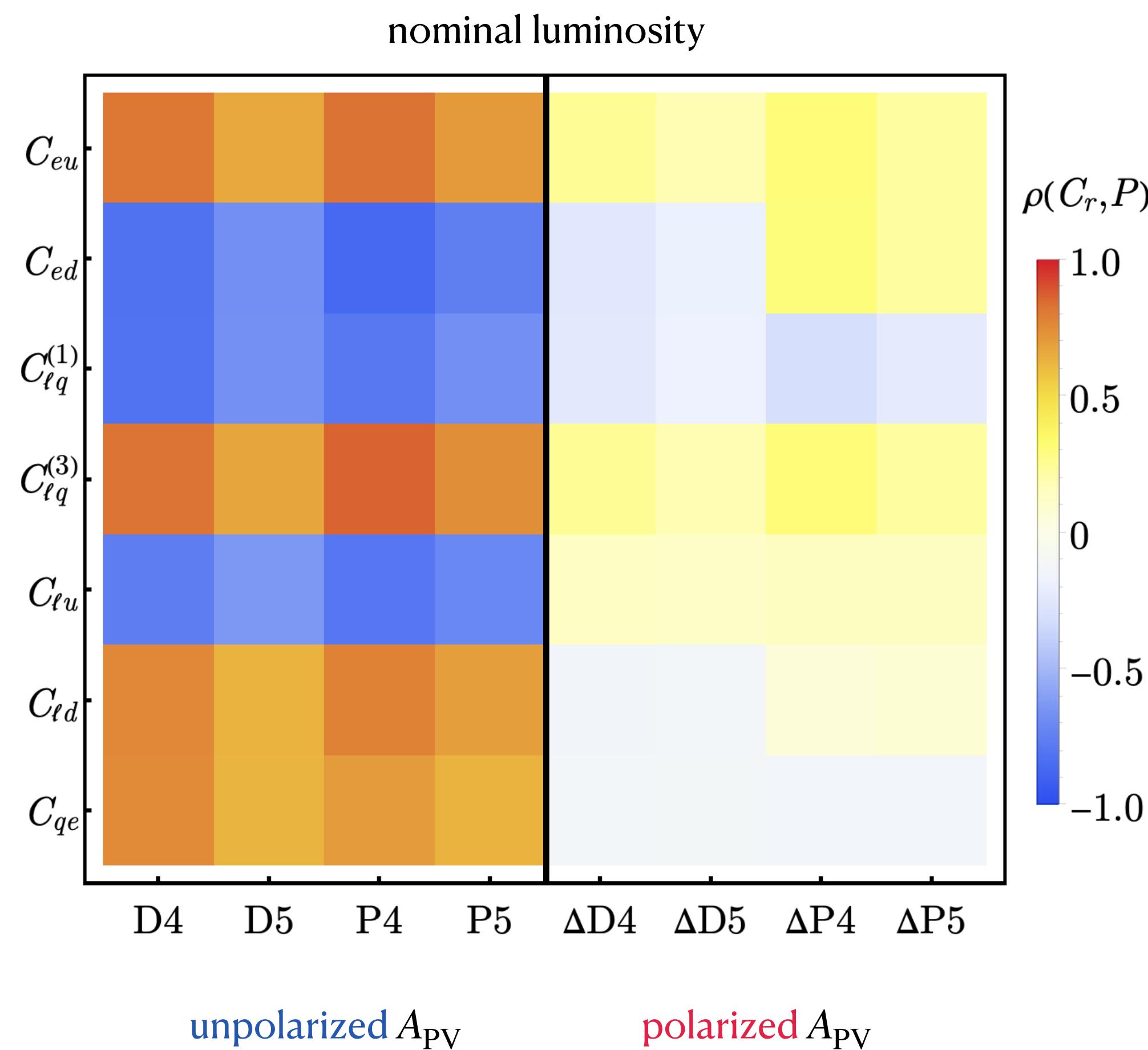
strong correlations between C_r and P
in unpolarized A_{PV} data sets

weak correlations between C_r and P
in polarized A_{PV} data sets

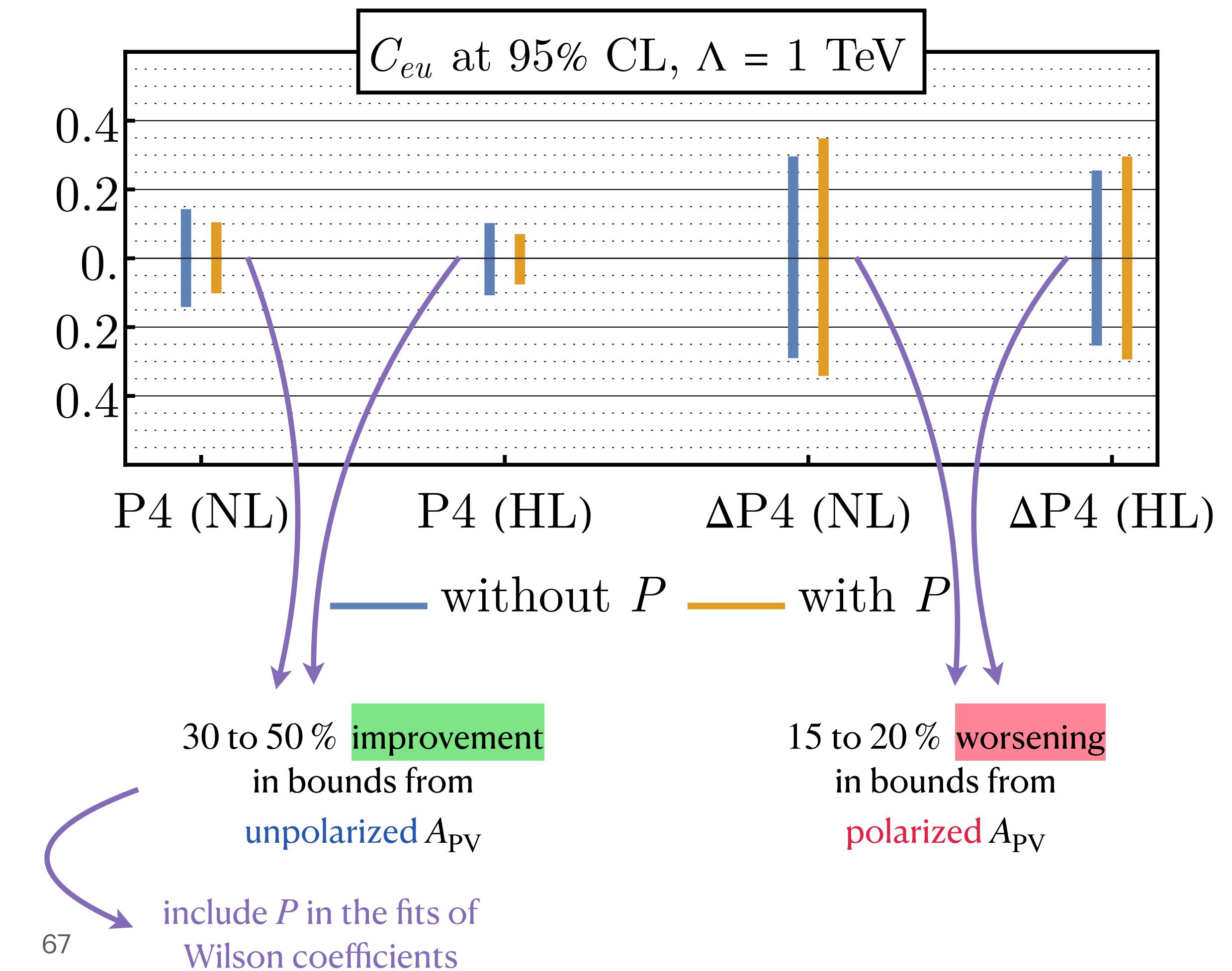
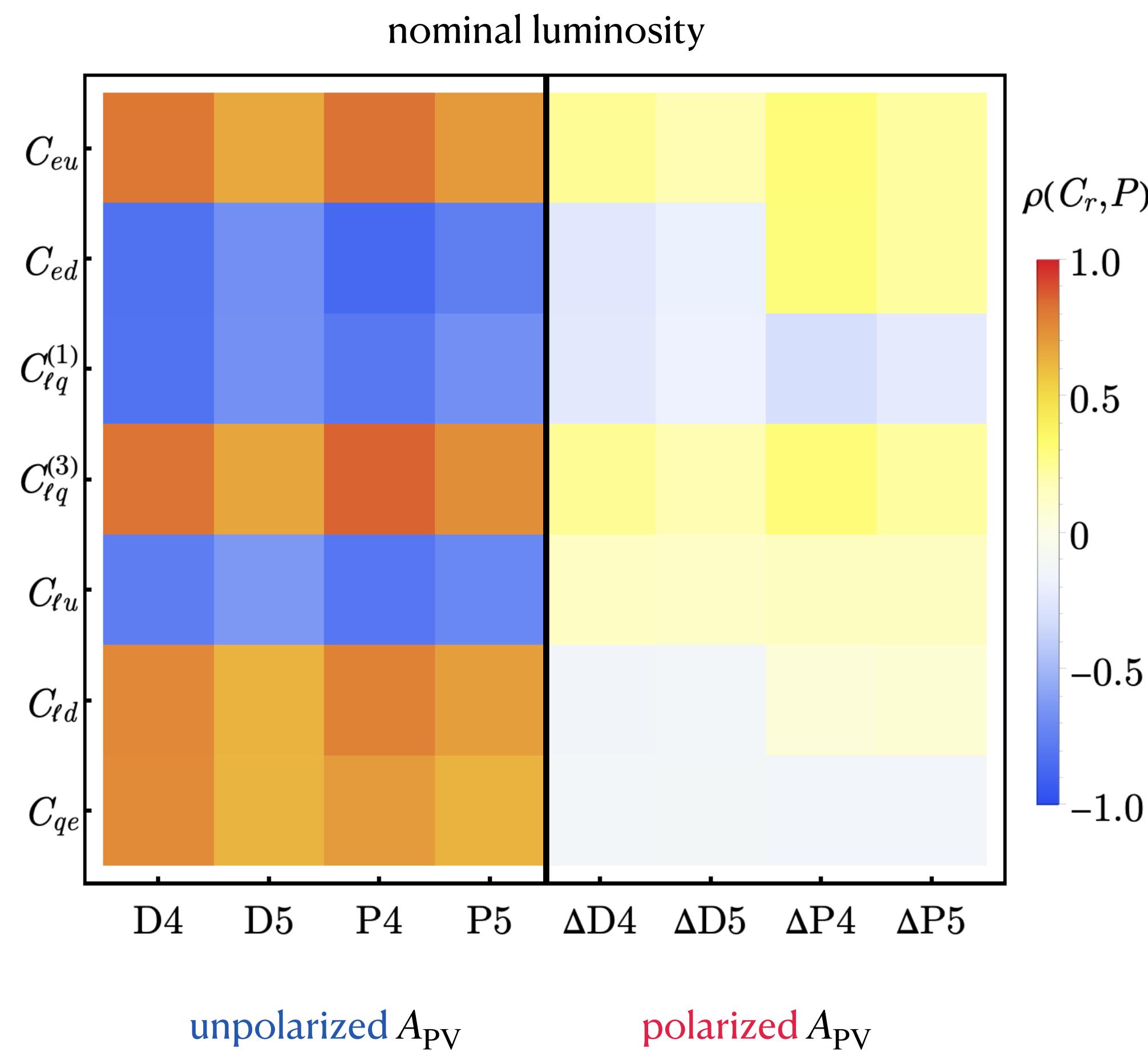
Simultaneous fits with polarization parameter



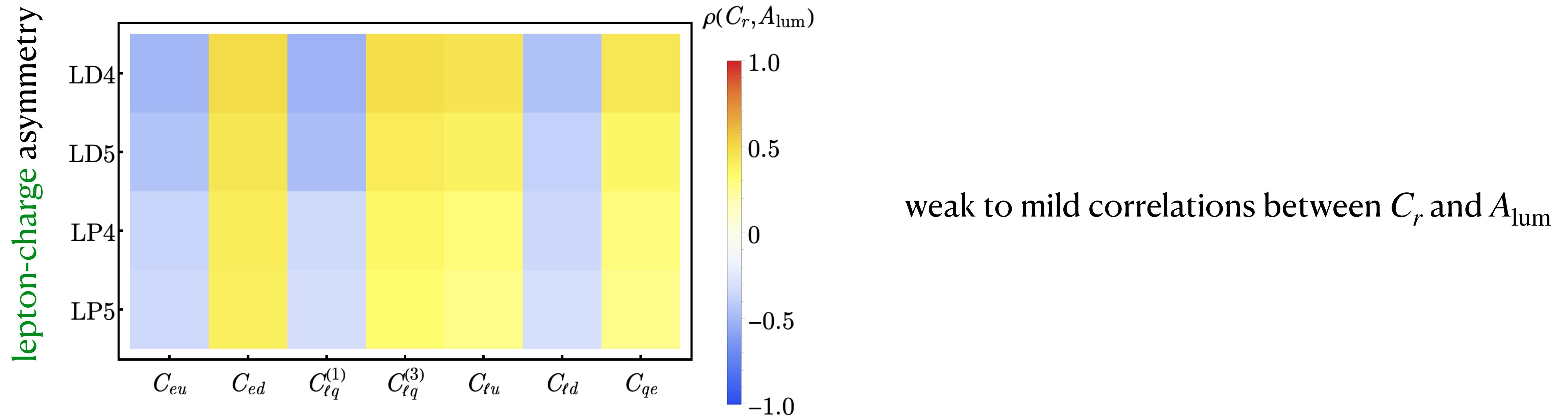
Simultaneous fits with polarization parameter



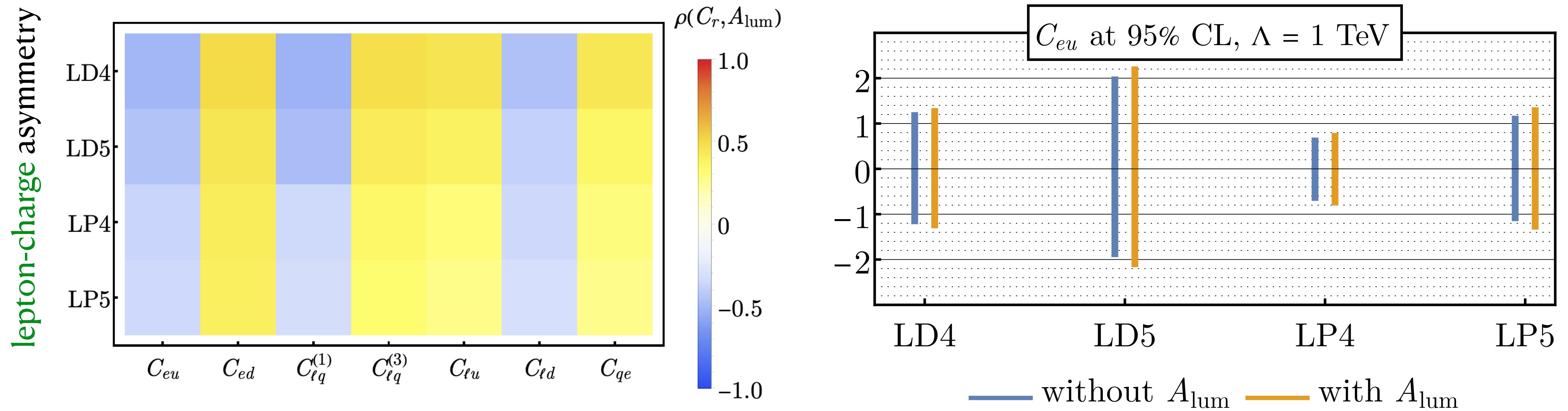
Simultaneous fits with polarization parameter



Simultaneous fits with luminosity difference



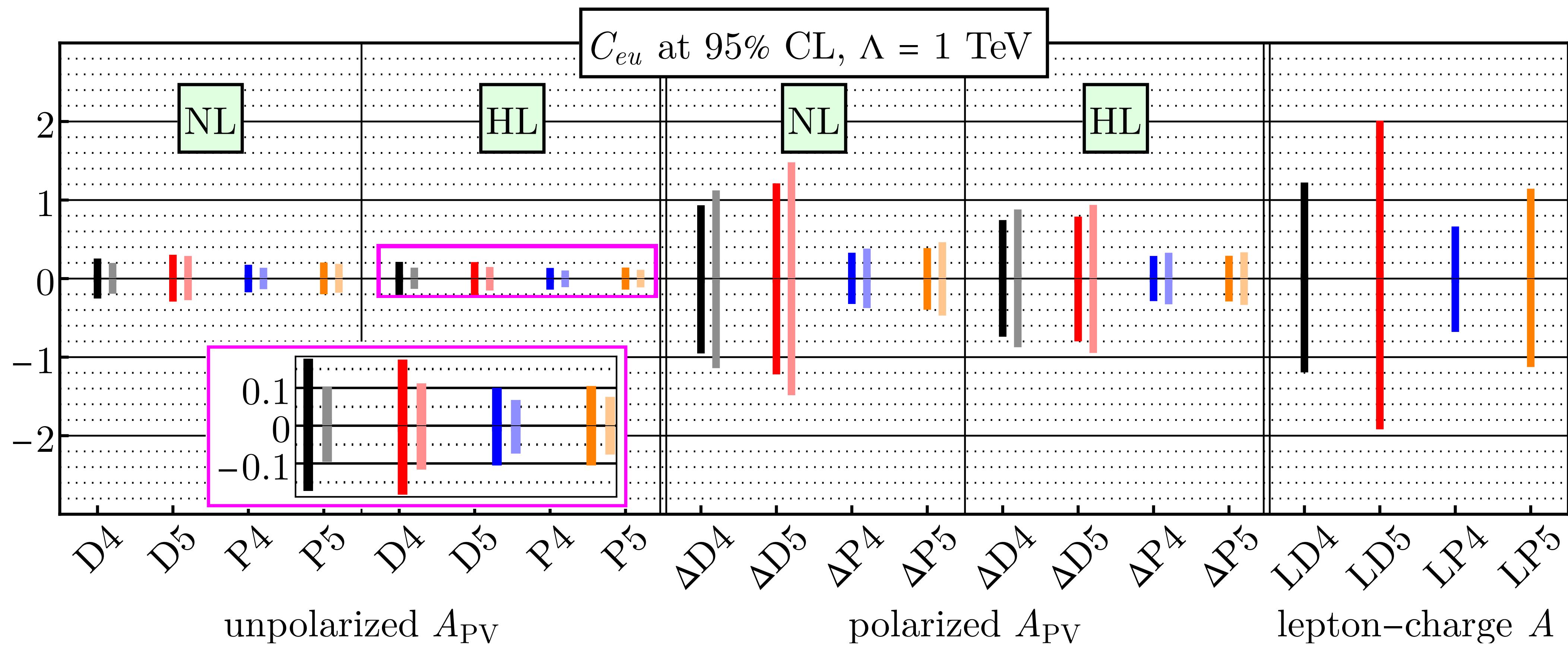
Simultaneous fits with luminosity difference



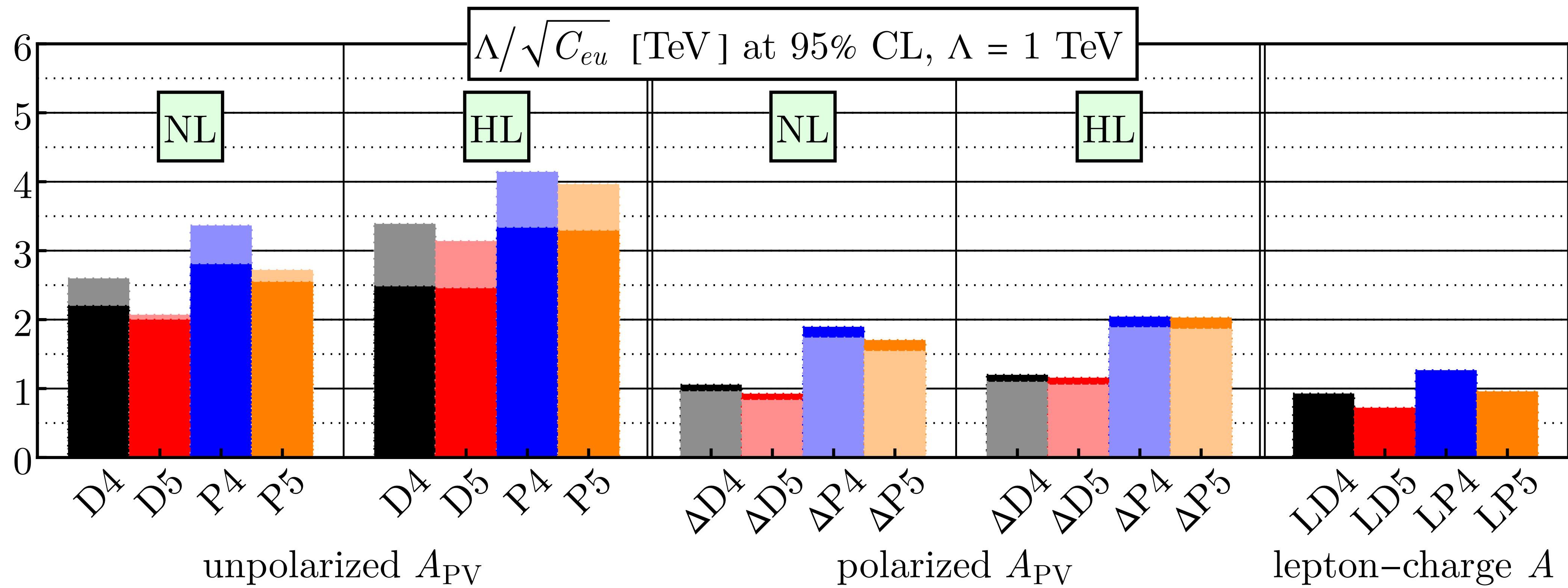
15 to 20 % worsening in
bounds from lepton-charge A
don't include A_{lum} in the fits of
Wilson coefficients

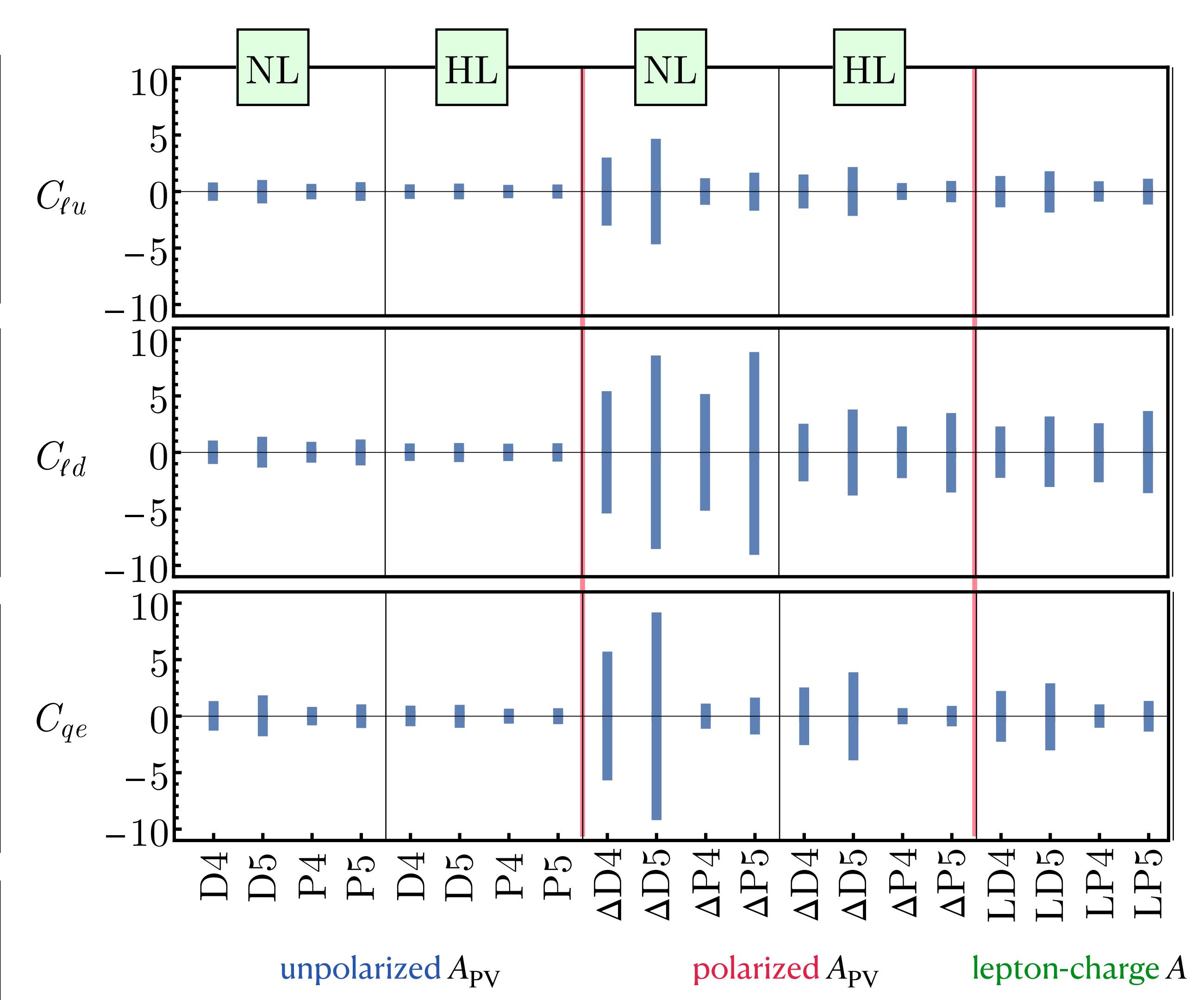
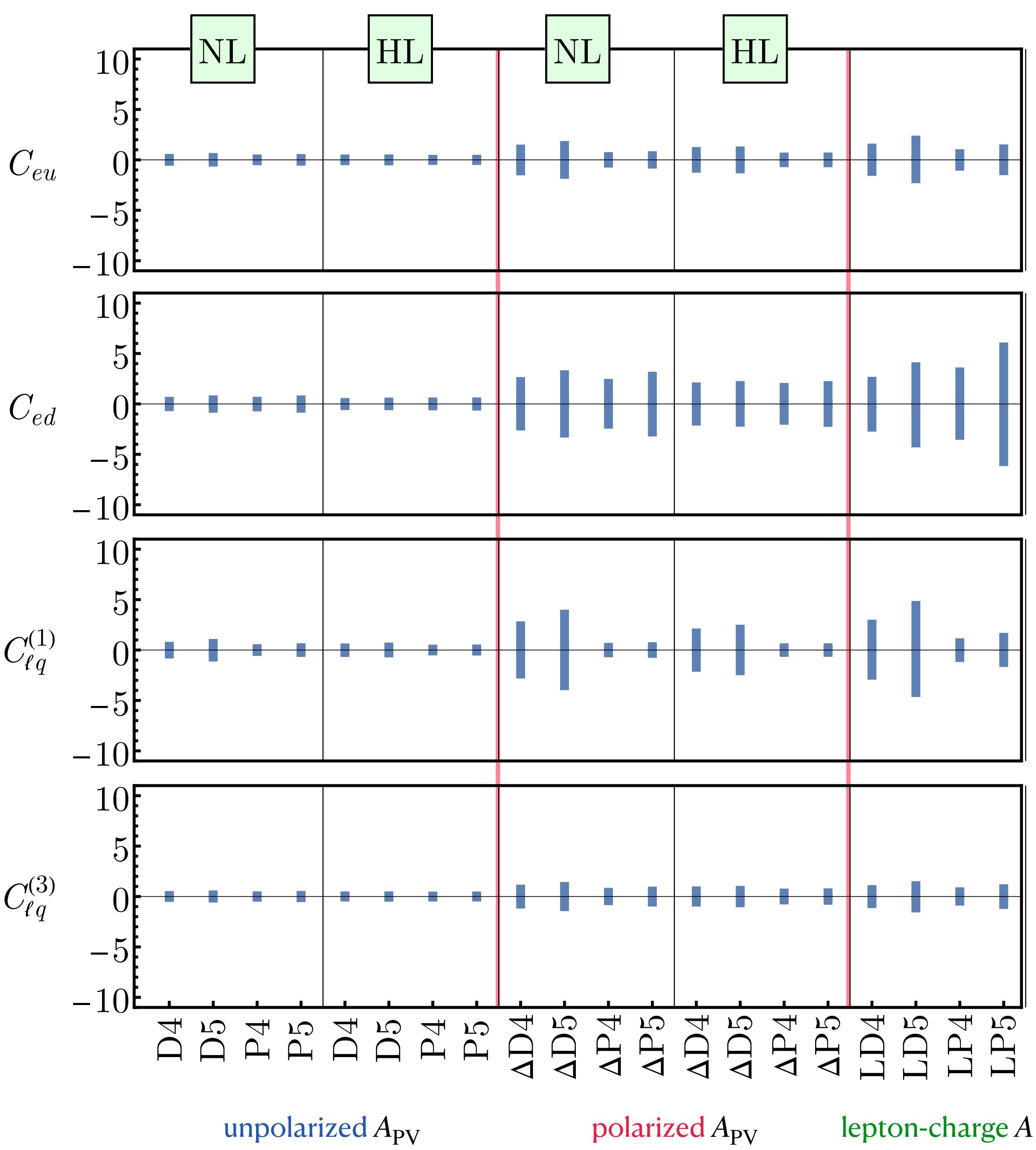
SMEFT fit results

Single Wilson coefficients



Single Wilson coefficients

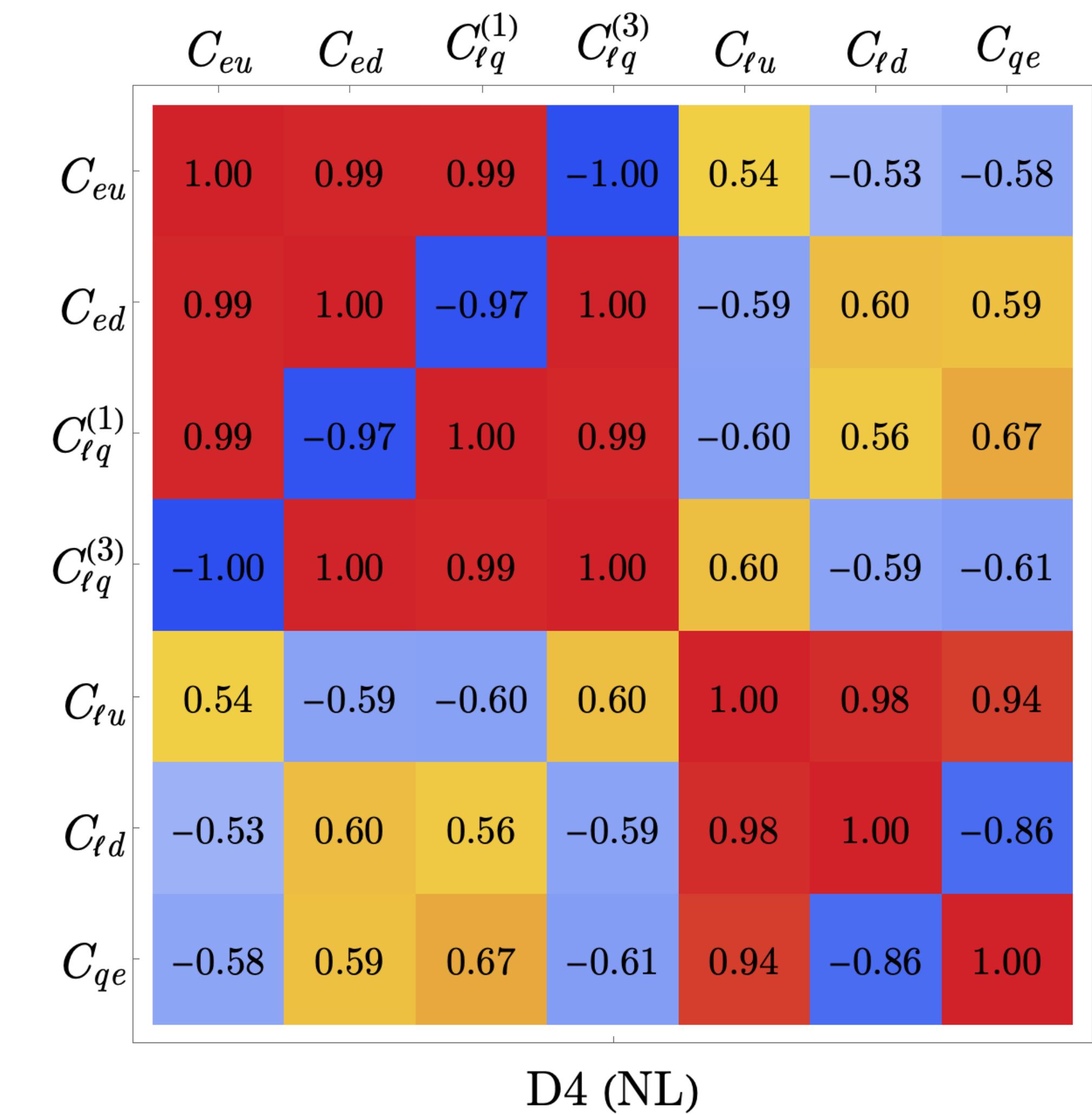
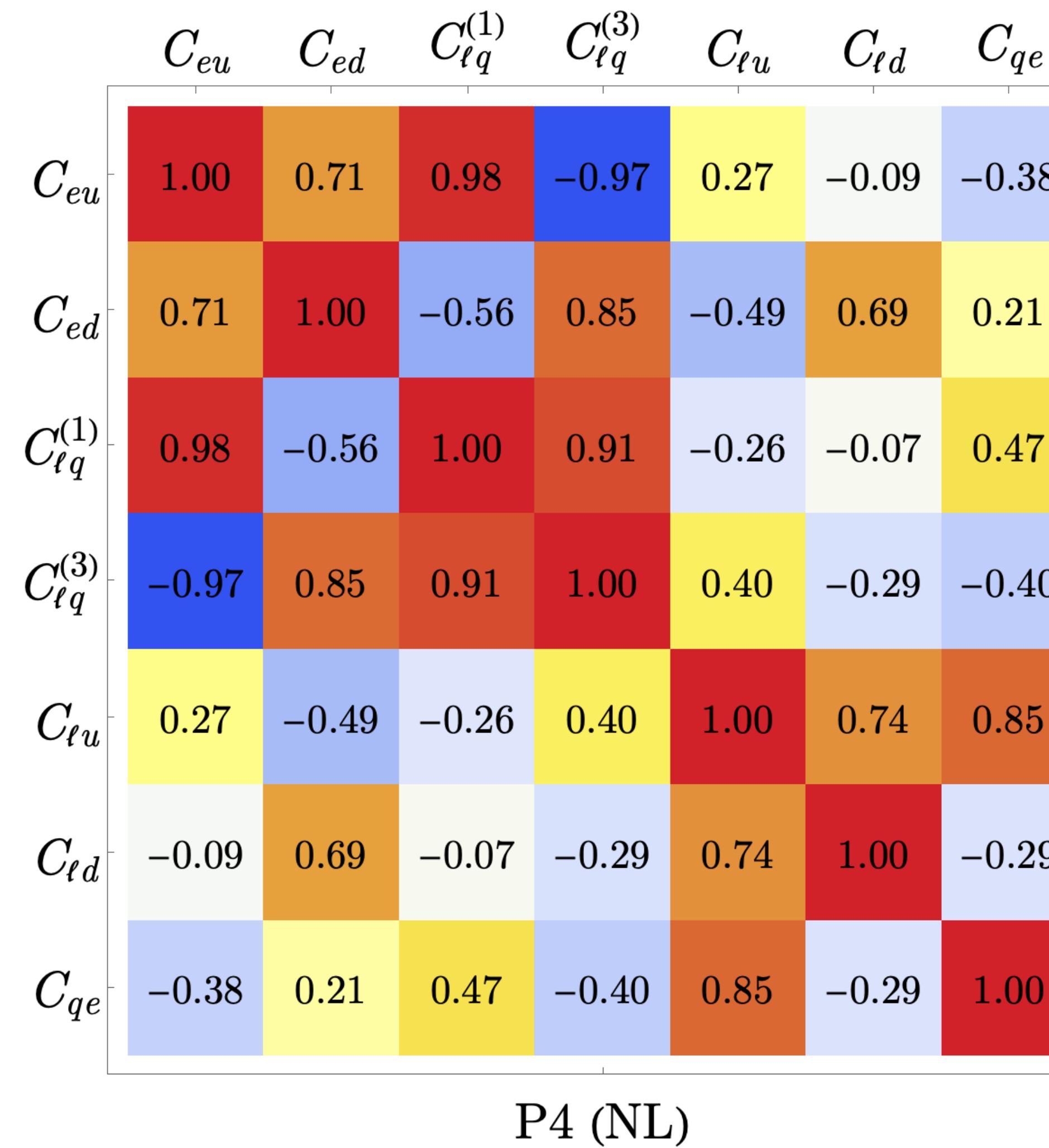




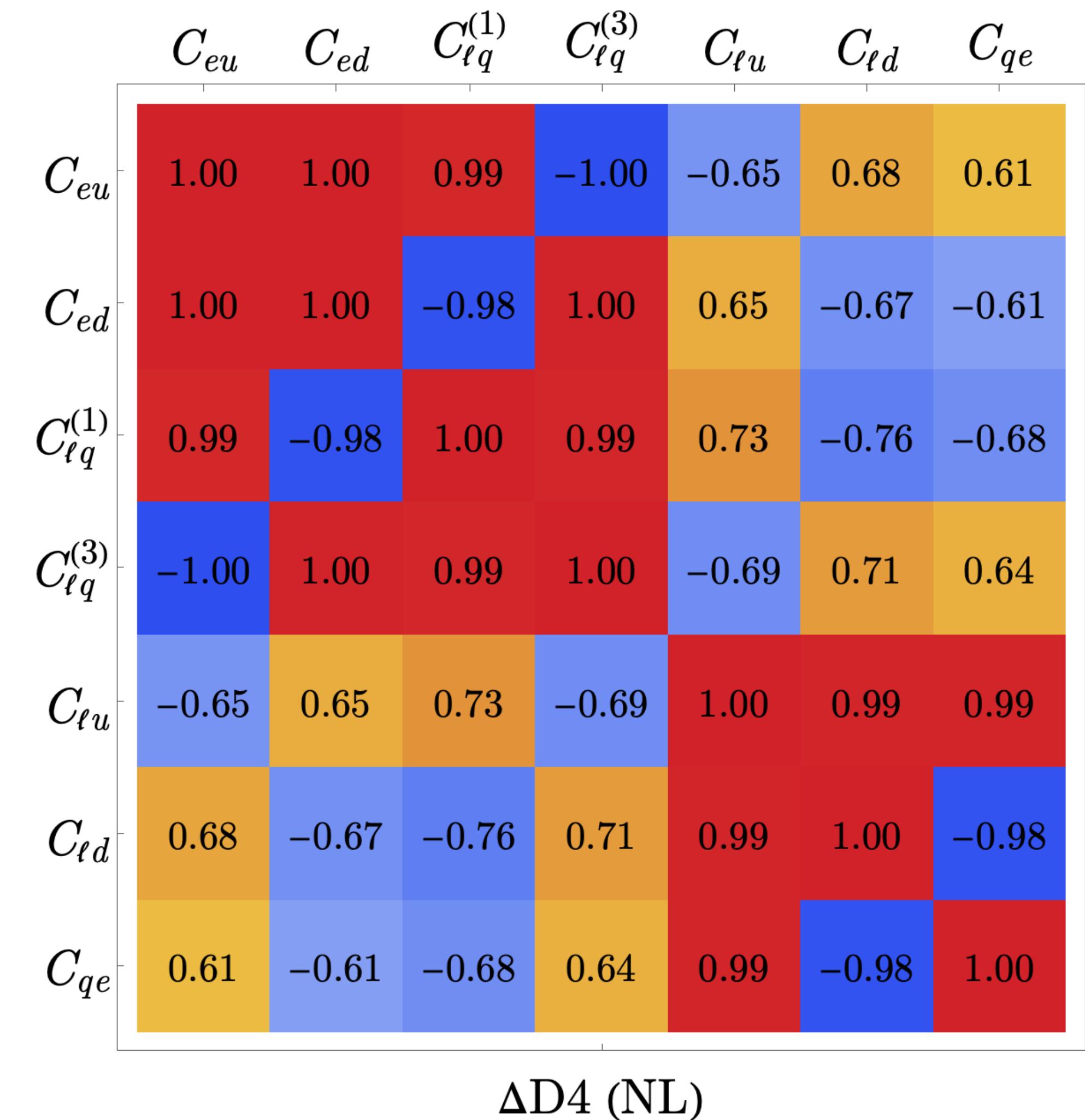
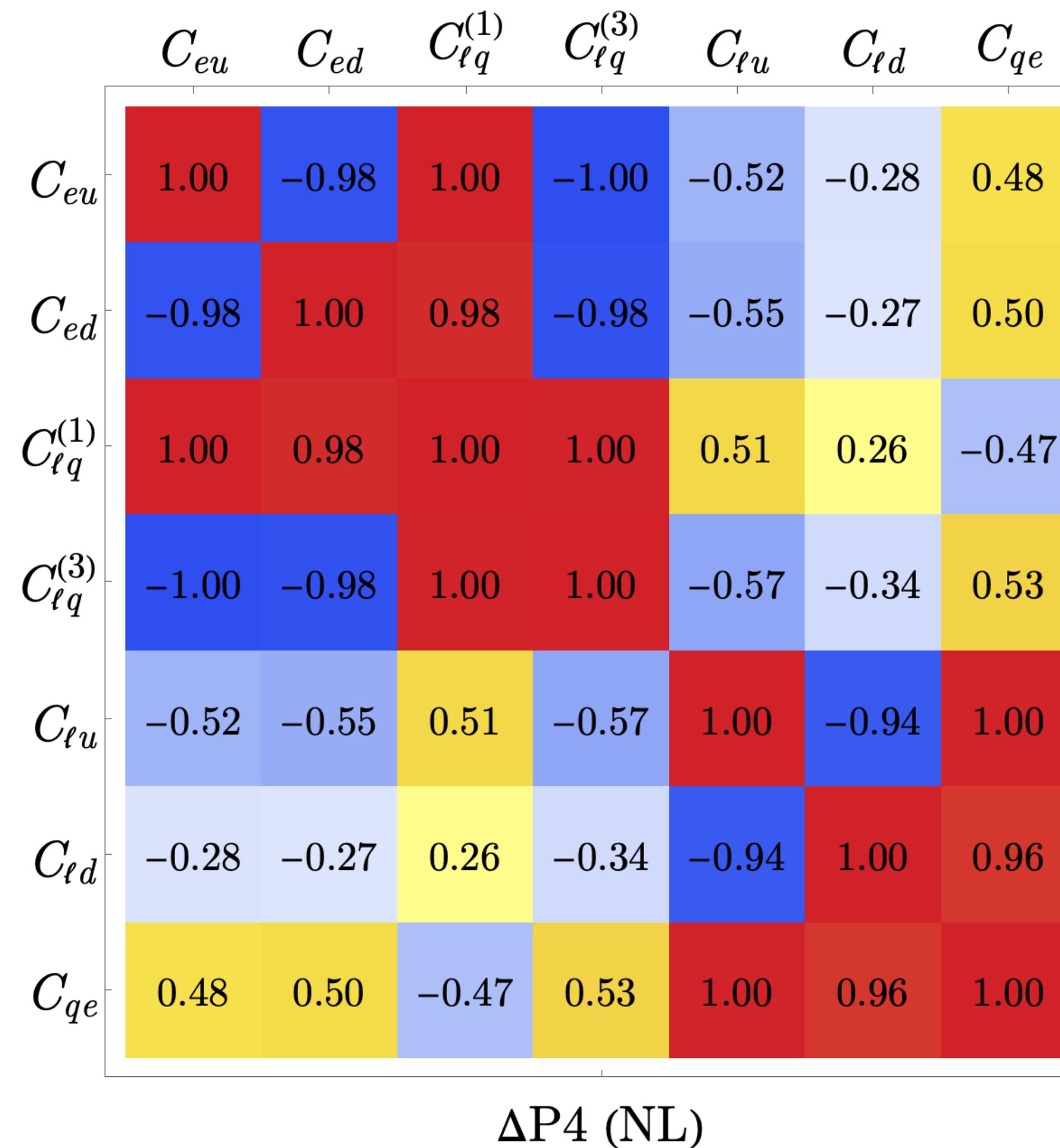
Comments:

- 95 % CL
- $\Lambda = 1$ TeV
- A_{PV} fits with beam polarization

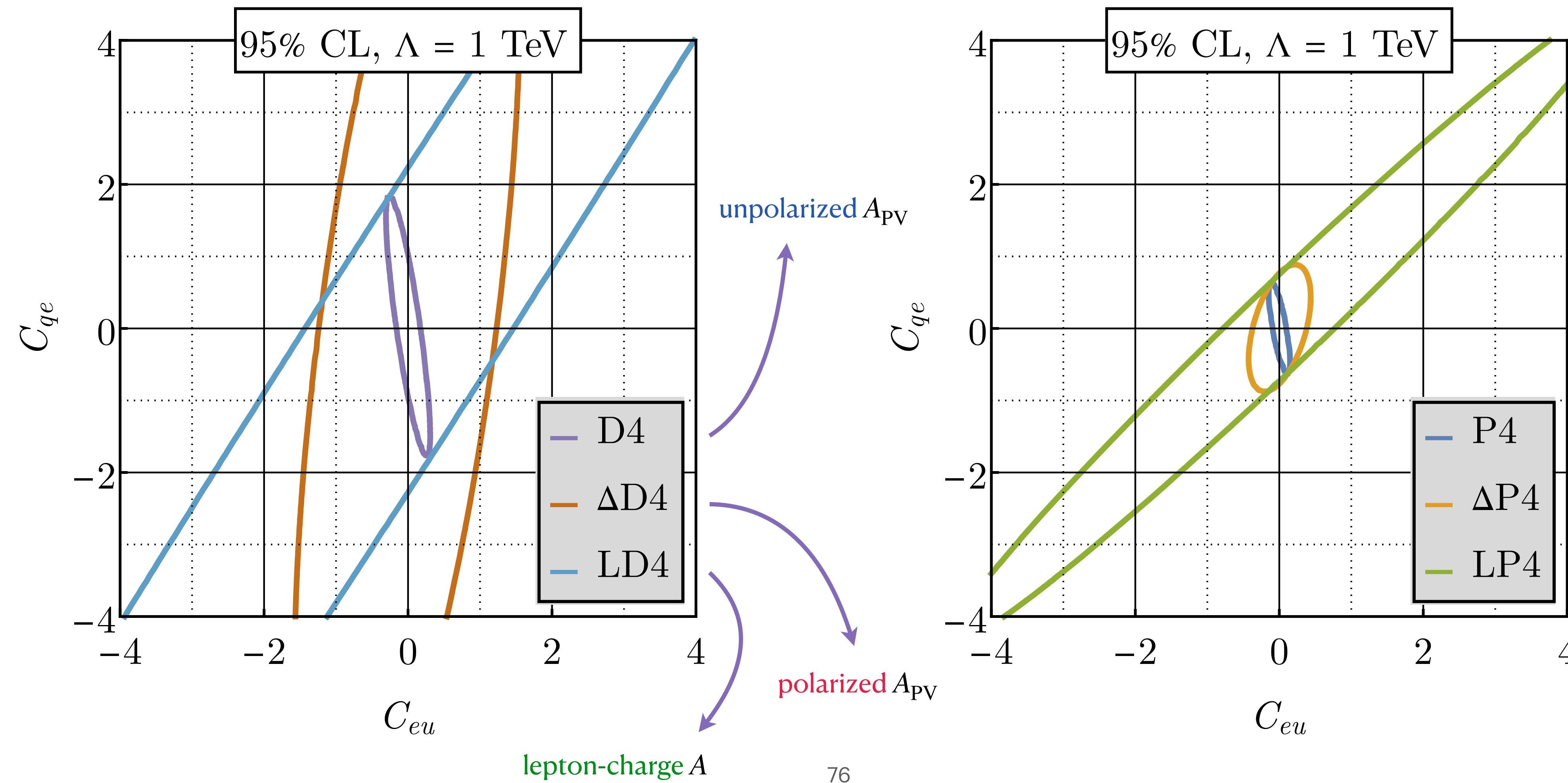
Double Wilson coefficients



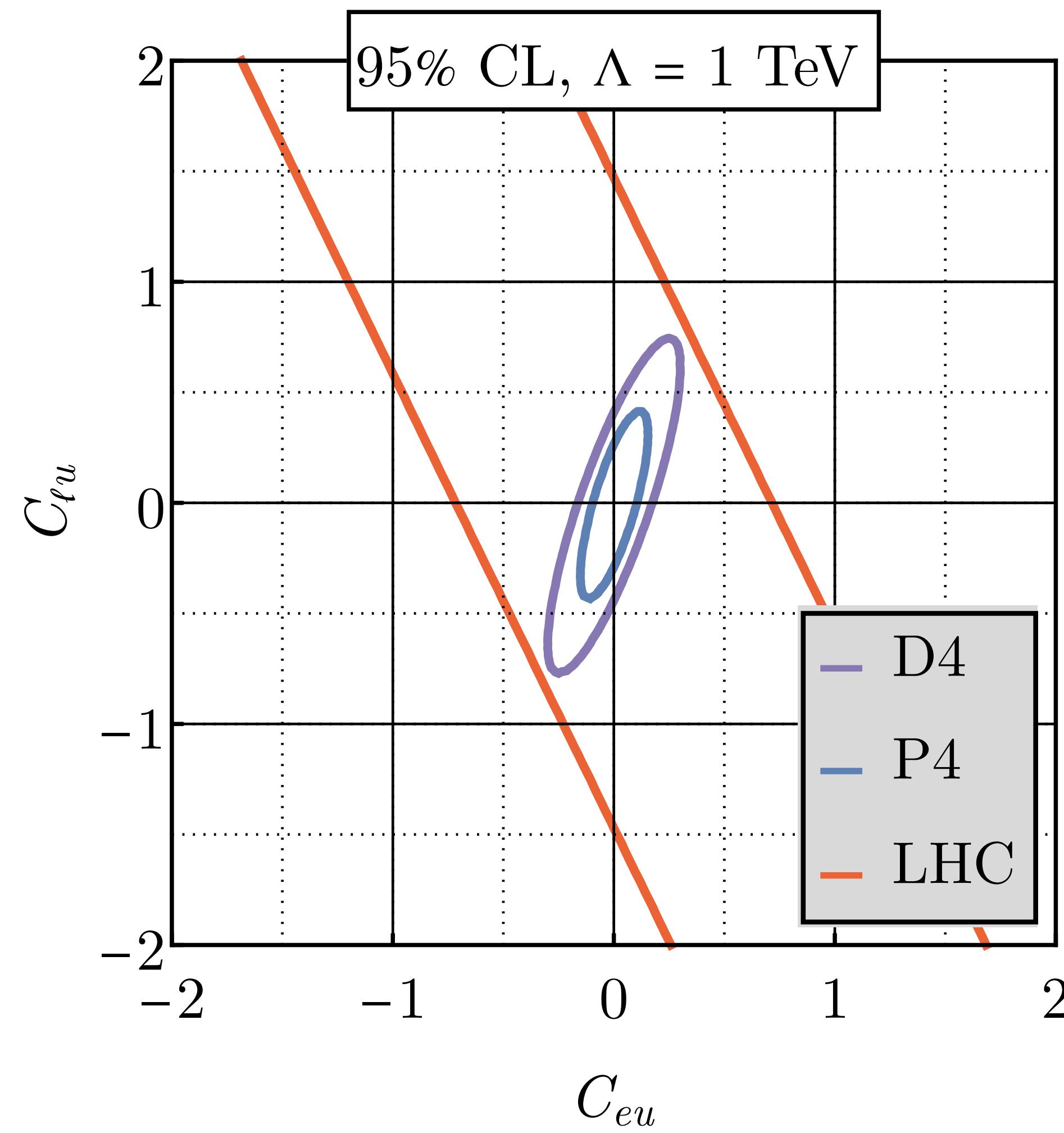
Double Wilson coefficients



Double Wilson coefficients



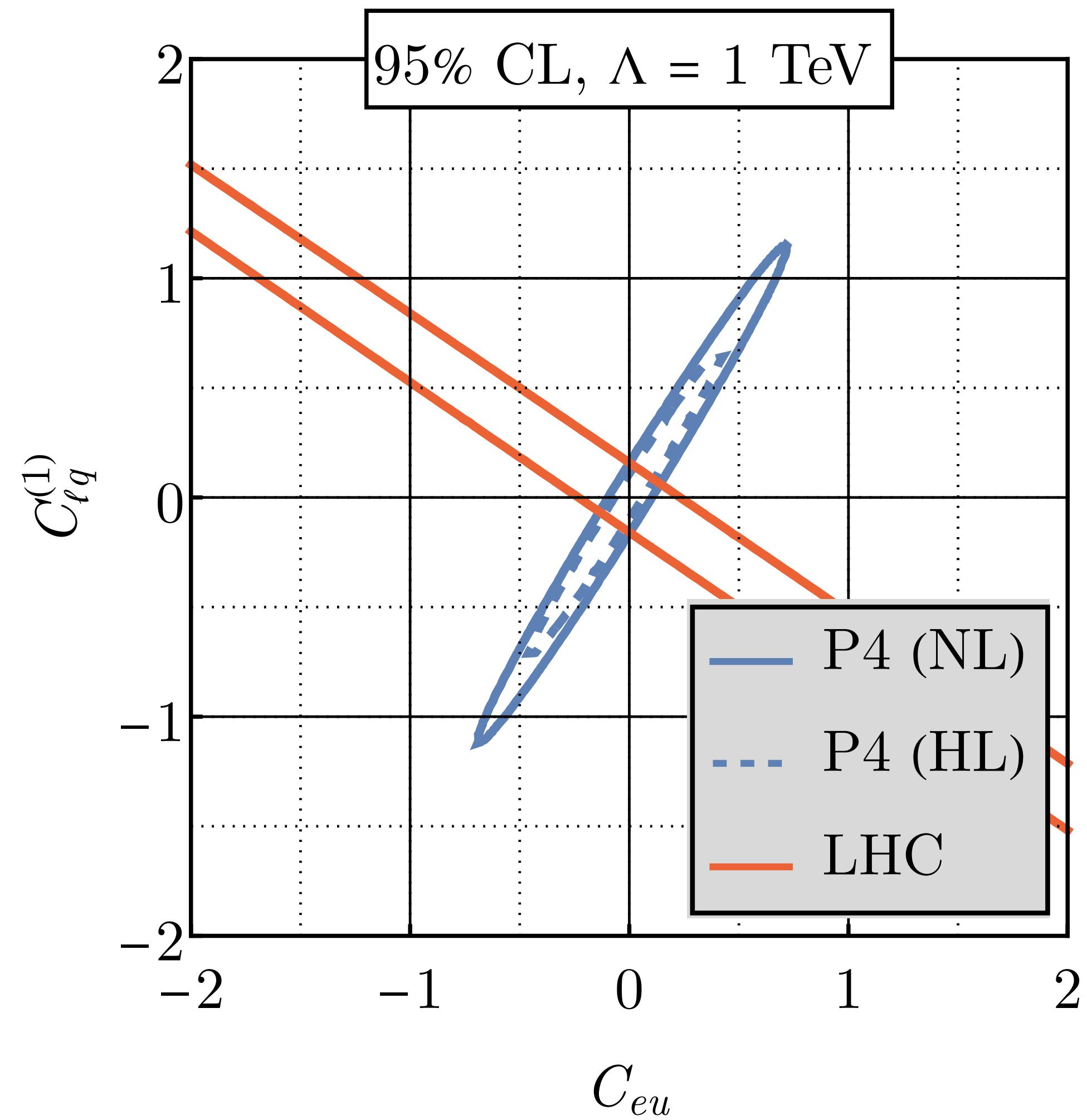
Double Wilson coefficients



LHC Drell-Yan adapted from
[RB, FP, DW 2104.03979](#)

LHC: 8 TeV and 20 fb^{-1}
not 13 TeV and HL

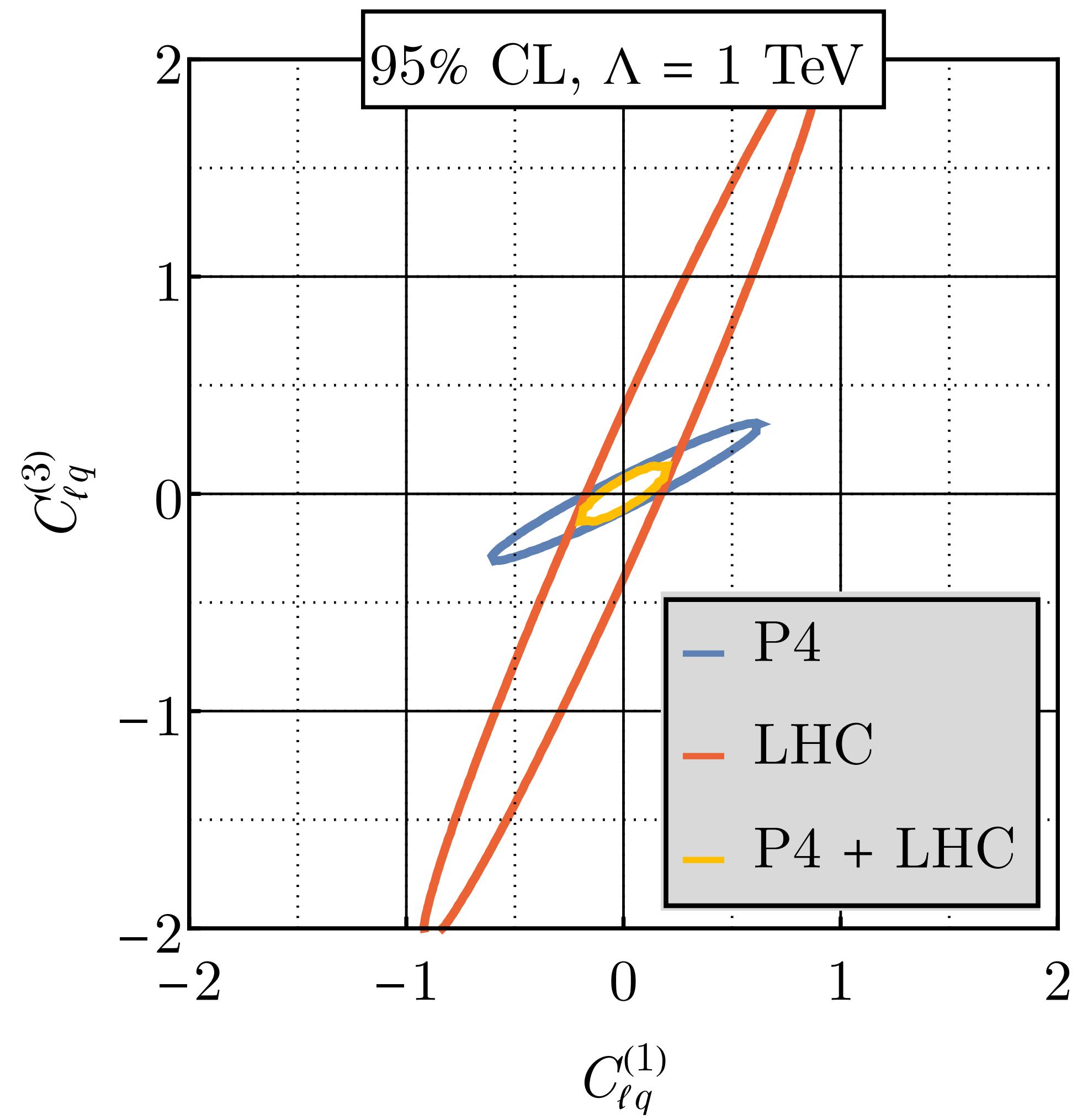
Double Wilson coefficients



LHC Drell-Yan adapted from
RB, FP, DW 2004.00748

LHC: 8 TeV and 20 fb^{-1}
not 13 TeV and HL

Double Wilson coefficients



LHC Drell-Yan adapted from
RB, FP, DW 2004.00748

LHC: 8 TeV and 20 fb^{-1}
not 13 TeV and HL

Summary of SMEFT fits

- Bound strength: proton > deuteron
- Bound strength: **unpolarized** A_{PV} > **polarized** A_{PV} > **lepton-charge** A
- All distinct correlations
- Complementary to LHC
- Flat directions in LHC Drell-Yan resolved by EIC
- Even stronger bounds by EIC than LHC

Conclusion

Philosophy and methodology

- BSM potential of EIC
- Model-independent SMEFT framework
- Semi-leptonic four-fermion operator sector
- Detailed accounting of anticipated uncertainties
- Simultaneous fits of Wilson coefficients with **beam polarization** and **luminosity difference**

Findings

- UV scales > 3 TeV with nominal annual luminosity
- UV scales > 4 TeV with $10 \times$ luminosity upgrade
- Strongest bounds from **polarized** electron + **unpolarized** proton scattering
- Complementary and competitive to LHC: EIC can
 - resolve degeneracies
 - impose even stronger bounds

EIC: designed as a QCD machine, also a powerful probe of BSM

The End